

# Thermodynamics of an ideal gas

eqn. of state for ideal gas:  $p = \rho RT$   
 $T$  in  $^{\circ}K$

comes from experiments showing

$$pdV + Vdp = RdT \quad \text{where } V \text{ is}$$

the volume of a unit mass,  $R$  is a constant,  $dT$  can be in  $^{\circ}C$  or  $^{\circ}K$  with  
 $^{\circ}K = ^{\circ}C + 273$

integrate to get  $pV = RT$  where  
 $T$  is in  $^{\circ}K$  so that the integration constant is zero

$$\Rightarrow p = \rho RT \quad \rho = \frac{1}{V}$$

$$R = C_p - C_v$$

$C_p$  = specific heat at constant pressure

$$C_p \equiv \left. \frac{\partial h}{\partial T} \right|_p \quad h = \hat{u} + \frac{p}{\rho}$$

(2)

$\hat{u}$  = internal energy per unit mass

$p$  = pressure = the (compressive) normal stress on a fluid at rest

$h$  = enthalpy per unit mass

$c_v$  = specific heat at constant volume

$$c_v \equiv \left. \frac{\partial \hat{u}}{\partial T} \right|_v$$

How can we understand  $c_p$  and  $c_v$ ?  
idealized situation

$c_p$  is the amount of energy required to raise the temperature of a unit mass by  $1^\circ\text{K}$ , keeping pressure fixed

$$c_p = \left[ \frac{\text{energy}}{\text{mass} \cdot ^\circ\text{K}} \right]$$

1st Law of Thermodynamics

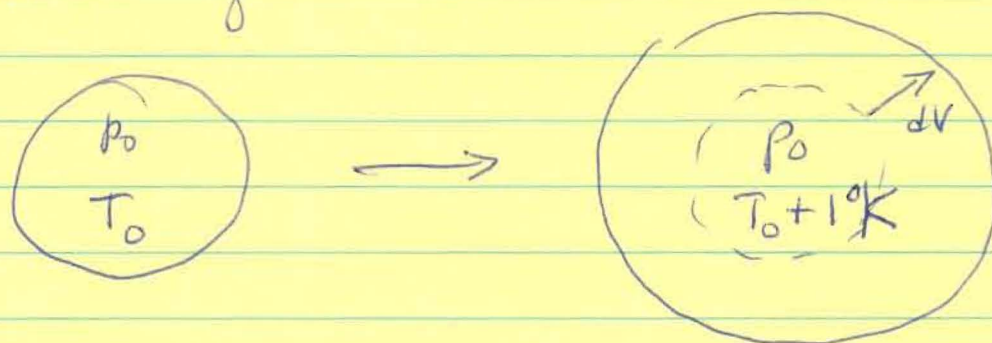
$$\delta Q = d\hat{u} + \delta W$$

where  $\delta Q$  is small amount of heat added to the system  
 $\delta W$  is " " of work done by the system

$d\hat{u}$  is the change in internal energy

[ $\delta Q$ ,  $\delta W$  are path-dependent]

balloon of gas



To heat the gas inside, keeping  $p_0$  fixed, some energy goes into expanding the balloon, doing work against the atmosphere

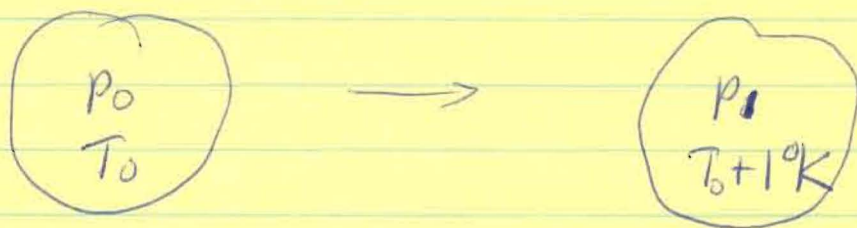
$$\delta Q = d\hat{u} + p dV$$

defining  $h = \hat{u} + pV = \hat{u} + P/\rho$

Then  $c_p = \left. \frac{\partial h}{\partial T} \right|_p$

(4)

