
Models of the Probability Distribution of Sea-Surface Wind Speeds

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Outline

- **Motivation**

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- **Characterisation of wind speed pdfs**

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- **Conclusions**

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- Sea surface winds (“eddy averaged”) are an important determinant of turbulence in both boundary layers
- Turbulence feeds back on surface winds, primarily through surface momentum flux



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- Calculation of averaged fluxes requires full surface wind pdf

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⇒ higher order vector wind moments may affect lower order wind speed moments

Skewness and Kurtosis

- **Skewness:** measure of asymmetry of pdf

$$\text{skew}(x) = \left\langle \left(\frac{x - \langle x \rangle}{\text{std}(x)} \right)^3 \right\rangle$$

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- **Kurtosis:** measure of flatness of pdf

$$\text{skew}(x) = \left\langle \left(\frac{x - \langle x \rangle}{\text{std}(x)} \right)^4 \right\rangle - 3$$

Sea-Surface Winds: Notation

Notation:

- \mathbf{u} vector wind (u, v)
- u along mean wind component
- v cross mean wind component
- w wind speed $(u^2 + v^2)^{1/2}$

Sea-Surface Winds: Data

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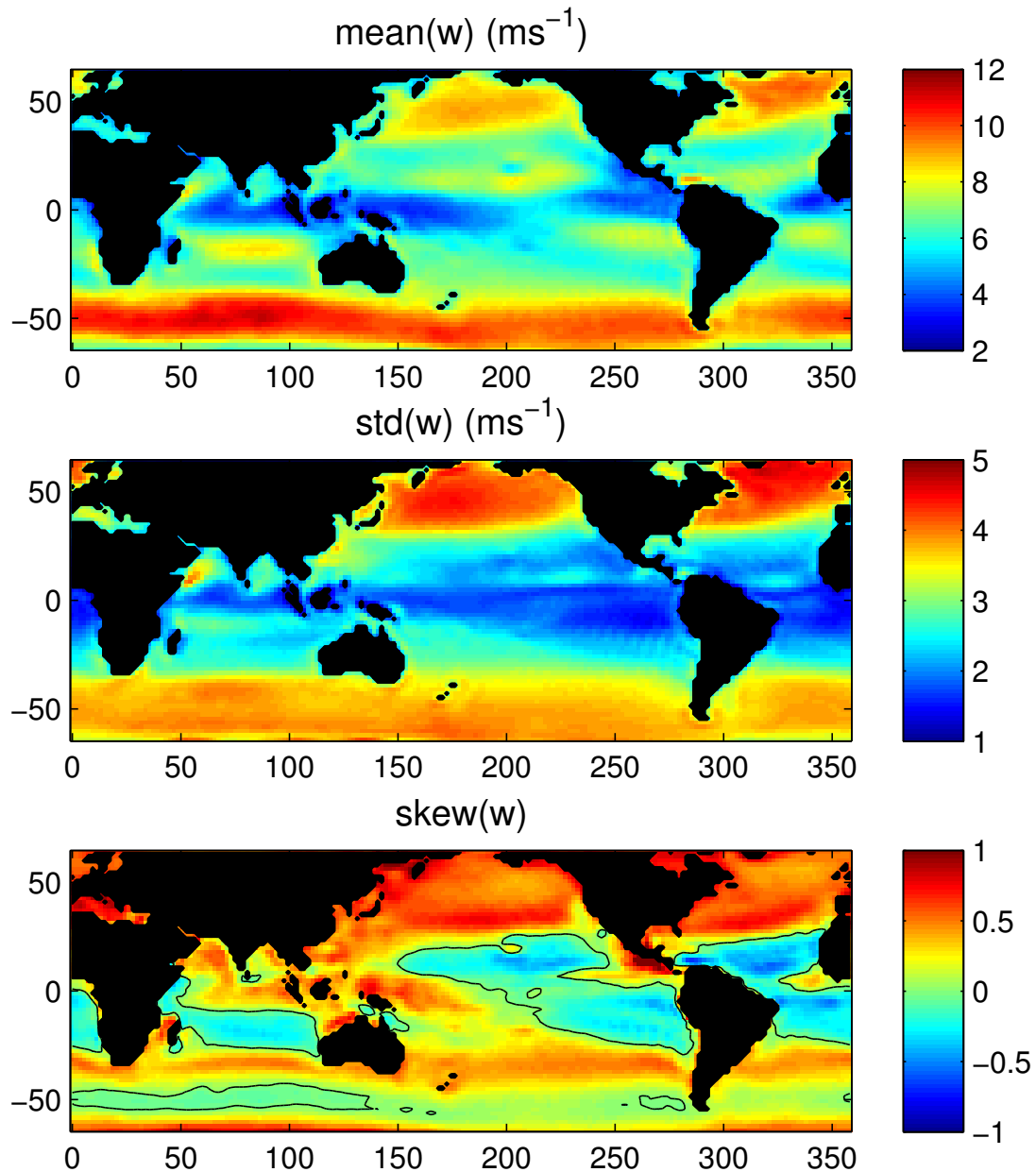
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 - Hourly buoy 10m winds

Sea-Surface Wind Speeds: Moments



Wind Speed pdfs: Weibull

- The pdf of wind speed w has traditionally been represented by 2-parameter Weibull distribution:

$$p_w(w) = \begin{cases} \frac{b}{a} \left(\frac{w}{a}\right)^{b-1} \exp\left(-\left(\frac{w}{a}\right)^b\right) & w > 0 \\ 0 & w < 0 \end{cases}$$

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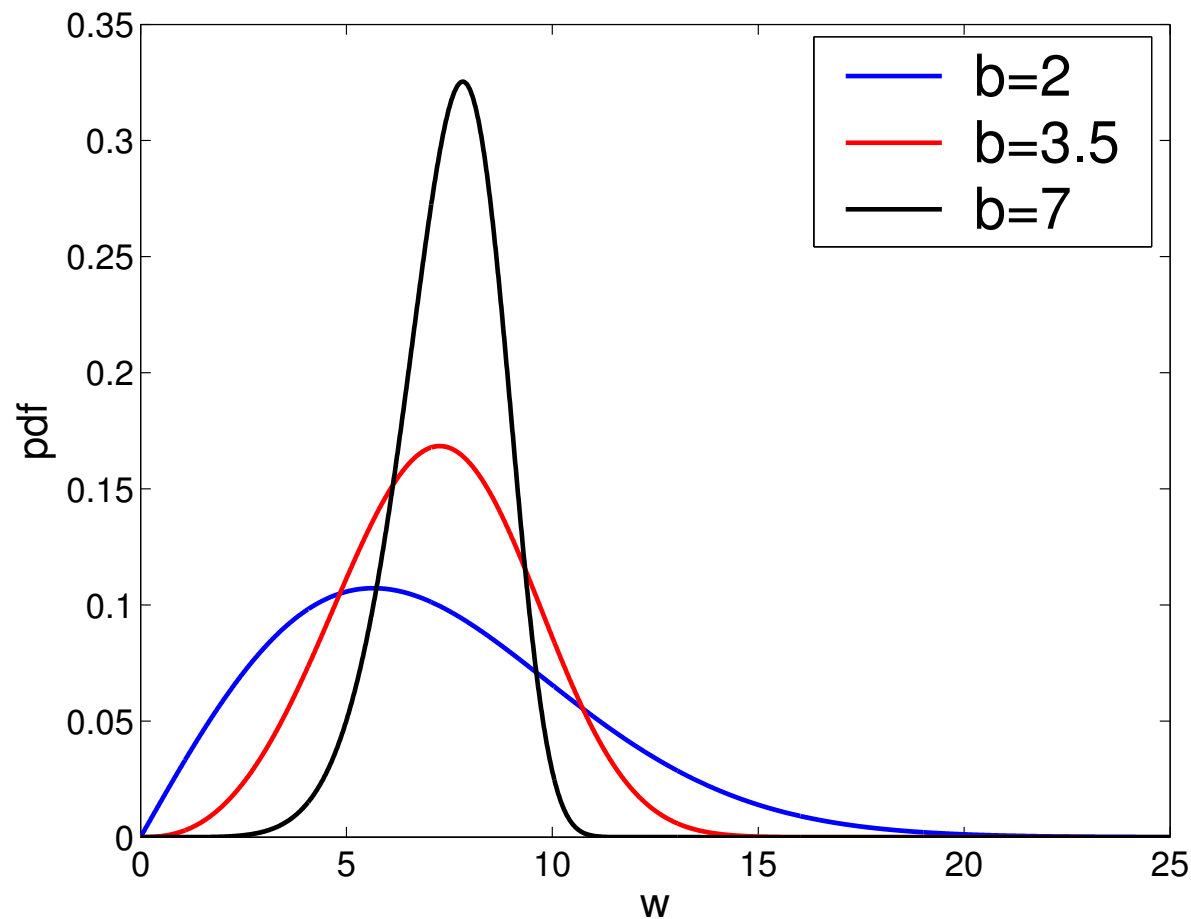
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- b is the shape parameter (pdf tilt)
- $p_w(w)$ is unimodal

