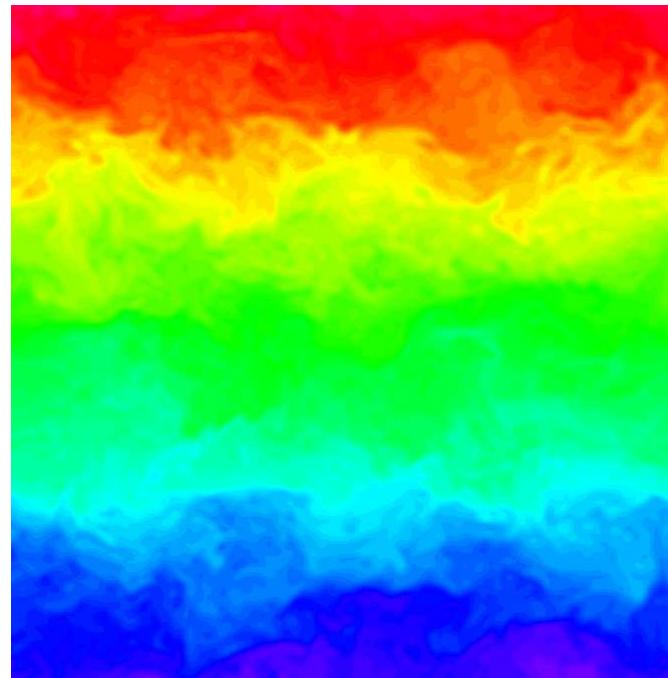


From Quasigeostrophic Turbulence to Stratified Turbulence

Peter Bartello¹ &
Michael L. Waite²

1. McGill University
2. Univ. of Victoria



Mesoscale atmospheric spectrum

- Gage (1979) then Lilly (1983) argued that there was an inverse cascade from convective scales. Rotation unimportant.
- Van Zandt (1982) said that it was a “gravity-wave spectrum”.

An early test of the Gage-Lilly idea

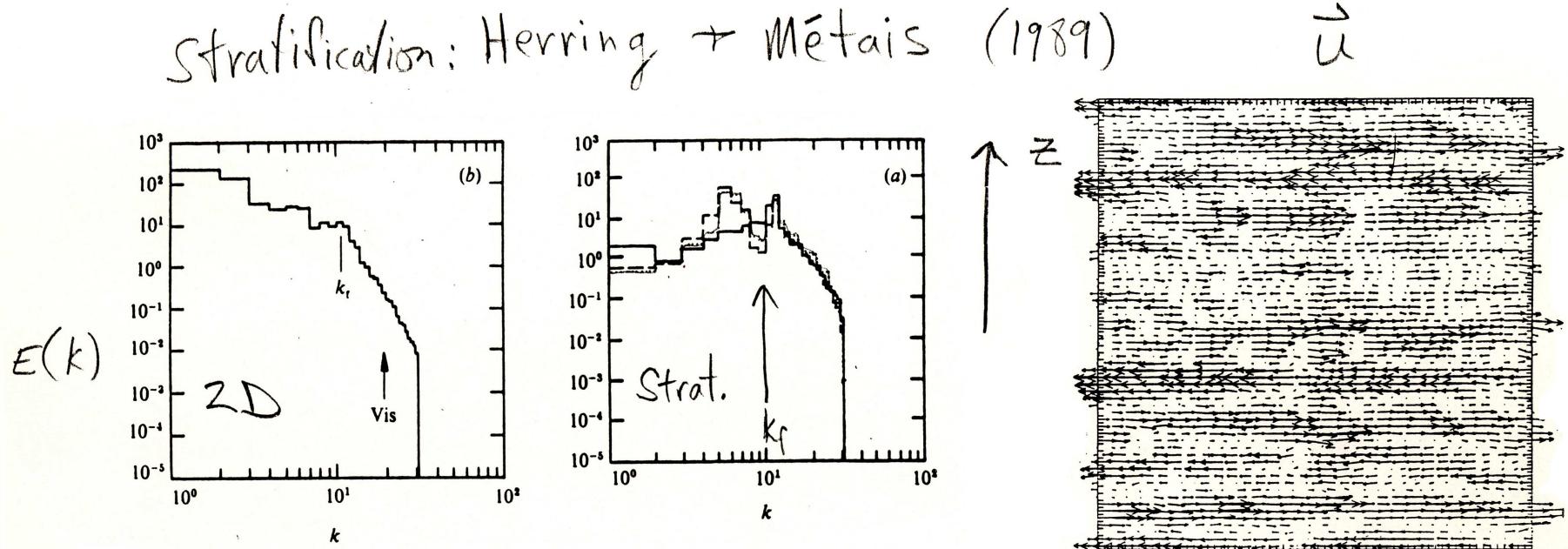


FIGURE 11. Vector plots of $u(x, y, z, t)$, $N = 80\pi$ for (x, z) -slice at the midplanes of the flow
Flow is statistically stationary.