Alternatively, keeping volume fixed



$$\delta Q = d\hat{u}$$

$$\text{then } C_V = \left. \frac{\partial \hat{u}}{\partial T} \right|_V$$

Note that heat is a form of energy

temperature is a measure of equilibrium

Heat depends on how much mass there is,  
temperature does not

Heat and work are not functions  
internal energy is a function

(5)

Ok, how do we get  $R = C_p - C_v$   
[obviously  $C_p > C_v$ ]

ideal gas law in differential form:

$$p dV + V dp = R dT$$

By definition  $h = \hat{u} + pV$

$$\Rightarrow dh = d\hat{u} + p dV + V dp$$

$$= d\hat{u} + R dT \quad \text{using ideal gas law above}$$

Also by definition

$$dh|_p = c_p dT$$

$$\Rightarrow c_p dT = d\hat{u} + R dT$$

Also by definition  $d\hat{u}|_v = c_v dT$

(6)

IF  $\hat{u} = \hat{u}(T)$  is only a function of temperature, then

$$C_p dT = C_v dT + R dT$$

$$\text{or } C_p = C_v + R$$


---

It remains to show that  $\hat{u} = \hat{u}(T)$   
for an ideal gas

To show  $\hat{u} = \hat{u}(T)$  use

ideal gas law  $V dp + p dV = R dT$

1<sup>st</sup> and 2<sup>nd</sup> Laws of Thermo

$$\delta Q = d\hat{u} + p dV = T dS$$

$S = \text{entropy}$  (a function)

## Manipulations

$$\delta G = d\hat{u} + R dT - V dp = T dS$$

$$\Rightarrow dS = \frac{d\hat{u}}{T} + R \frac{dT}{T} - \frac{V dp}{T}$$

$$= \frac{d\hat{u}}{T} + R d(\ln T) - \frac{R}{p} dp$$

$$= \frac{d\hat{u}}{T} + R d(\ln T) - R d(\ln p)$$

$$= \frac{d\hat{u}}{T} + R d[\ln T - \ln p]$$

$$= \frac{d\hat{u}}{T} + R d[\ln(T/p)]$$

$$= \frac{d\hat{u}}{T} + d[R \ln(T/p)] \quad R_{\text{constant}}$$

$$\Rightarrow \frac{d\hat{u}}{T} = d[S - R \ln(T/p)]$$

$$= dG \quad [\text{exact differential}]$$

(8)

$\frac{d\hat{u}}{T} = dG(p, T)$  bc. two thermodynamic quantities  $p, T$  determine  $S$

$$\frac{1}{T} \frac{\partial \hat{u}}{\partial p} dp + \frac{1}{T} \frac{\partial \hat{u}}{\partial T} dT = \frac{\partial G}{\partial p} dp + \frac{\partial G}{\partial T} dT$$

$$\Rightarrow \frac{1}{T} \frac{\partial \hat{u}}{\partial p} = \frac{\partial G}{\partial p} \quad \frac{1}{T} \frac{\partial \hat{u}}{\partial T} = \frac{\partial G}{\partial T}$$

Now since  $\frac{\partial^2 G}{\partial T \partial p} = \frac{\partial^2 G}{\partial p \partial T}$

$$\Rightarrow \frac{\partial}{\partial T} \left[ \frac{1}{T} \frac{\partial \hat{u}}{\partial p} \right] = \frac{\partial}{\partial p} \left[ \frac{1}{T} \frac{\partial \hat{u}}{\partial T} \right]$$

$$-\frac{1}{T^2} \frac{\partial \hat{u}}{\partial p} + \frac{1}{T} \frac{\partial^2 \hat{u}}{\partial T \partial p} = \frac{1}{T} \frac{\partial^2 \hat{u}}{\partial p \partial T} \quad \text{cancel out}$$

$$\Rightarrow \frac{\partial \hat{u}}{\partial p} = 0 \quad \text{Q.E.D.}$$

$$\Rightarrow \hat{u} = \hat{u}(T) \quad \text{only}$$