

## Radar observations in DART

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### 1. Introduction

The purpose of this document is to link theory on radar measurements and how these measurements are represented for the assimilation in the Data Assimilation Research Testbed (DART).

### 2. Radar observation location

What is usually known about radar observation is the position of the radar  $(\lambda_r, \phi_r, h_r)$ , the length of the path  $r$  between the target and the radar (also referred to as the range), and the azimuth and elevation angles of the electromagnetic beam as it leaves the radar  $(\alpha, \theta_e)$ . This information has to be translated into longitude, latitude, and height at the observation location. In the case of Doppler radial velocities, the orientation of the beam at the observation location also has to be estimated. This is the purpose of what follows.

Doviak and Zrnic (1993) give in their (2.28b-c) approximations for the height of the observation  $h$  (above sea level) and the great circle distance  $s$  (see Fig. 1):

$$h = \left[ r^2 + (k_e a)^2 + 2 r k_e a \sin(\theta_e) \right]^{1/2} - k_e a + h_r \quad (1)$$

$$s = k_e a \sin^{-1} \left( \frac{r \cos(\theta_e)}{k_e a + h} \right) \quad (2)$$

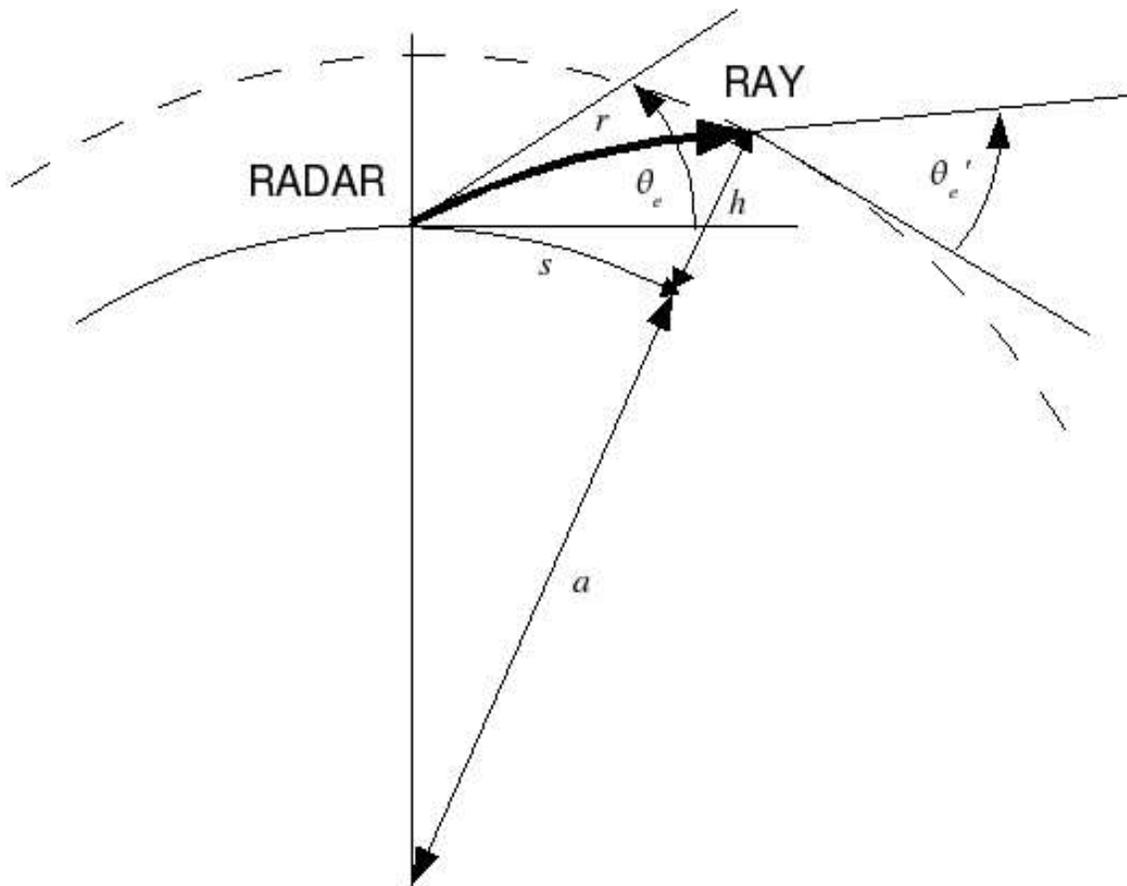
where  $r$  is the range,  $k_e = 4/3$  assumes that the vertical gradient of the refractive index is constant and equal to  $-1/4a$ ,  $a$  is the radius of the earth,