- No inverse cascade with strong stratification and no rotation.
- Flow reached statistical stationarity (except for shear modes). Therefore energy went downscale.
- Later, Métais *et al.* added strong rotation and obtained an inverse cascade.

Boussinesq Normal Mode Decomposition

Let $G_{\mathbf{k}} = B_{\mathbf{k}}^{(0)}$ and $A_{\mathbf{k}}^{(\pm)}(\epsilon t) e^{\pm i\omega_{\mathbf{k}} t} = B_{\mathbf{k}}^{(\pm)}$, where $\epsilon = \min(Ro, Fr)$.

$$\frac{\partial}{\partial(\epsilon t)} \langle |G_{\mathbf{k}}|^2 \rangle = \sum_{\Delta} N_1 \langle G_{\mathbf{k}} G_{\mathbf{p}} G_{\mathbf{q}} \rangle +$$

$$N_2 \langle G_{\mathbf{k}} G_{\mathbf{p}} A_{\mathbf{q}} \rangle e^{i\omega_{\mathbf{q}} t} + N_3 \langle G_{\mathbf{k}} A_{\mathbf{p}} A_{\mathbf{q}} \rangle e^{i(\omega_{\mathbf{q}} + \omega_{\mathbf{p}})t}$$

1. conservation

$$\frac{\partial}{\partial(\epsilon t)} \langle |A_{\mathbf{k}}|^2 \rangle = \sum_{\Delta} M_1 \langle A_{\mathbf{k}} G_{\mathbf{p}} G_{\mathbf{q}} \rangle e^{i\omega_{\mathbf{k}} t} + M_2 \langle A_{\mathbf{k}} A_{\mathbf{p}} G_{\mathbf{q}} \rangle e^{i(\omega_{\mathbf{k}} + \omega_{\mathbf{p}})t}$$

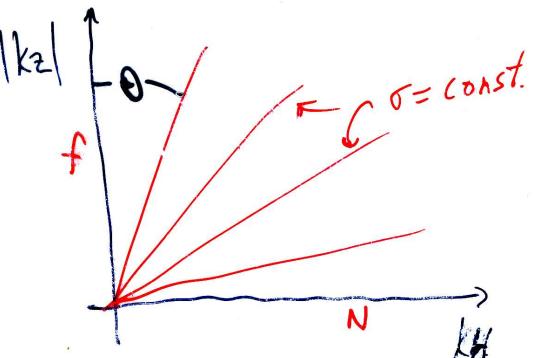
2. resonance

$$+ M_3 \langle A_{\mathbf{k}} A_{\mathbf{p}} A_{\mathbf{q}} \rangle e^{i(\omega_{\mathbf{k}} + \omega_{\mathbf{p}} + \omega_{\mathbf{q}})t}$$

Energy, $E = \sum_{\mathbf{k}} |G_{\mathbf{k}}|^2 + |A_{\mathbf{k}}|^2 = E_G + E_A$.

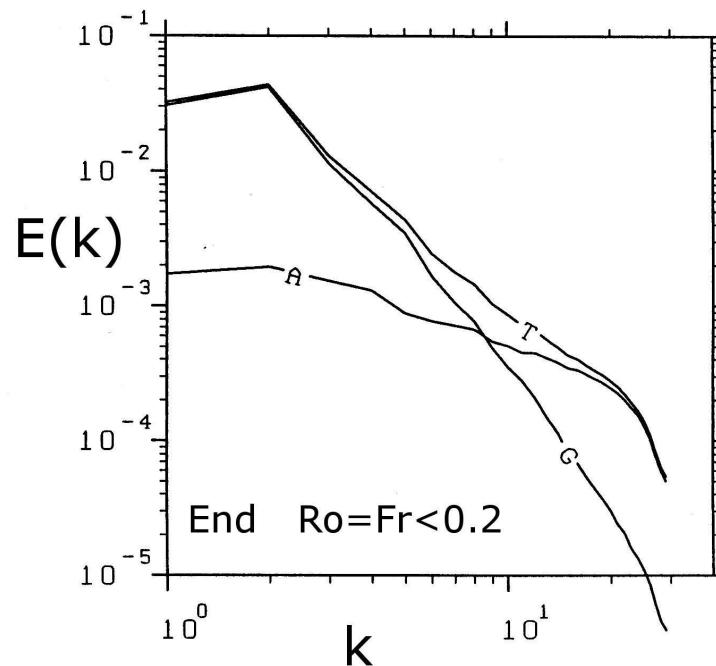
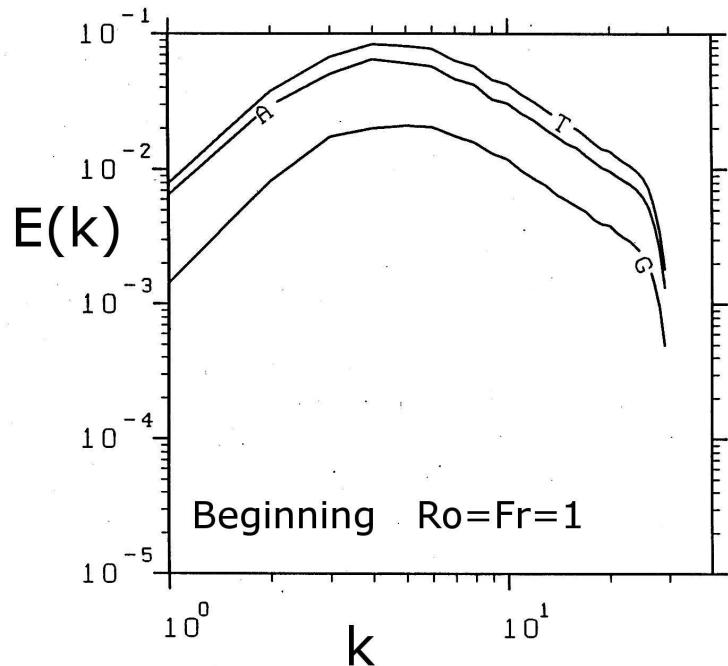
Potential Enstrophy, $V \doteq \sum_{\mathbf{k}} \sigma_{\mathbf{k}}^2 k^2 |G_{\mathbf{k}}|^2$ as