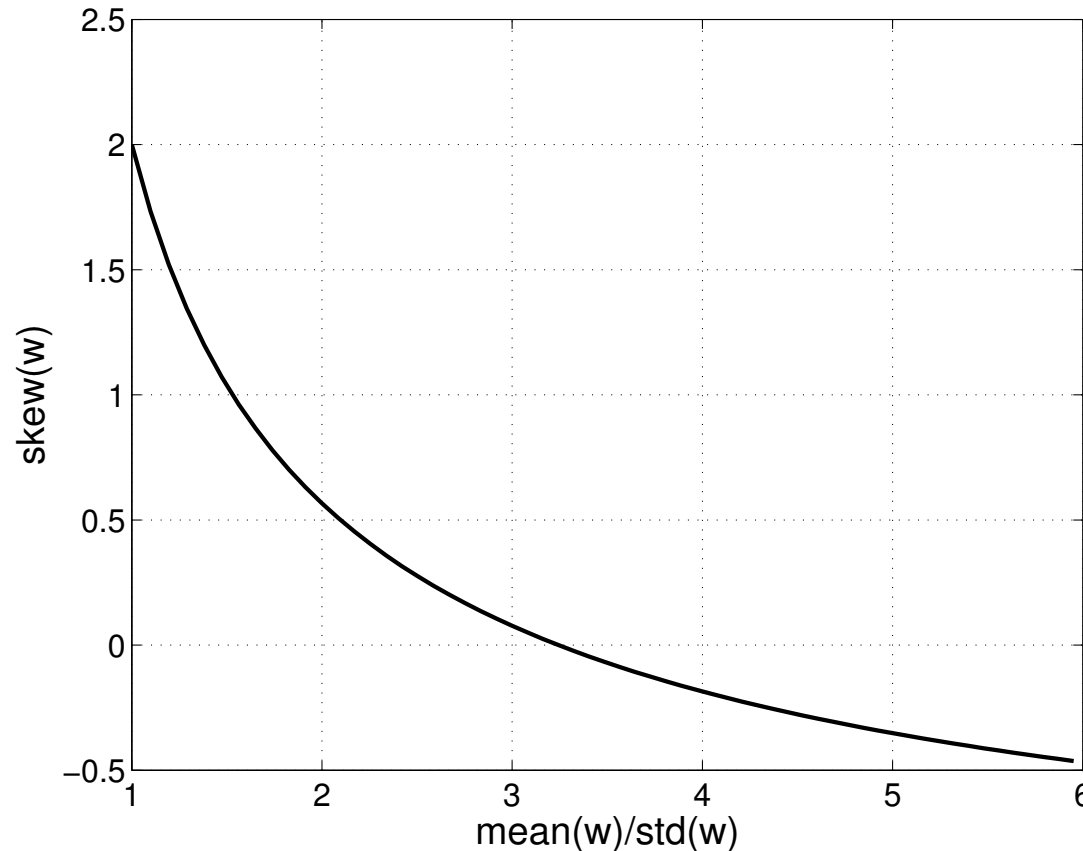
Wind Speed pdfs: Weibull

- Weibull pdfs for $a = 8$



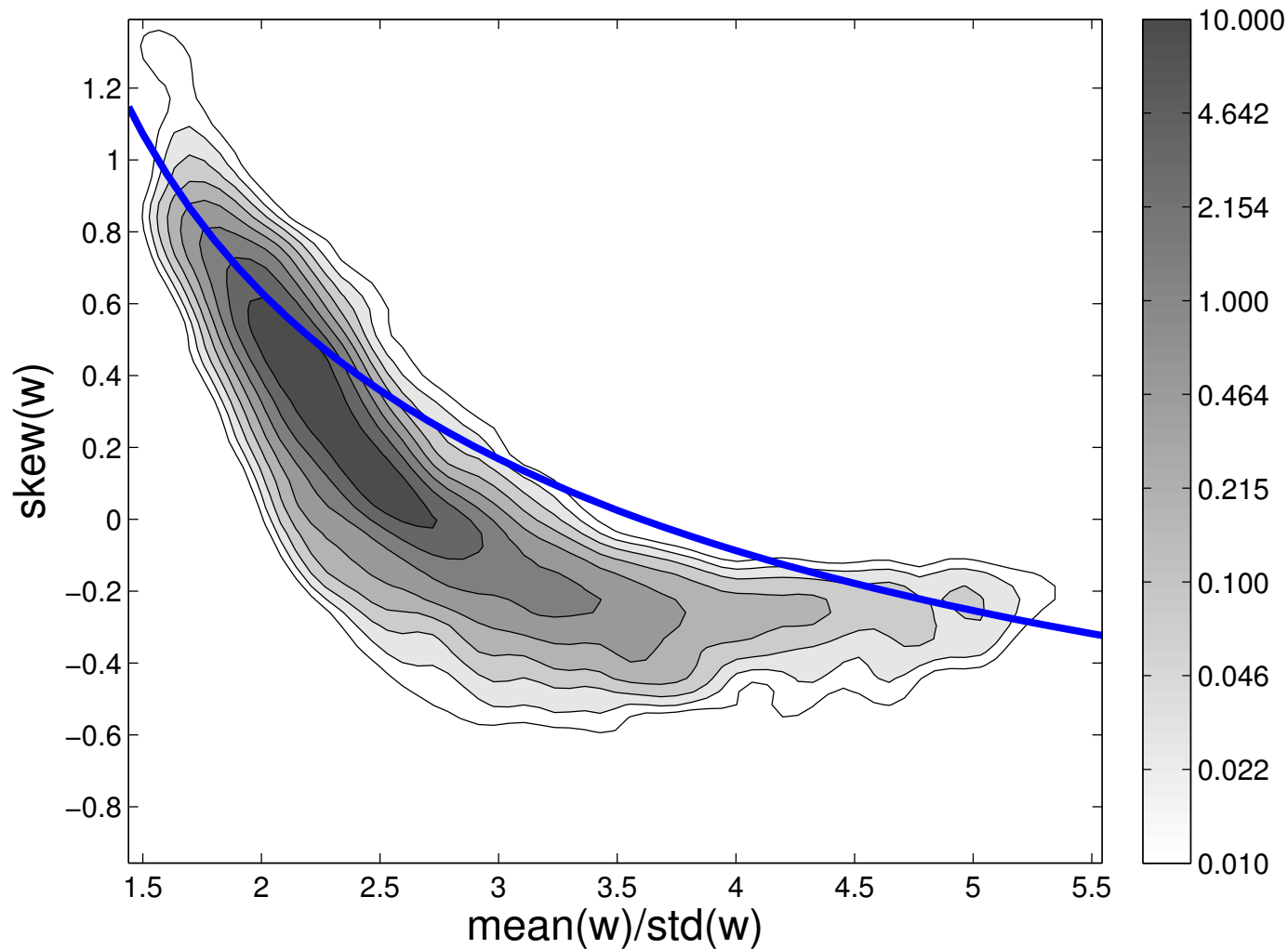
Wind Speed pdfs: Weibull

- For a Weibull distribution, $\text{skew}(w)$ is a decreasing function of $\text{mean}(w)/\text{std}(w)$