and  $\theta_e$  is the elevation angle of the beam as it leaves the radar. Note that Fig. 1 depicts the particular case  $h_r = 0$ ,  $h_r$  being the elevation of the radar above sea level. Let  $(\lambda_r, \phi_r)$  be the longitude and latitude at the radar location and  $(\lambda_o, \phi_o)$  be the longitude and latitude of the observation. For  $s \ll a$ , the following approximations hold:

$$a \cos\left(\frac{(\phi_r + \phi_o)}{2}\right)(\lambda_o - \lambda_r) \approx s \sin(\alpha) \quad (3)$$

$$a(\phi_o - \phi_r) \approx s \cos(\alpha) \quad (4)$$

where  $\alpha$  is the azimuth angle of the beam as it leaves the radar. The azimuth angle here is zero northward and positive clockwise. These expressions can be easily inverted to give  $(\lambda_o, \phi_o)$ .



*Figure 1 Circular path of a ray in an atmosphere in which the refractive index is linearly dependent on height.*

### **3. Radar observation operators**

To assimilate observations, a model is required to represent the observations in terms of state variables of the system. In the present section, the observation operators for radar reflectivity and Doppler velocity are described. It is assumed that radar data are point measurements but one has to keep in mind that real radar data represents measurements over a finite volume, weighted by the radar antenna pattern.

*a. Radar reflectivity*

The reflectivity factor for each hydrometeor category is calculated separately. The total reflectivity is simply the sum of the contributions from each hydrometeor category. The reflectivity contributed by ice crystals and clouds is assumed negligible. The reflectivity factors here are appropriate when the radar wavelength is much larger than the particle sizes.

The size distribution of the  $i^{\text{th}}$  hydrometeor class is assumed to be well approximated by an exponential function:

$$n_i(D) = n_{0i} \exp(-\lambda_i D) \quad (5)$$

where  $D$  is the particle diameter and the intercept parameter  $n_{0i}$  and the slope parameter of the size distribution  $\lambda_i$  are related to the mixing ratio of the species  $q_i$  by:

$$\lambda_i = \left( \frac{\pi \rho_i n_{0i}}{\rho q_i} \right)^{0.25} \quad (6)$$

where  $\rho_i$  is the density of the species and  $\rho$  is the air density. There are large uncertainties on  $\rho_i$  and the intercept parameters  $n_{0i}$  values. See Gilmore et al. (2004) for a discussion on this topic. Table 1 gives specific values for  $\rho_i$  and  $n_{0i}$ .

The effective reflectivity factor (the observed quantity) is

$$Z_e = C r^2 \frac{P_r}{P_t} \quad (7)$$

where  $P_t$  is the transmitted power,  $P_r$  is the received power,  $r$  is the range to the target, and  $C$  is the radar calibration coefficient. The effective reflectivity factor is modeled by a constant times the 6<sup>th</sup> moment of the size distribution:

$$Z_e \approx \sum_i c_i \int_0^{\infty} n_i(D) D^6 dD \quad (8)$$

According to theory, the backscattered energy is proportional to the 6<sup>th</sup> moment of the size distribution when all scatterers are spherical and the scattering is in the Rayleigh regime (radar wavelength  $\gg$  particle diameter). If all scatterers are spherical, then  $c_i = 1$ . In general, the measurement  $Z_e$  is associated with a mixture of hydrometeor types. In what follows, explicit expressions for reflectivity factor are given for rain, dry and wet graupel/hail, and dry and wet snow.

*i. Rain:*

$$c_r = 1; Z_r = \frac{7.2 \times 10^{20} (\rho q_r)^{1.75}}{\pi^{1.75} n_{0r}^{0.75} \rho_r^{1.75}} \quad (9)$$

*ii. Wet snow:*

$$c_{s, wet} = 1; Z_{s, wet} = \frac{7.2 \times 10^{20} (\rho q_s)^{1.75}}{\pi^{1.75} n_{0s}^{0.75} \rho_s^{1.75}} \quad (10)$$