Gravity-wave dispersion relation, $\sigma_{\mathbf{k}} = (f^2 c$



Strong rotation and strong stratification

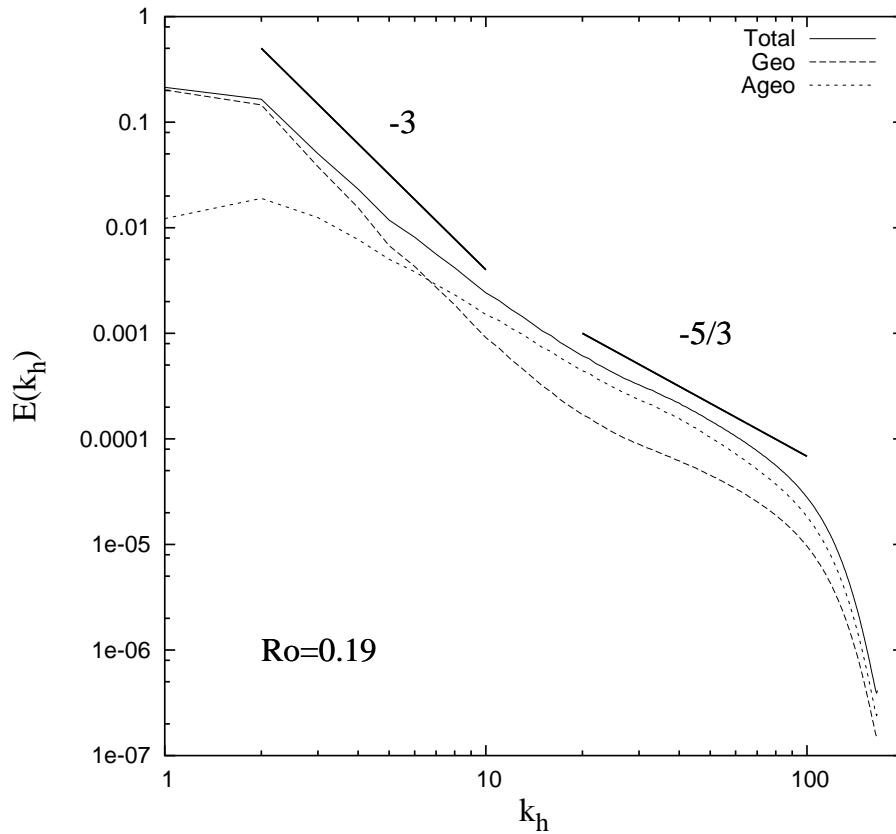
Geostrophic, G and ageostrophic, A spectra



- G modes cascade upscale, A modes downscale \rightarrow balance
- Balance requires dissipation of balanced modes to be less than that of unbalanced ones (e.g. inverse cascades).

The atmosphere?

- Boussinesq with $N/f = 10$ + Rayleigh damping
- Grid: $500 \times 500 \times 50$ with $\Delta x = \Delta y = \Delta z$
- A set of baroclinic QG modes with $L/H = N/f$ are forced.



Stratification only: Dimensionless equations

$$Fr_h = \frac{U}{NL}, \quad Fr_z = \frac{U}{NH}.$$

Two scaling ideas:

- Integral scale vertical Froude number $\rightarrow O(1)$, i.e. $H_i \sim \sqrt{E}/N$
 - Vertical Froude is unity at all vertical scales, i.e.
 $H(k_z) \sim [k_z E(k_z)]^{1/2}/N$ or equivalently $E(k_z) \sim N^2 k_z^{-3}$

Limit of Strong Stratification ($\Omega = 0$)