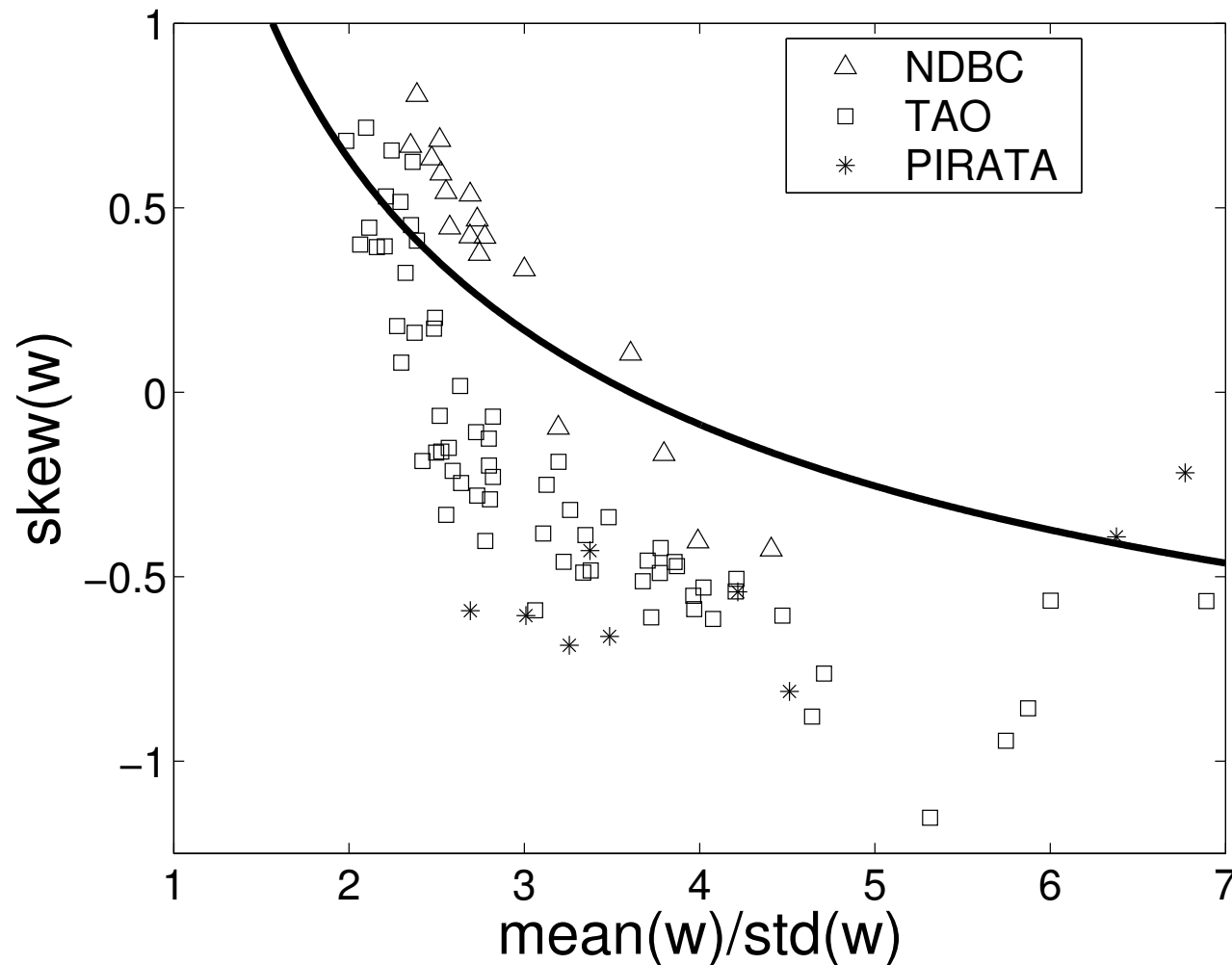
Wind Speed pdfs: Observed

- Observed speed moments fall around Weibull curve



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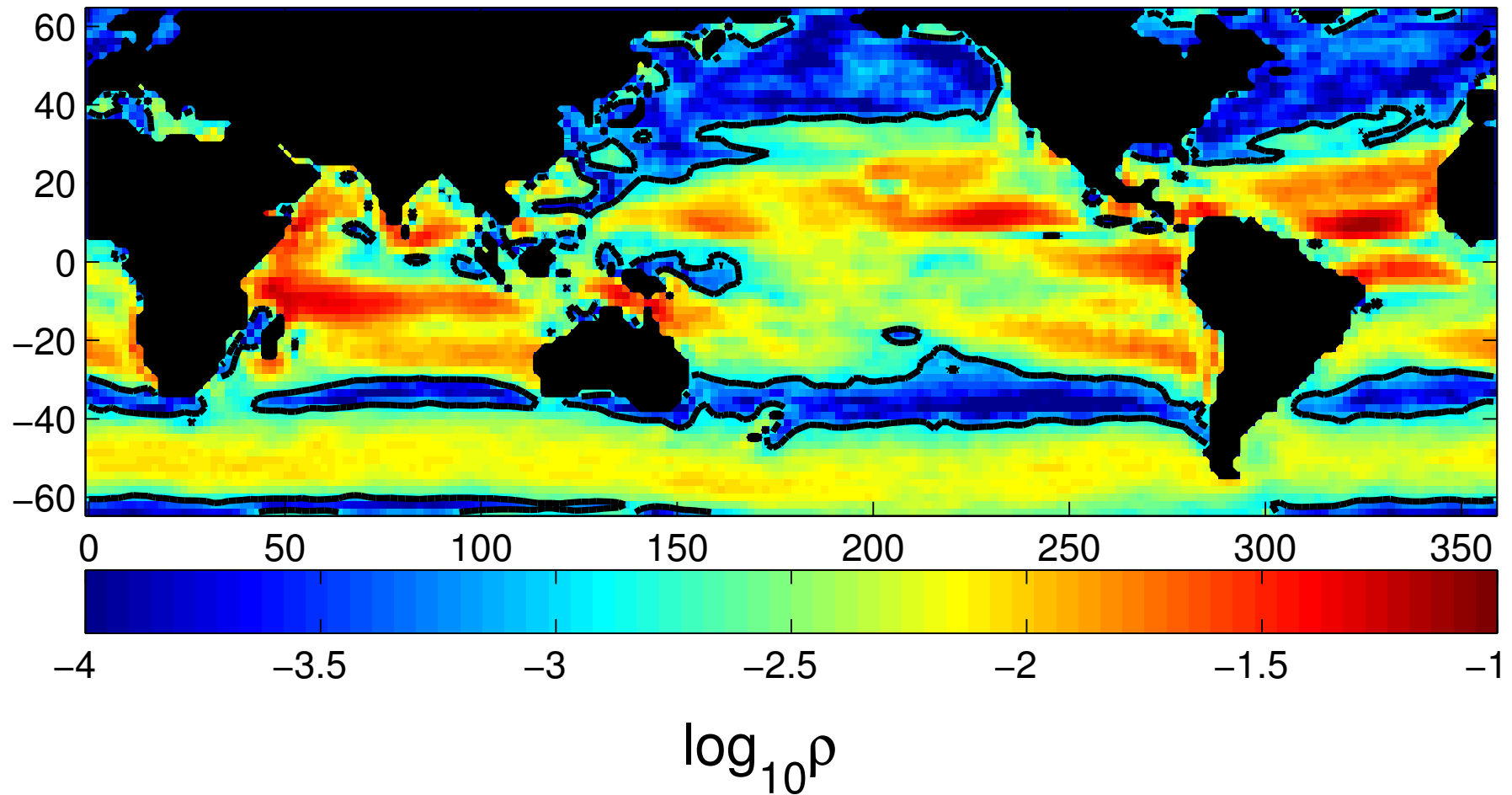
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where:

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$q_w(w)$ = best-fit Weibull pdf

Wind Speed pdfs: Relative Entropy



Wind Speed pdfs: Empirical Models

- Strategy: systematically construct wind speed pdfs from joint pdf of vector components, $p_{uv}(u, v)$

$$p_w(w) = w \int_{-\pi}^{\pi} p_{uv}(w \cos \theta, w \sin \theta) d\theta$$

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- Approach simplified by assuming u, v independent, so

$$p_{uv}(u, v) = p_u(u)p_v(v)$$

Wind Speed pdfs: Gaussian Vector Winds

- Simplest model assumes isotropic Gaussian fluctuations in vector winds:

$$p_u(u) = \frac{1}{\sqrt{2\pi\sigma^2}} \exp\left(-\frac{(u - \bar{u})^2}{2\sigma^2}\right)$$
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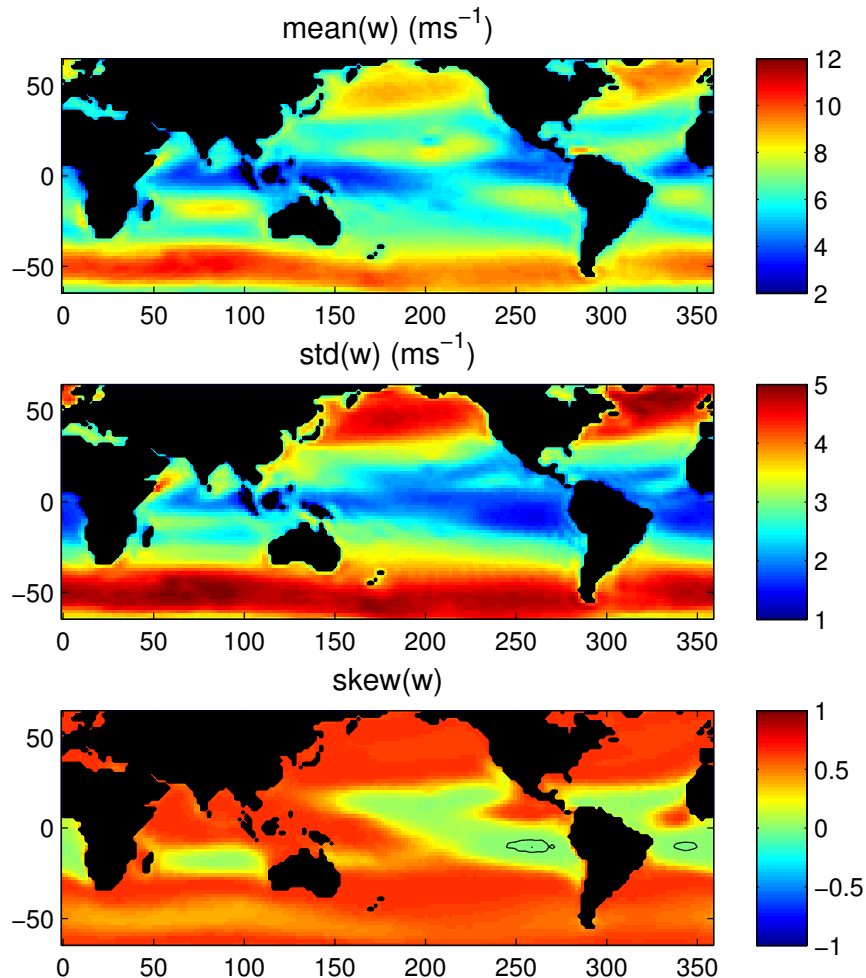
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- Integrating over wind direction:

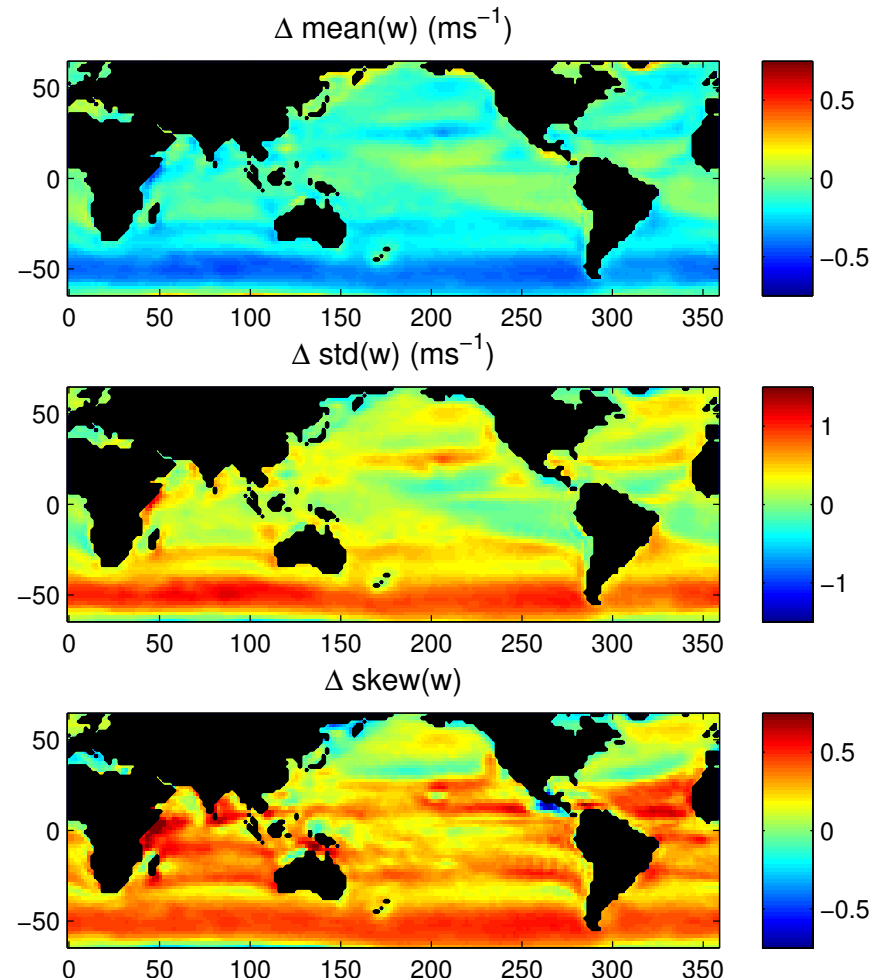
$$p_w(w) = \frac{w}{\sigma^2} \exp\left(-\frac{w^2 + \bar{u}^2}{2\sigma^2}\right) I_0\left(\frac{w\bar{u}}{\sigma^2}\right)$$