*iii. Dry snow:*

$$c_{s,dry} = \left( \frac{|K_{ice}|^2}{|K_w|^2} \right) \left( \frac{\rho_s^2}{\rho_r^2} \right); Z_{s,dry} = \frac{c_{s,dry} 7.2 \times 10^{20} (\rho q_s)^{1.75}}{\pi^{1.75} n_{0s}^{0.75} \rho_s^{1.75}} \quad (11)$$

*iv. Wet graupel/hail:*

$$c_{g,wet} = \frac{(7.2 \times 10^{20})^{0.95}}{\Gamma(7)} (n_{0g})^{0.0375} \left( \frac{\pi \rho_g}{\rho q_g} \right)^{0.0875}; Z_{g,wet} = \left( \frac{7.2 \times 10^{20} (\rho q_g)^{1.75}}{\pi^{1.75} n_{0g}^{0.75} \rho_g^{1.75}} \right)^{0.95} \quad (12)$$

*v. Dry graupel/hail:*

$$c_{g,dry} = \left( \frac{|K_{ice}|^2}{|K_w|^2} \right) \left( \frac{\rho_g^2}{\rho_r^2} \right); Z_{g,dry} = \frac{c_{g,dry} 7.2 \times 10^{20} (\rho q_g)^{1.75}}{\pi^{1.75} n_{0g}^{0.75} \rho_g^{1.75}} \quad (13)$$

Here  $|K^2|$  is the dielectric factor. Expressions for dry particles are applied for temperatures below 0°C and expressions for wet particles are applied for temperatures above 0°C. Hence, a vertical discontinuity in reflectivity is expected at the melting layer. To avoid this problem, one can define a transition zone where the amount of dry and wet particles changes continuously, from dry particles only to wet particles only as the temperature increases and crosses the melting point. The expression for wet graupel/hail (12) is elevated to the 0.95 power to take into account Mie effect (Smith, 1984) and is appropriate for 10-cm radars. For details about these 5 expressions, see Smith et al. (1975) and Smith (1984). Additional references can be found in Ferrier (1994).

*b. Doppler velocity*

The radial wind sample from the model is computed as follows:

$$v_r = d_x u + d_y v + d_z (w - w_t) \quad (14)$$

where  $(u, v, w)$  are the zonal, meridional, and vertical wind components,  $w_t$  is the terminal fall speed of the radar scatterers, and

$$d_x \equiv \sin(\alpha') \cos(\theta_e') \quad (15)$$

$$d_y \equiv \cos(\alpha') \cos(\theta_e') \quad (16)$$

$$d_z \equiv \sin(\theta_e') \quad (17)$$

where the primes refer to the azimuth and elevation angles at the observation location. The direction vector  $d$  is calculated before the assimilation for each observation and stored in the observation file. This strategy is employed to avoid superfluous calculation during the assimilation. The azimuth angle of the observed velocity component is approximated to the azimuth angle of the beam as it leaves the radar ( $\alpha' = \alpha$ ). For the elevation angle of the observed velocity component, we follow Battan's (1973) equation (3.18a) and assume the same vertical gradient of the refractive index as in section 2:

$$\theta_e' \approx \left[ \frac{1.5h}{a} + \theta_e^2 \right]^{1/2} \quad (18)$$

This approximation is appropriate for small elevation angles.

The terminal fall speeds for precipitating particle of diameter  $D$  for rain, snow, and graupel are, respectively,

$$U_r = aD^b \left( \frac{\rho_0}{\rho} \right)^{1/2} \quad (19)$$

$$U_s = cD^d \left( \frac{\rho_0}{\rho} \right)^{1/2} \quad (20)$$

$$U_g = \left( \frac{4g\rho_g}{3C_D\rho} \right)^{1/2} D^{1/2} . \quad (21)$$

For references, see Lin et al, (1983). We define  $w_t$  the reflectivity-weighted mean terminal speeds as

$$w_t = \frac{\sum_i c_i \int_0^\infty U_i n_i(D) D^6 dD}{\sum_i Z_i} = \frac{\sum_i \omega_i}{\sum_i Z_i} \quad (22)$$

where the  $\omega_i$  are the reflectivity-convoluted mean terminal speeds for each hydrometeor category. The 5 terms in the denominator are given by equations (9) to (13). For the numerator in (22), we have

*i. Rain:*

$$\omega_r = 10^{18} n_{0r} a \left( \frac{\rho_0}{\rho} \right)^{1/2} \Gamma(b+7) \left( \frac{\rho q_r}{\pi n_{0r} \rho_r} \right)^{\left( \frac{b+7}{4} \right)} \quad (23)$$