- The vertical scale collapses until dissipation scale reached.
- Two limits:
 - 1. $Fr_h \rightarrow 0$, Fr_h/Fr_z fixed (RMW 1981)
 - Appropriate if flow is initially isotropic or when vertical scale determined by dissipation
 - Potential enstrophy approximately quadratic
 - 2. $Fr_h \rightarrow 0$, Fr_z fixed
 - Implies $H \sim U/N$ set by stratification. Due to
 - i) Thorpe (1977) (introduced characteristic scale for overturning. U/N is an upper bound.)
 - ii) Munk (1981) (identified it as small-scale end of Garrett-Munk spectrum)
 - iii) Hines (1996) (transition scale between unsaturated and saturated waves)
 - iv) Billant & Chomaz (2001)

Stratification only

- Vectors and ∇ , etc. are horizontal components
- Use vertical time scale $T \sim L/U$:

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} + Fr_z^2 w \frac{\partial \mathbf{u}}{\partial z} = -\nabla p,$$

$$Fr_h^2 \left(\frac{\partial w}{\partial t} + \mathbf{u} \cdot \nabla w + Fr_z^2 w \frac{\partial w}{\partial z} \right) = -\frac{\partial p}{\partial z} + b,$$

$$\nabla \cdot \mathbf{u} + Fr_z^2 \frac{\partial w}{\partial z} = 0,$$

$$\frac{\partial b}{\partial t} + \mathbf{u} \cdot \nabla b + Fr_z^2 w \frac{\partial b}{\partial z} = -w.$$

- Decoupled layers when $Fr_h \rightarrow 0$, $Fr_z \rightarrow 0$.
- 3D hydrostatic turbulence if $H \sim U/N$

Wave/vortical decomposition without rotation

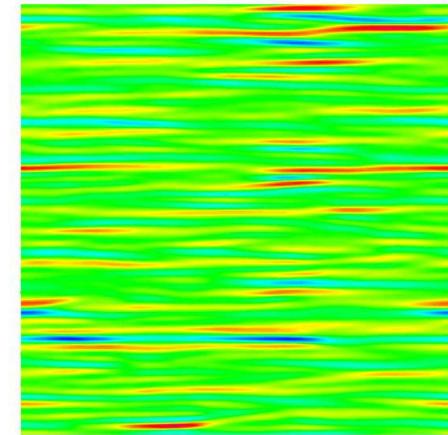
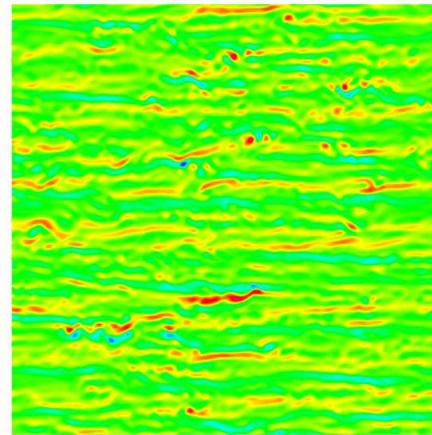
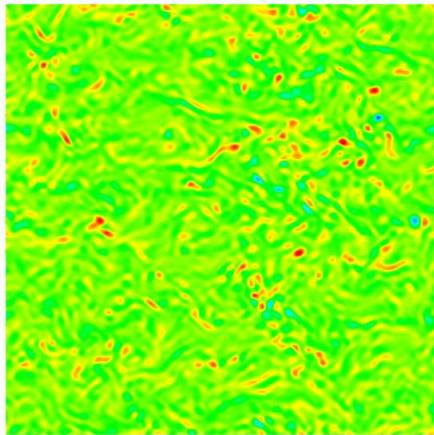
- Need to be careful: potential enstrophy only approximately quadratic when $Fr_z \ll 1$, but this may not be relevant.
- We need to restrict application to scales much larger than H_i or to particular problems, e.g. early decay of isotropic i.c.'s.
- If $Fr_z \ll 1$ we can examine inviscid truncated statistical equilibrium

$$\langle |G_{\mathbf{k}}|^2 \rangle = \frac{1}{\lambda_1 + \lambda_2 k_h^2} \quad \langle |A_{\mathbf{k}}|^2 \rangle = \frac{2}{\lambda_1}$$