Wind Speed pdfs: Gaussian Vector Winds

Model



Model-Observations



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- Relax assumption of isotropy:
 $\sigma_u = \text{std}(u), \sigma_v = \text{std}(v) \Rightarrow$

$$p_w(w) = \frac{w}{\sigma_u \sigma_v} \exp\left(-\frac{w^2 + \bar{u}^2}{2\sigma_u^2}\right) \times \left\{ I_0\left(\frac{w\bar{u}}{\sigma_u^2}\right) + \sum_{k=1}^{\infty} \left[\frac{w}{\bar{u}} \left(1 - \frac{\sigma_u^2}{\sigma_v^2}\right) \right]^k \frac{\Gamma(k + 1/2)}{\sqrt{\pi} k!} I_k\left(\frac{\bar{u}w}{\sigma_u^2}\right) \right\}$$

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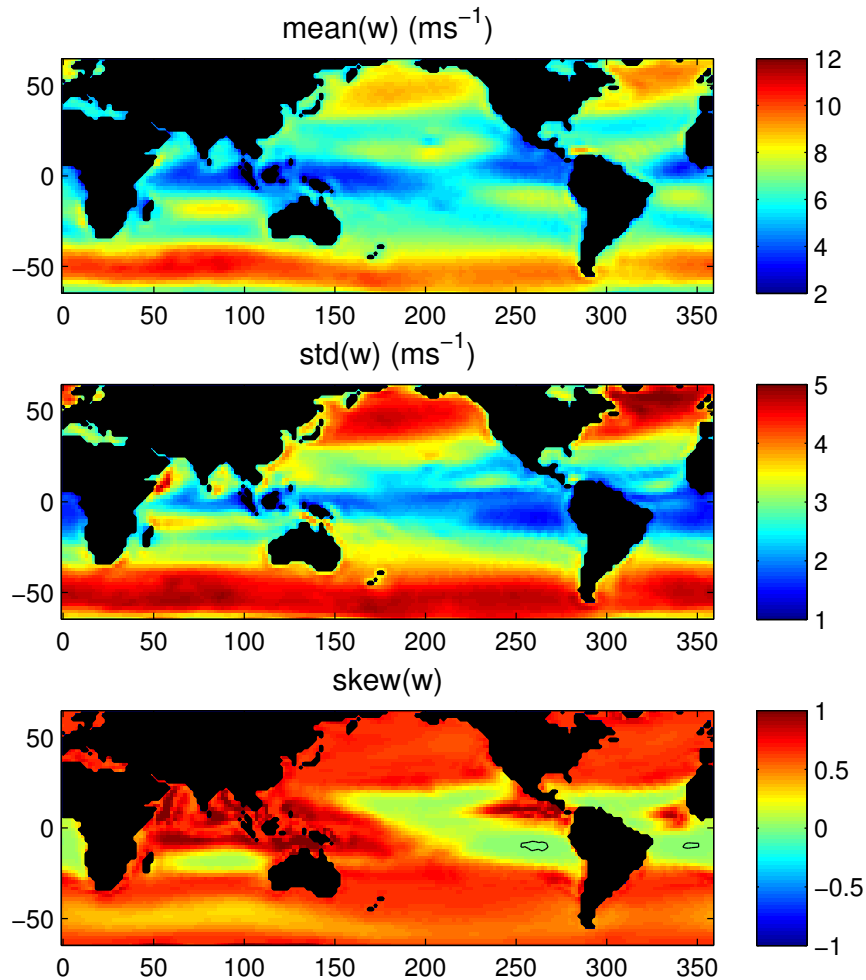
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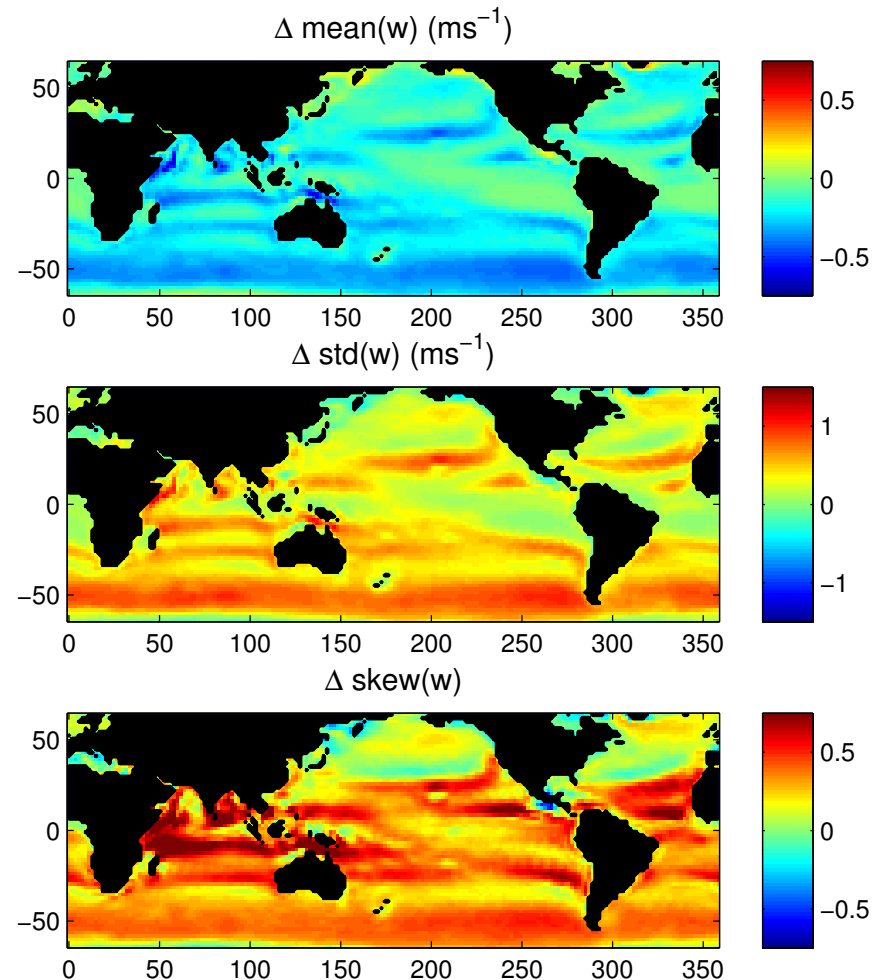
- Actually makes approximation somewhat *worse*