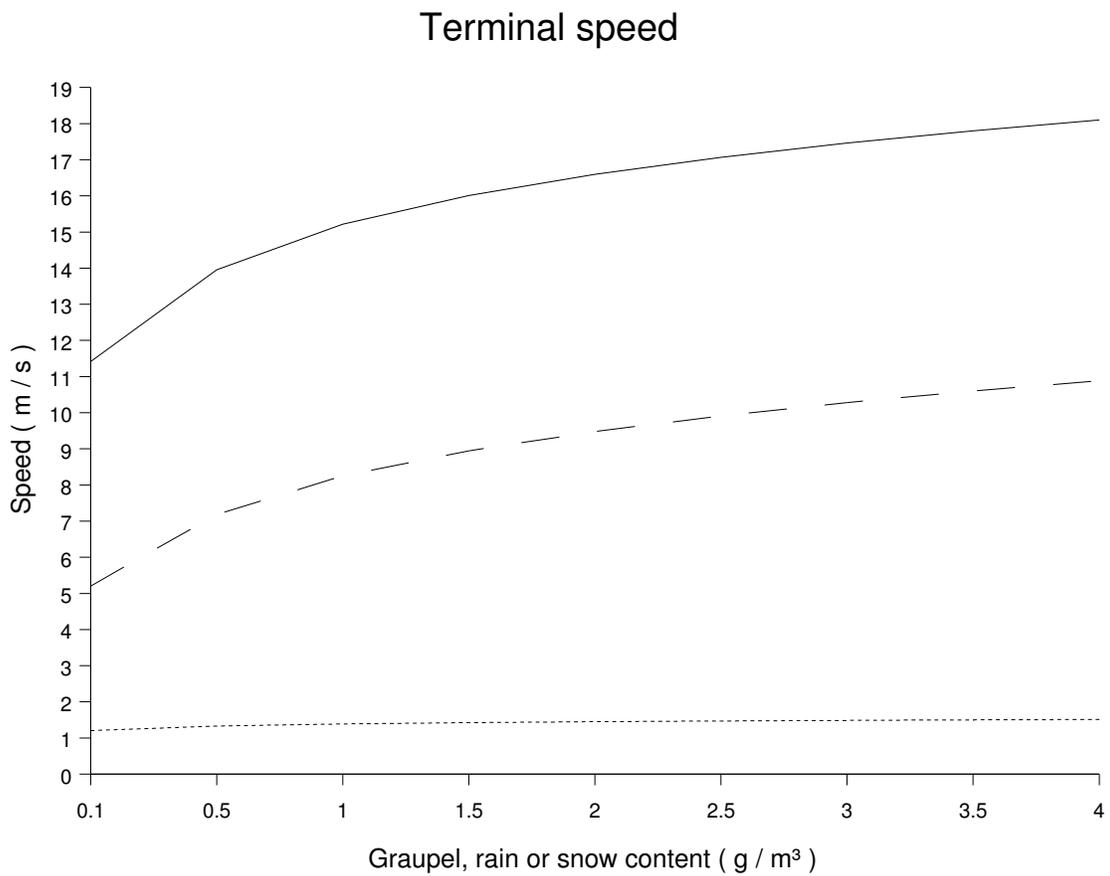
*ii. Snow:*

$$\omega_{s, dry/wet} = 10^{18} c_{s, dry/wet} n_{0s} c \left( \frac{\rho_0}{\rho} \right)^{1/2} \Gamma(d+7) \left( \frac{\rho q_s}{\pi n_{0s} \rho_s} \right)^{\left( \frac{d+7}{4} \right)} \quad (24)$$

iii. *Graupel / Hail:*

$$\omega_{g, dry/wet} = 10^{18} c_{g, dry/wet} n_{0g} \left( \frac{4g\rho_g}{3C_D\rho} \right)^{1/2} \Gamma(7.5) \left( \frac{\rho q_g}{\pi n_{0g} \rho_g} \right)^{(1.875)} \quad (25)$$

where the factor  $10^{18}$  takes into account that the reflectivity in the denominator has  $\text{mm}^6 \text{m}^{-3}$  units. The reflectivity-weighted mean terminal speeds as functions of precipitation content are shown in Fig. 2. For references on reflectivity-weighted mean terminal speeds derivations, see Doviak and Zrnica (1993) p.275, Hauser and Amayenc (1981), and Conway and Zrnica (1993) in their appendix.