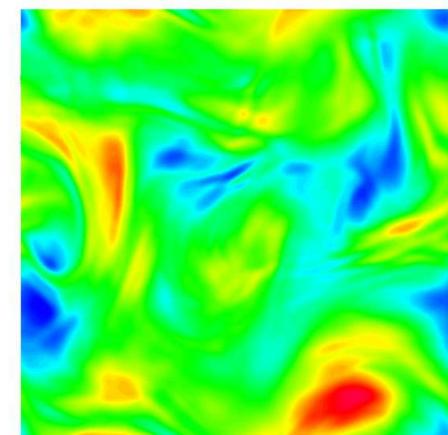
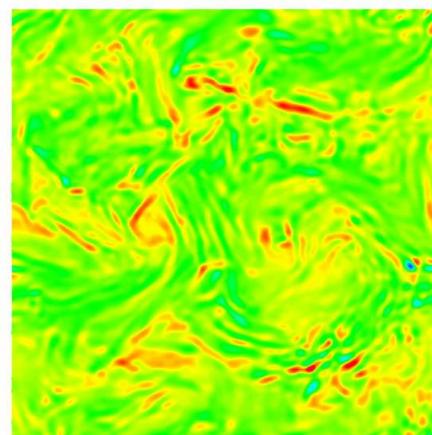
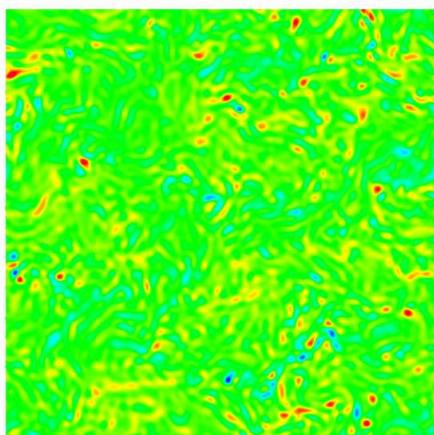
- Independent of k_z
- Can show that if $k_i \ll k_{max}$, where k_i defined by the ratio of potential enstrophy to energy, then as $k_{max} \rightarrow \infty$, the spectrum is peaked at a wavenumber $\rightarrow \sqrt{2}k_i$.
- In 2D turbulence, this wavenumber goes to zero. This argument suggest that there is no inverse cascade.