Wind Speed pdfs: Gaussian Vector Winds

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Model-Observations



UVic

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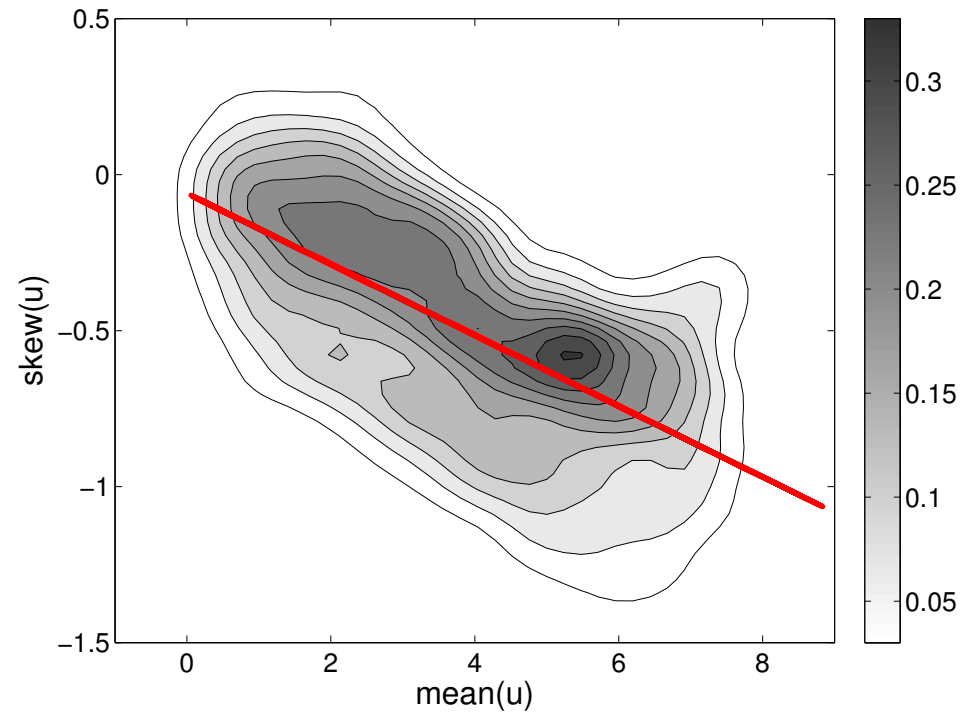
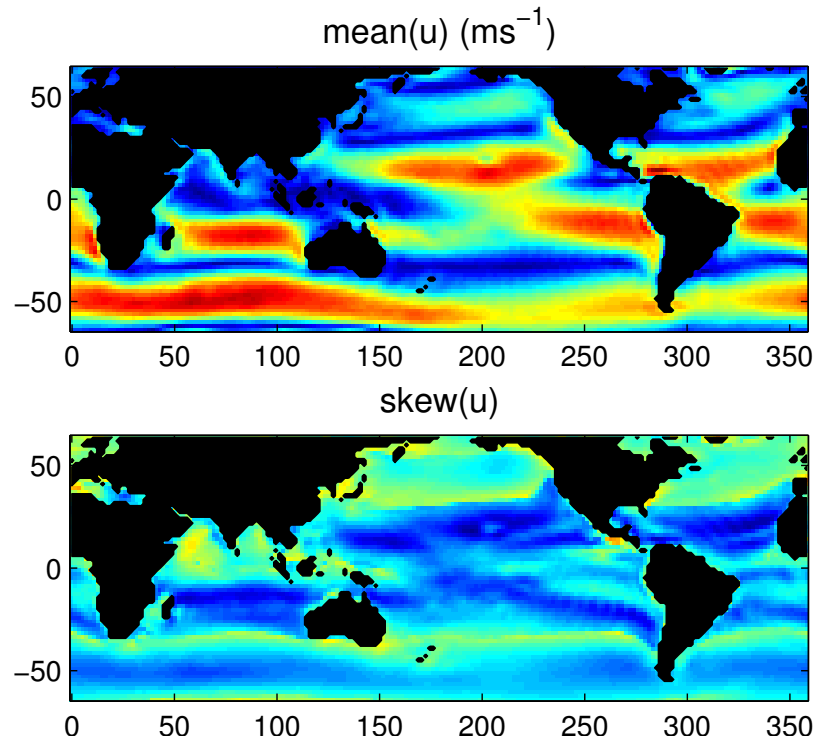
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- Coupling of moments follows from nonlinear dependence of surface drag on surface winds
- Symmetric fluctuations in forcing
⇒ asymmetric response, skewed toward rest

Vector Wind Moments



$$\text{skew}(u) \simeq (-0.11 \text{ s m}^{-1}) \text{mean}(u) - 0.06$$

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- Question is: how important are details of skewed pdf of u for pdf of w ?

Bigaussian pdf

- Two half-Gaussians:

$$p_u(u) = \frac{1}{\sqrt{2\pi\sigma^2}} \begin{cases} \exp\left(-\frac{(u-\mu)^2}{2\sigma^2(1-\epsilon)^2}\right) & u < \mu \\ \exp\left(-\frac{(u-\mu)^2}{2\sigma^2(1+\epsilon)^2}\right) & u > \mu \end{cases}$$

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- Moments:

$$\text{mean}(u) = \mu + \sqrt{\frac{8}{\pi}} \sigma \epsilon$$

$$\text{std}(u) = \sigma \left[1 + \left(3 - \frac{8}{\pi} \right) \epsilon^2 \right]^{1/2}$$

$$\text{skew}(u) \approx \sqrt{\frac{8}{\pi}} \frac{\epsilon}{\text{std}^3(u)}$$

Centred Gamma pdf

- With $z = (u - \bar{u})/\sigma$:

$$p_u(u) = \begin{cases} \frac{|\beta|}{\sigma\Gamma(\beta^2)} [\beta(z + \beta)]^{\beta^2-1} \exp[-\beta(z + \beta)] & z + \beta > 0 \\ 0 & z + \beta < 0 \end{cases}$$

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- Moments:

$$\begin{aligned} \text{mean}(u) &= \bar{u} \\ \text{std}(u) &= \sigma \\ \text{skew}(u) &= 2/\beta \end{aligned}$$

Gram-Charlier Expansion

- With $z = (u - \bar{u})/\sigma$:

$$p_u(u) = \frac{1}{\sqrt{2\pi\sigma^2}} \left[1 + \frac{\nu}{6} H_3(z) + \frac{\kappa}{24} H_4(z) \right] \exp\left(-\frac{z^2}{2}\right)$$

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$$\text{kurt}(u) = \kappa$$

Maximum Entropy pdf

- Subject to the constraints:

$$\begin{aligned}\text{mean}(u) &= \bar{u} \\ \text{std}(u) &= \sigma \\ \text{skew}(u) &= \nu \\ \text{kurt}(u) &= \kappa\end{aligned}$$

find the pdf $p_u(u)$ which maximises the entropy

$$H = - \int p_u(u) \ln p_u(u) du$$

Maximum Entropy pdf

- Solution takes the form

$$p_u(u) = \frac{1}{Z} \exp \left(\sum_{i=1}^4 \lambda_i u^i \right)$$

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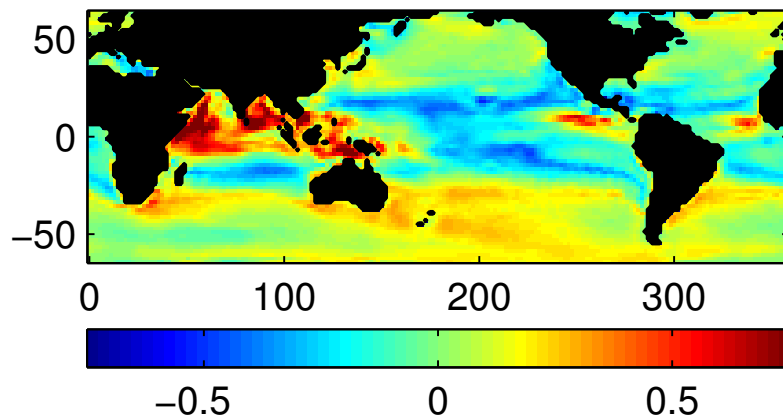
$$p_u(u) = \frac{1}{Z} \exp \left(\sum_{i=1}^4 \lambda_i u^i \right)$$

- Lagrange multipliers $\{\lambda_i\}$ found as solutions to unconstrained dual variational problem
- Maximum entropy pdf is “least biased” among all pdfs with given moments, in a rigorous information theoretic sense

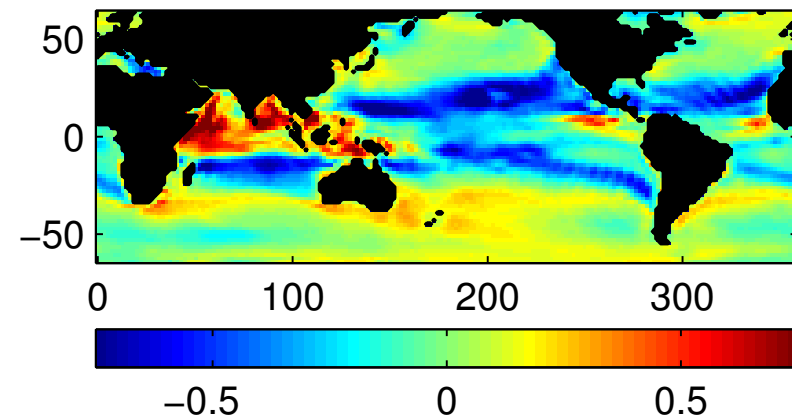
Empirical Wind Speed pdfs: Intercomparison