*Figure 2 Reflectivity-weighted mean terminal speeds for graupel (full line), rain (dashed line) and snow (dotted line).*

#### LIST OF SYMBOLS

<i>Notation</i>	<i>Description</i>	<i>DART value [range]</i>	<i>Units</i>
<i>a</i>	constant in empirical formula for $U_r$	842	$\text{m}^{(1-b)} \text{s}^{-1}$
<i>b</i>	constant in empirical formula for $U_r$	0.8	
<i>c</i>	constant in empirical formula for $U_s$	4.84	$\text{m}^{(1-d)} \text{s}^{-1}$

<i>Notation</i>	<i>Description</i>	<i>DART value [range]</i>	<i>Units</i>
$C_D$	drag coefficient for hailstone	0.6	
$d$	constant in empirical formula for $U_s$	0.25	
$g$	gravitational acceleration	9.81	$\text{m s}^{-2}$
$K_{ice}^2 / K_w^2$	dielectric factor ratio	0.224 (0.189 or 0.224) <sup>†</sup>	
$n_{or}$	intercept parameter of the raindrop size distribution	$8 \times 10^6$	$\text{m}^{-4}$
$n_{og}$	intercept parameter of the graup size distribution	$4 \times 10^4$ ‡ [ $4 \times 10^3, 4 \times 10^6$ ]	$\text{m}^{-4}$
$n_{os}$	intercept parameter of the snow size distribution	$3 \times 10^6$	$\text{m}^{-4}$
$q$	mixing ratio		$\text{kg kg}^{-1}$
$Z_e$	effective reflectivity factor		$\text{mm}^6 \text{m}^{-3}$
$\rho$	air density		$\text{kg m}^{-3}$
$\rho_0$	surface air density	1	$\text{kg m}^{-3}$
$\rho_g$	density of graupel	$917$ ‡ [400, 900]	$\text{kg m}^{-3}$
$\rho_s$	density of snow	100	$\text{kg m}^{-3}$
$\rho_r$	density of water	1000	$\text{kg m}^{-3}$

† According to Smith (1984), there are two choices for the dielectric factor, depending on how the snow particle sizes are specified. If melted raindrop diameters are used, then the factor is 0.224. If equivalent ice sphere diameters are used, then the factor is 0.189.

‡ From Lin et al. (1983).

## APPENDIX

$$\int_0^{\infty} x^{\nu-1} e^{-\mu x} dx = \frac{1}{\mu^{\nu}} \Gamma(\nu), R(\mu) > 0, R(\nu) > 0 \quad (\text{A1})$$

The gamma function  $\Gamma$  is related to factorials when the argument is an integer:

$$\Gamma(n) = (n-1)! \quad (\text{A2})$$

## REFERENCES

- Battan, L. J., 1973: *Radar Observation of the Atmosphere*. Univ. of Chicago Press, 324 pp.
- Conway, J. W., and D. S. Zrnice, 1993: A study of embryo production and hail growth using dual-Doppler and multiparameter radars. *Mon. Wea. Rev.*, **121**, 2511-2528.
- Doviak, R. J., and D. S. Zrnice, 1993: *Doppler Radar and Weather Observations*. Academic Press, 562 pp.
- Ferrier, B. S., 1994: A double-moment multiple-phase four-class bulk ice scheme. Part I: Description. *J. Atmos. Sci.*, **51**, 249-280.
- Gilmore, S. G., Straka J. M., and E. R. Rasmussen, 2004: Precipitation uncertainty due to variations in precipitation particle parameters within a simple microphysics scheme. *Mon. Wea. Rev.*, **132**, 2610-2627.
- Hauser, D., and P. Amayenc, 1981: A new method for deducing hydrometeor size distributions and vertical air motions from Doppler radar measurements at vertical incidence. *J. Appl. Meteor.*, **20**, 547-555.
- Lin, Y.-L., Farley R. D., and H. D. Orville, 1983: Bulk parameterization of

- the snow field in a cloud model. *J. Climate Appl. Meteor.*, **22**, 1065-1092.
- Smith, P. L. Jr., 1984: Equivalent radar reflectivity factors for snow and ice particles. *J. Climate Appl. Meteor.*, **23**, 1258-1260.
- Smith, P. L. Jr., Myers C. G., and H. D. Orville, 1975: Radar reflectivity factor calculations in numerical cloud models using bulk parameterization of precipitation. *J. Appl. Meteor.*, **14**, 1156-1165.