Vortical forcing ($\Omega = 0$) : perpendicular vorticity

$x - z$ plane

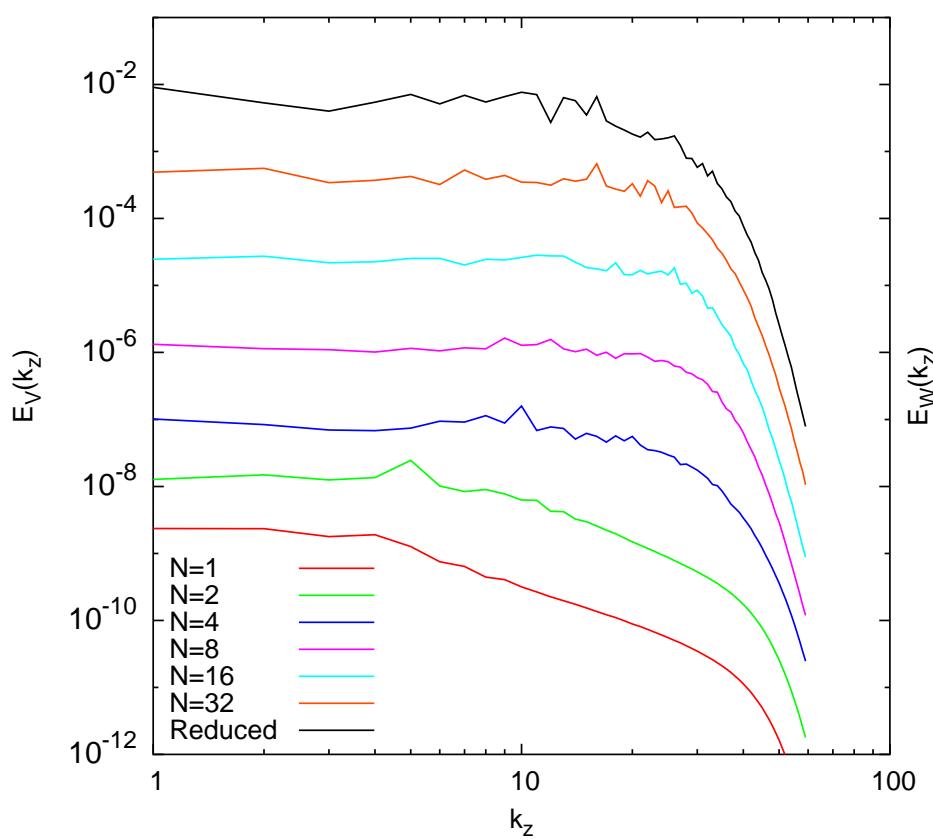


$x - y$ plane

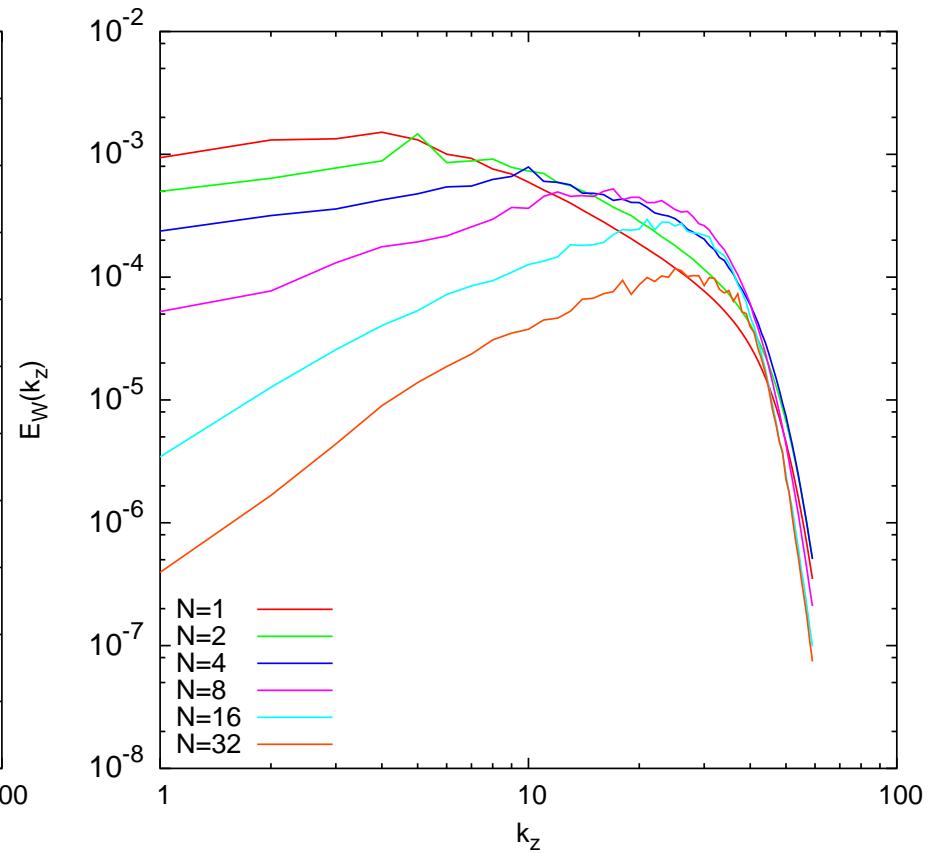


Vortical forcing ($\Omega = 0$): k_z spectra

Vortical energy



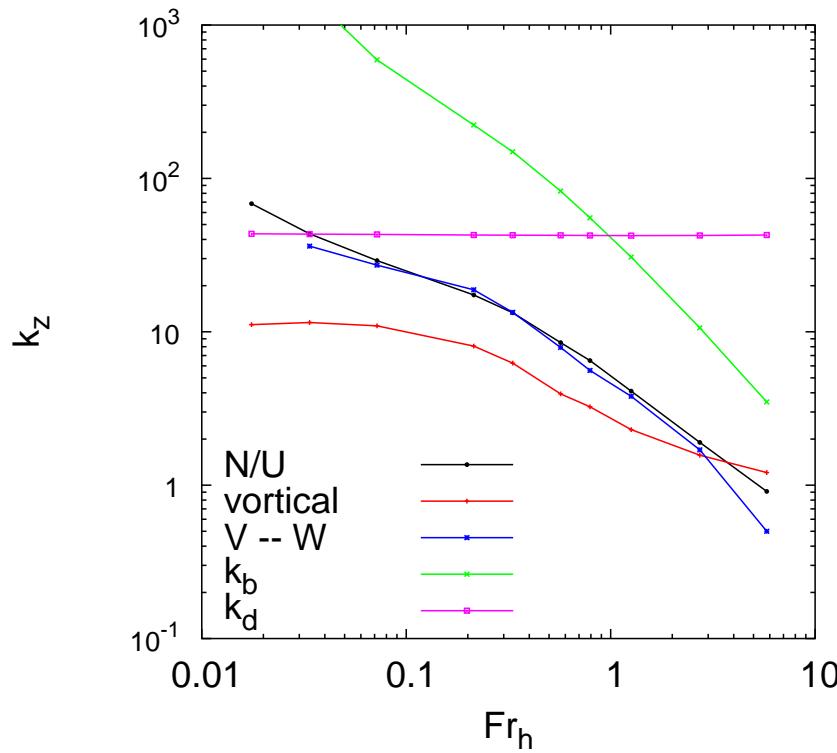
Wave energy



offset by factors of 10

Vortical Forcing ($\Omega = 0$): Length scales

Vortical forcing:



- $k_b \sim \sqrt{N^3/\epsilon}$ Ozmidov
- k_d is dissipation wavenumber

Wave forcing ($\Omega = 0$): $N^2 k_z^{-3}$ spectra?

Dashed line: $0.1N^2 k_z^{-3}$ (Bouruet-Aubertot, Sommeria & Staquet 1996)