- Skew(u) from observations, no kurt(u) information

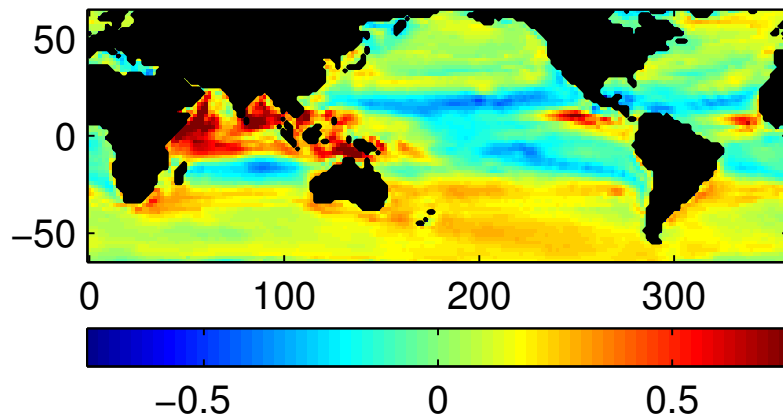
Bigaussian



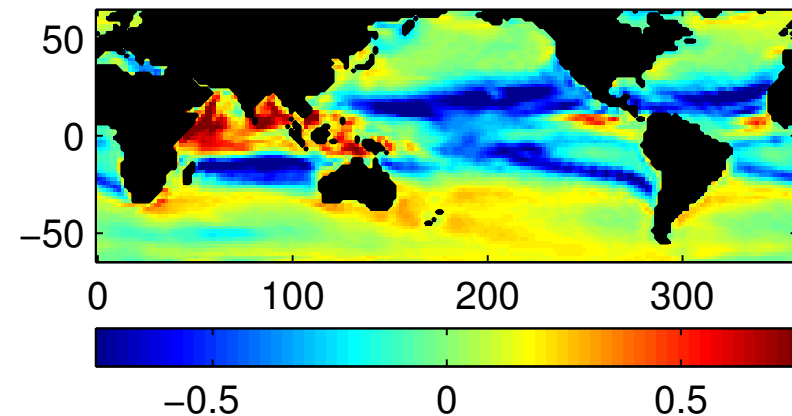
Gram-Charlier



Centred Gamma



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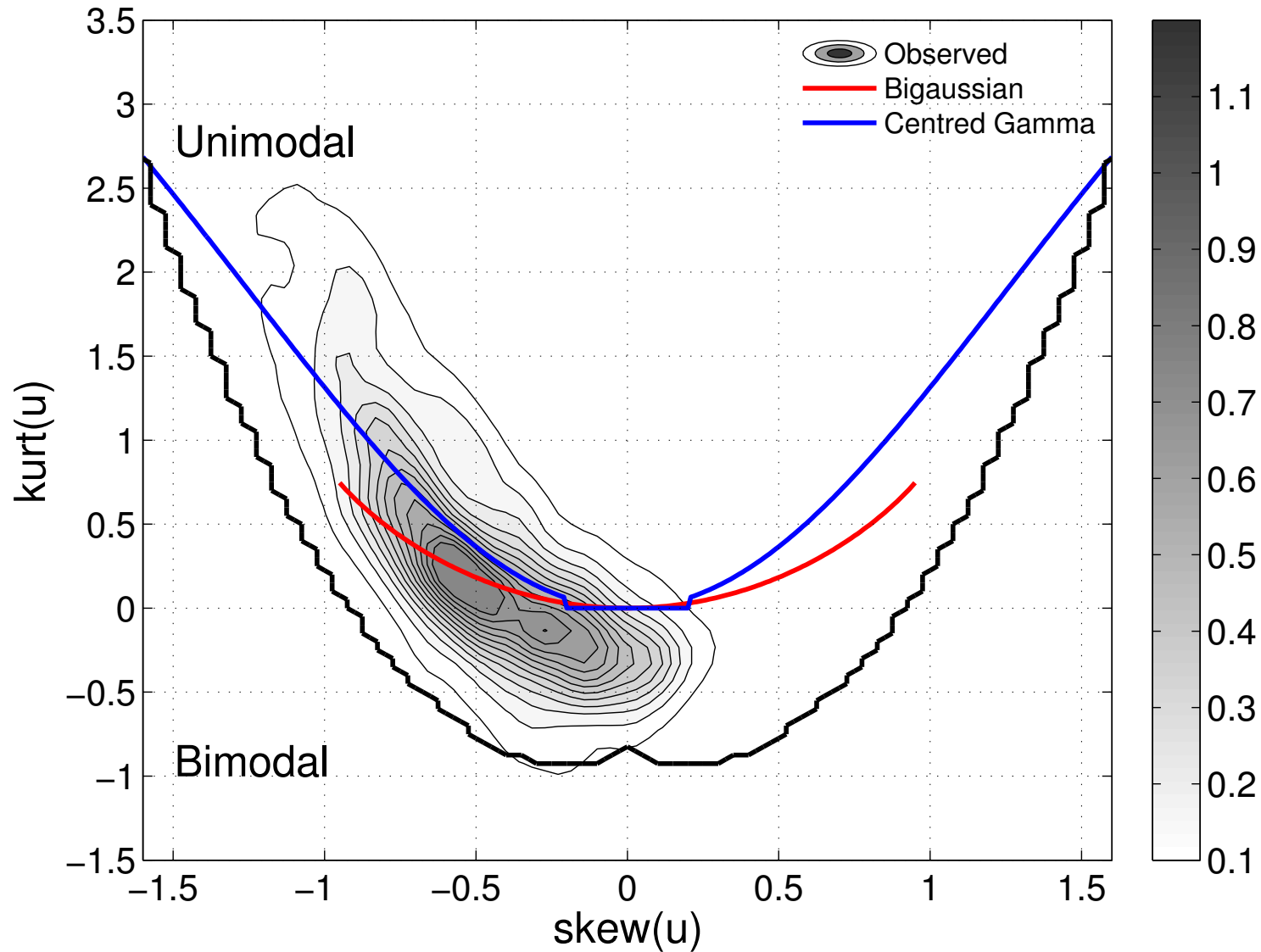
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- **BUT:** for Gram-Charlier & maximum entropy pdfs, have set $\text{kurt}(u) = 0$

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- In fact, kurtosis of along-mean-wind component non-trivial

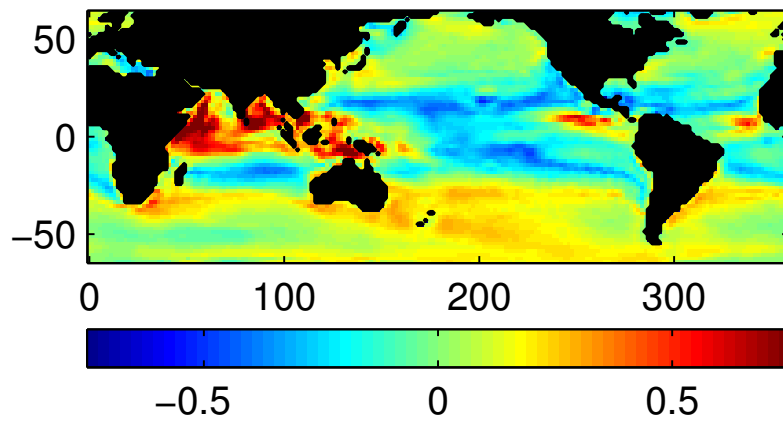
Skew(u) vs. Kurt(u)



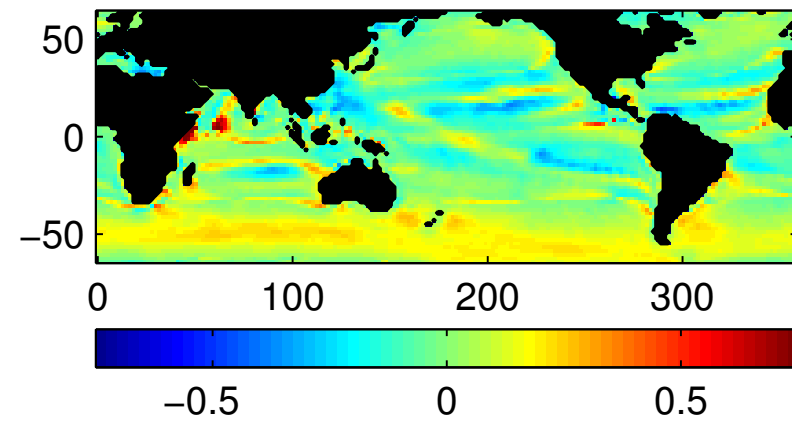
Empirical Wind Speed pdfs: Intercomparison

■ Skew(u) and kurt(u) from observations

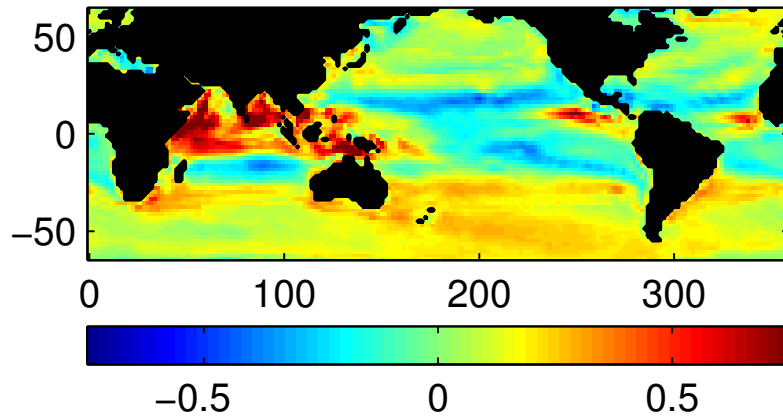
Bigaussian



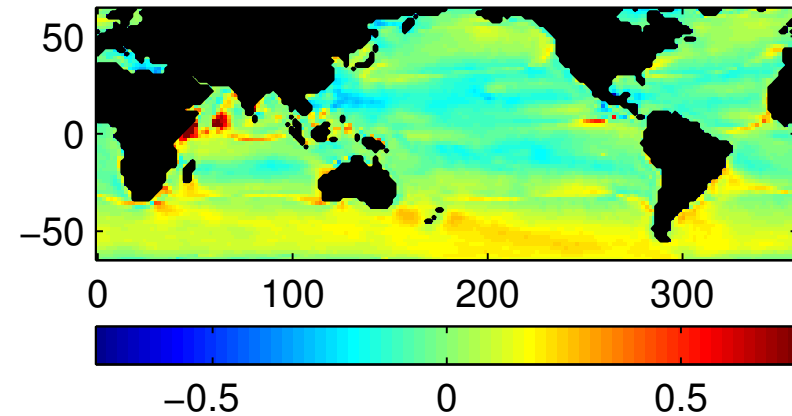
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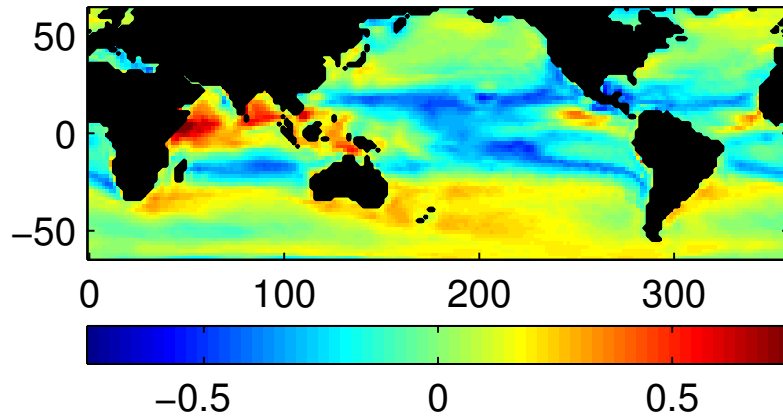
Maximum Entropy