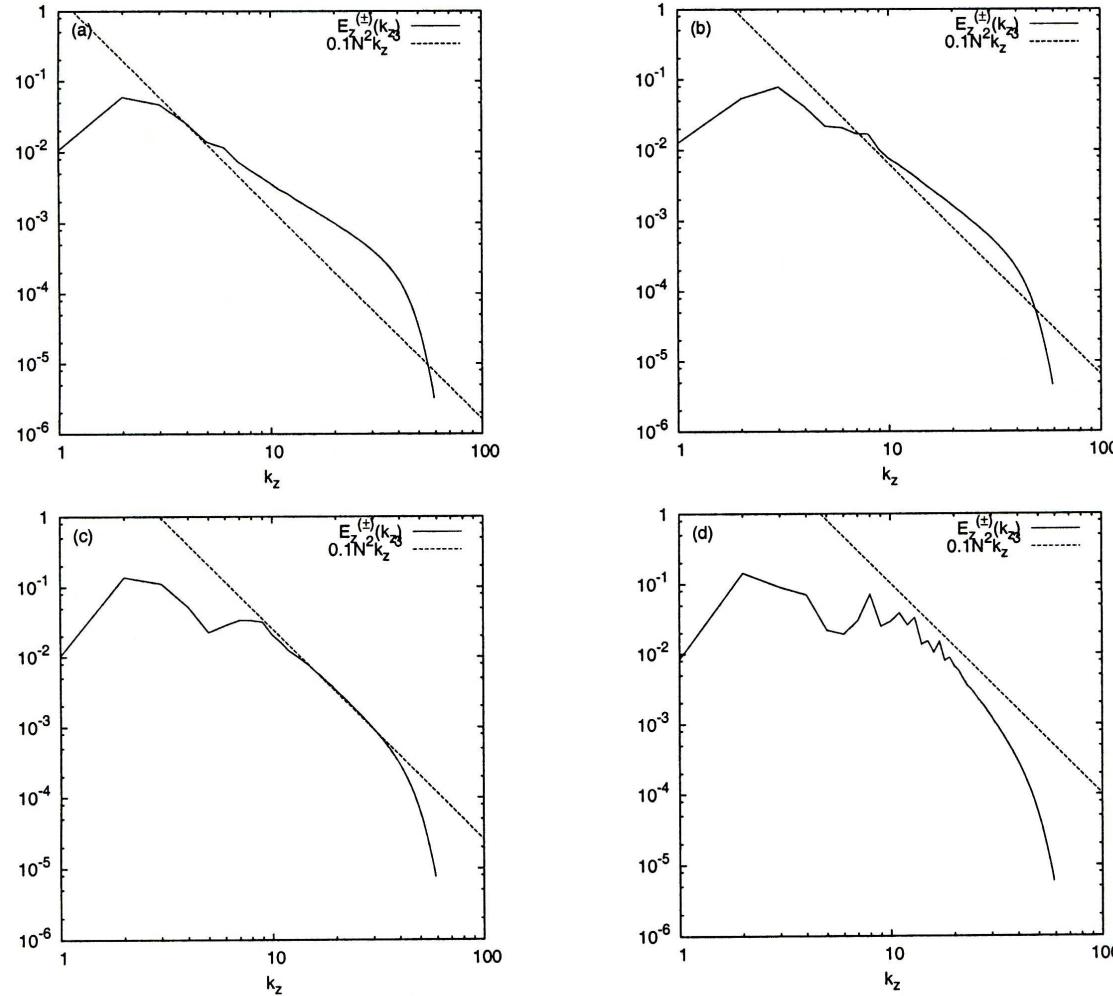
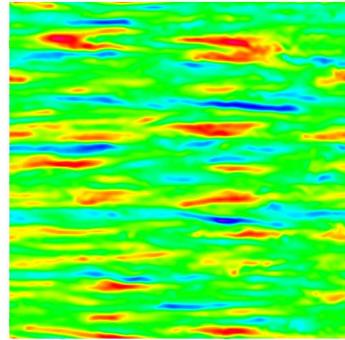
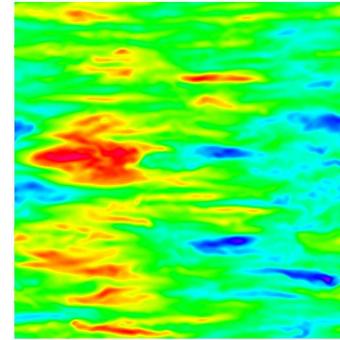
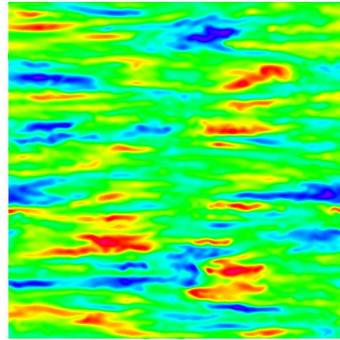


FIGURE 9. Vertical wavenumber spectra of wave energy, along with the hypothetical saturation spectrum $0.1N^2 k_z^{-3}$, for (a) $N = 4$, (b) $N = 8$, (c) $N = 16$ and (d) $N = 32$ when $M = 180$ and $R = 0$.