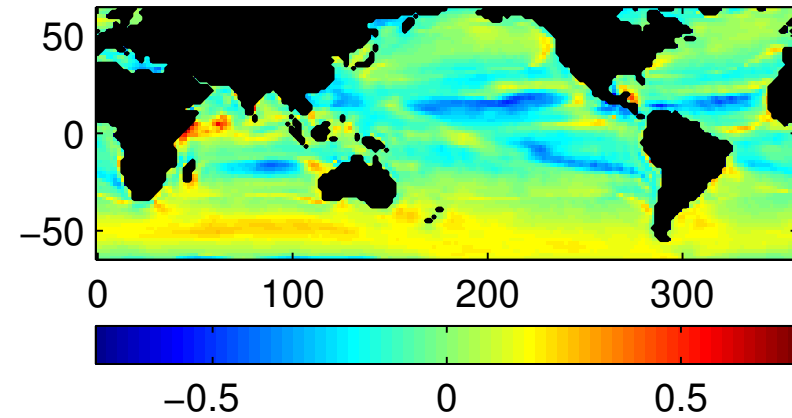
Empirical Wind Speed pdfs: Intercomparison

■ Isotropic fluctuations

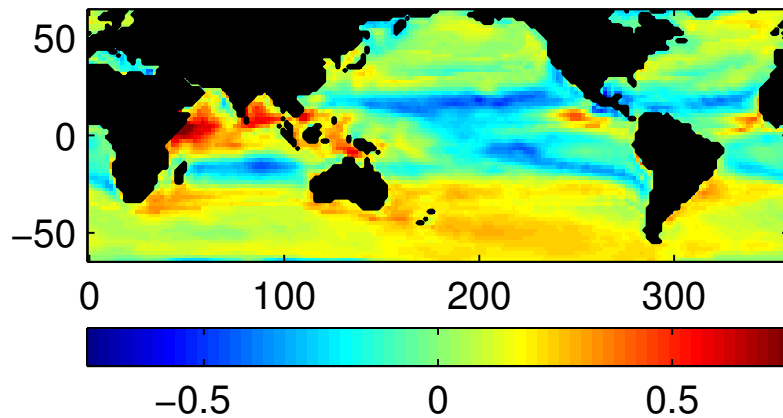
Bigaussian



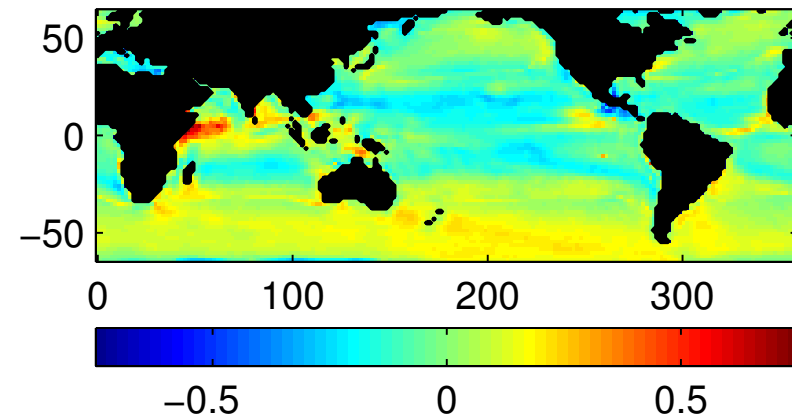
Gram-Charlier



Centred Gamma

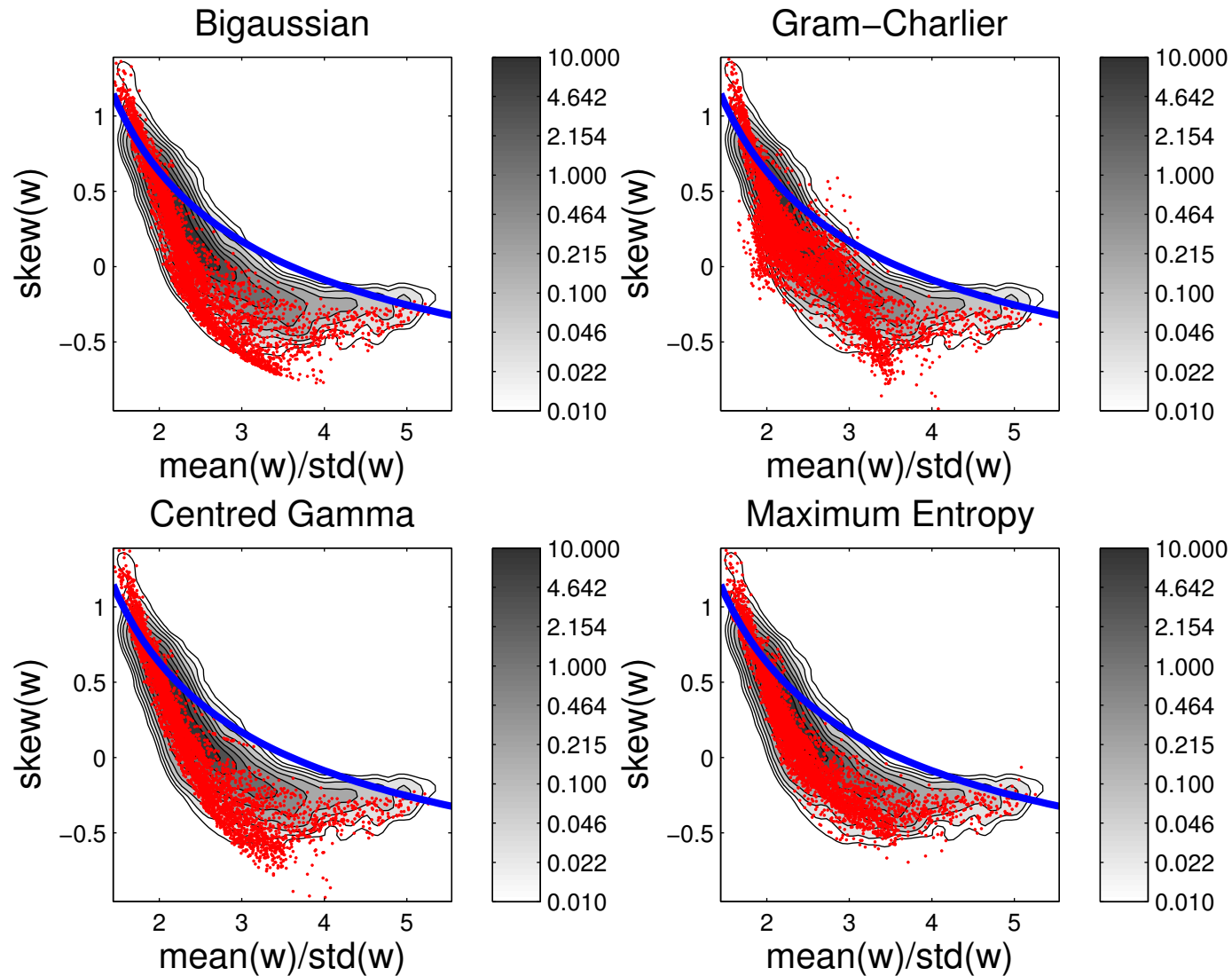


Maximum Entropy



Empirical Wind Speed pdfs: Intercomparison

■ Relationships between moments



Wind Speed pdfs

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- Relationship between moments can be found **empirically** or **mechanistically**

Mechanistic Model

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- Horizontal momentum equation:

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\frac{1}{\rho} \nabla p - f \hat{\mathbf{k}} \times \mathbf{u} - \frac{1}{\rho} \frac{\partial (\overline{\rho \mathbf{u}' u'_3})}{\partial z}$$

Mechanistic Model: Assumptions

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$$\frac{d\mathbf{u}}{dt} = -\frac{1}{\rho}\nabla p - f\hat{\mathbf{k}} \times \mathbf{u} - \frac{c_d}{h}w\mathbf{u} + \frac{K}{h^2}(\mathbf{U} - \mathbf{u})$$



Mechanistic Model: SDE

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- Finally, we obtain stochastic differential equation

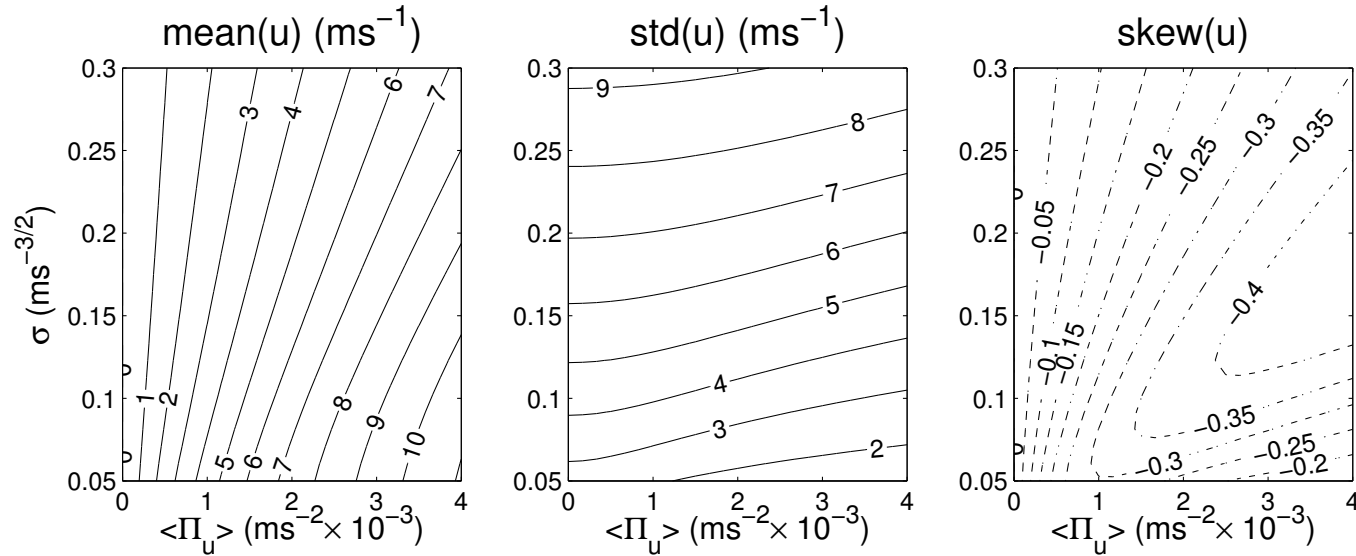
$$\begin{aligned}\frac{du}{dt} &= \langle \Pi_u \rangle - \frac{c_d}{h}wu - \frac{K}{h^2}u + \sigma \dot{W}_1 \\ \frac{dv}{dt} &= -\frac{c_d}{h}wv - \frac{K}{h^2}v + \sigma \dot{W}_2\end{aligned}$$

Mechanistic Model: pdf

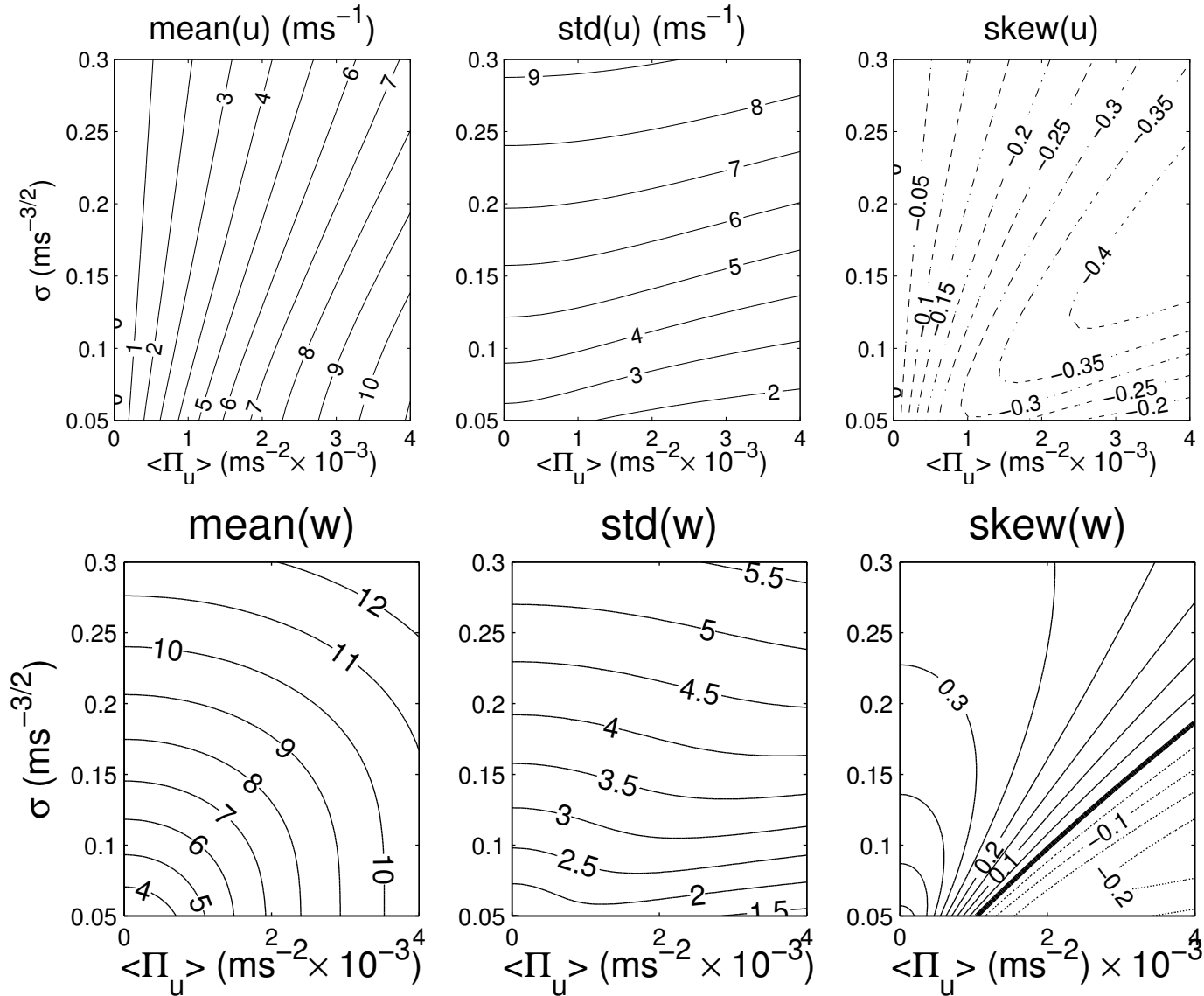
- The stochastic differential equation has an associated Fokker-Planck equation for the stationary pdf, which yields the analytic solution:

$$p_{uv}(u, v) = \mathcal{N}_1 \exp \left(\frac{2}{\sigma^2} \left\{ \langle \Pi_u \rangle u - \frac{K}{2h^2} (u^2 + v^2) - \frac{1}{h} \int_0^{\sqrt{u^2 + v^2}} c_d(w') w'^2 dw' \right\} \right)$$

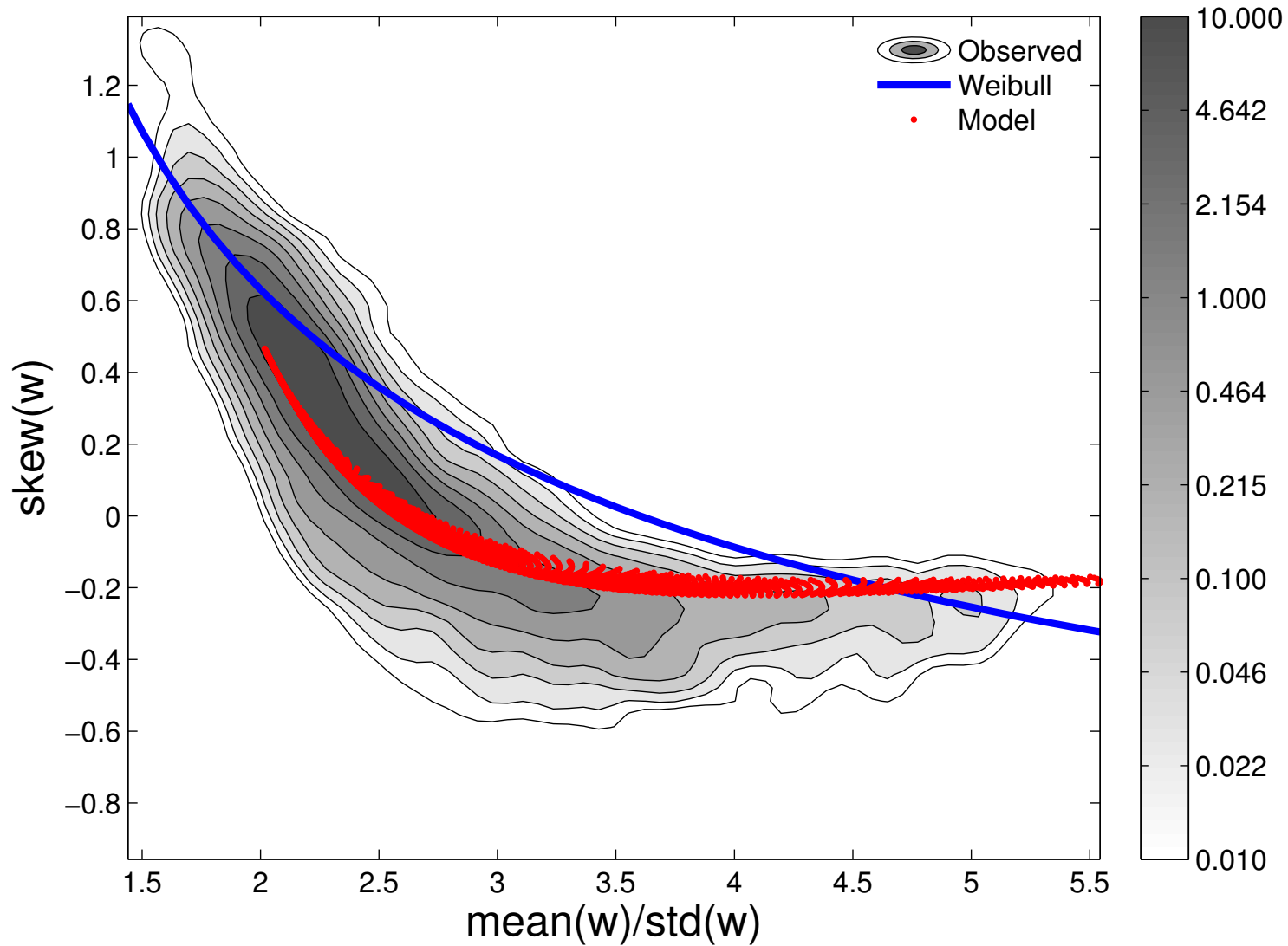
Mechanistic Model: Predictions



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Mechanistic Model: Comparison with Observations



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- Qualitative success of model suggests it has captured essential physics: still improvements to be made

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