Strong stratification at various rotations



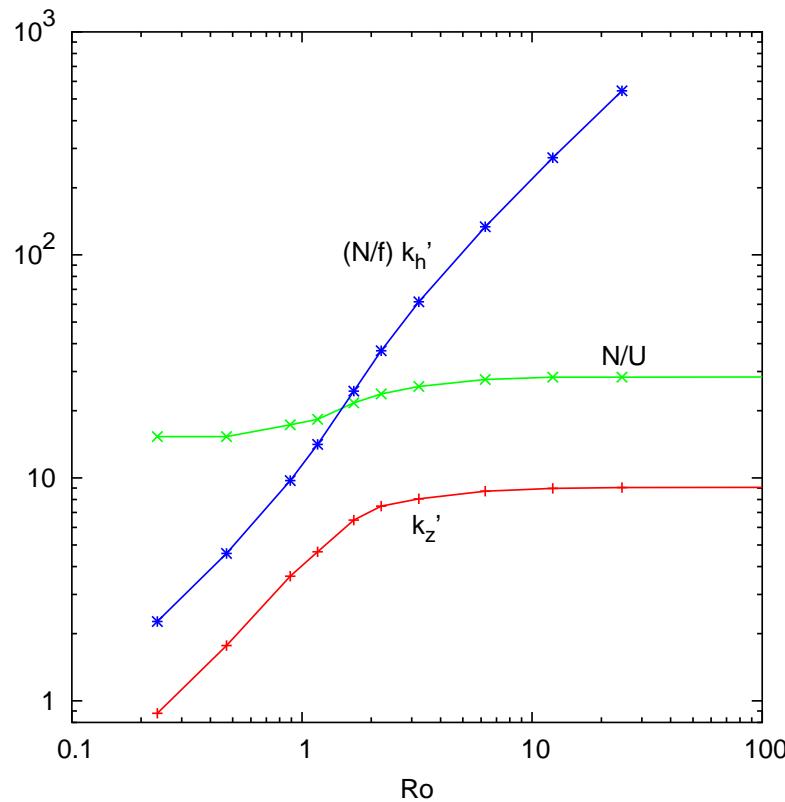
No rotation



lots of rotation

This is the horizontal velocity component perpendicular to the screen.

Strong stratification at various rotations

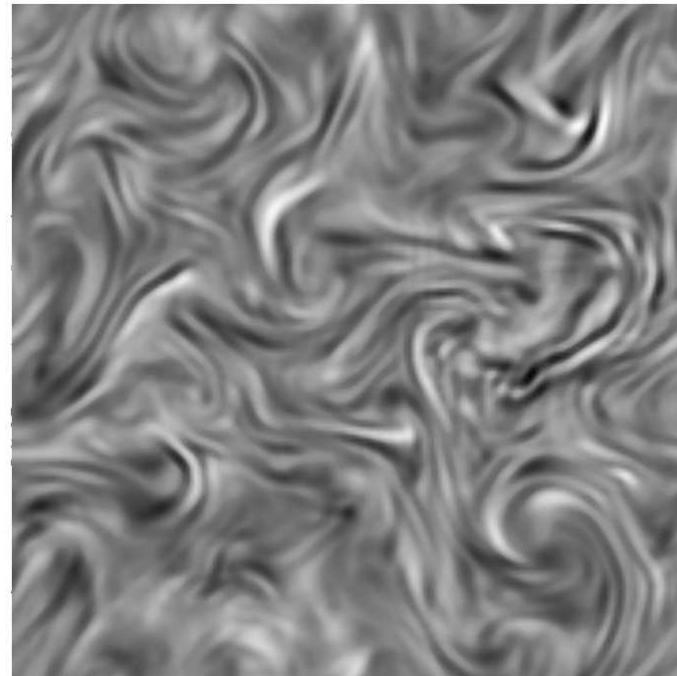
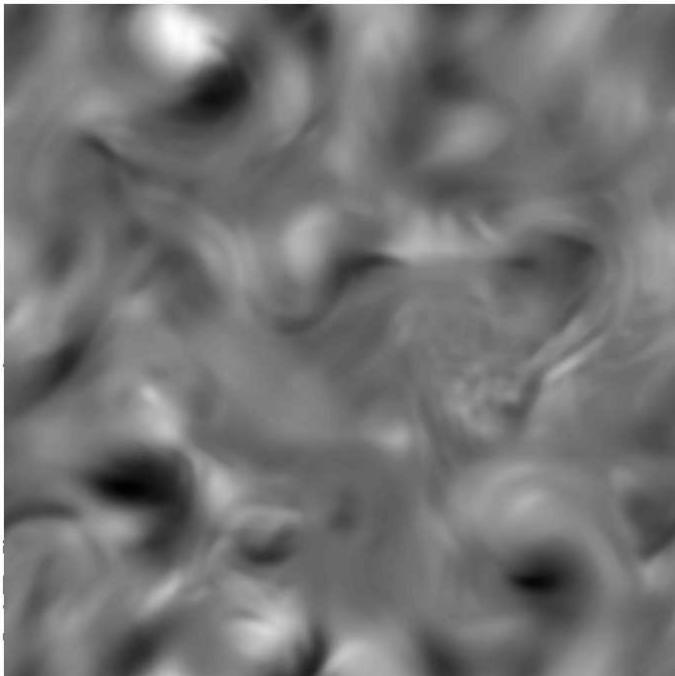


Stratified turbulence. Vertical lengthscale = U/N

Quasigeostrophic turbulence. Vertical lengthscale = fL/N
(Charney 1971)

This is how atmosphere/ocean modellers will need to set the ratio $\Delta z / \Delta x$ as resolutions increase.

Back to strong stratification and rotation



Balance at $\text{Ro}=0.09$?

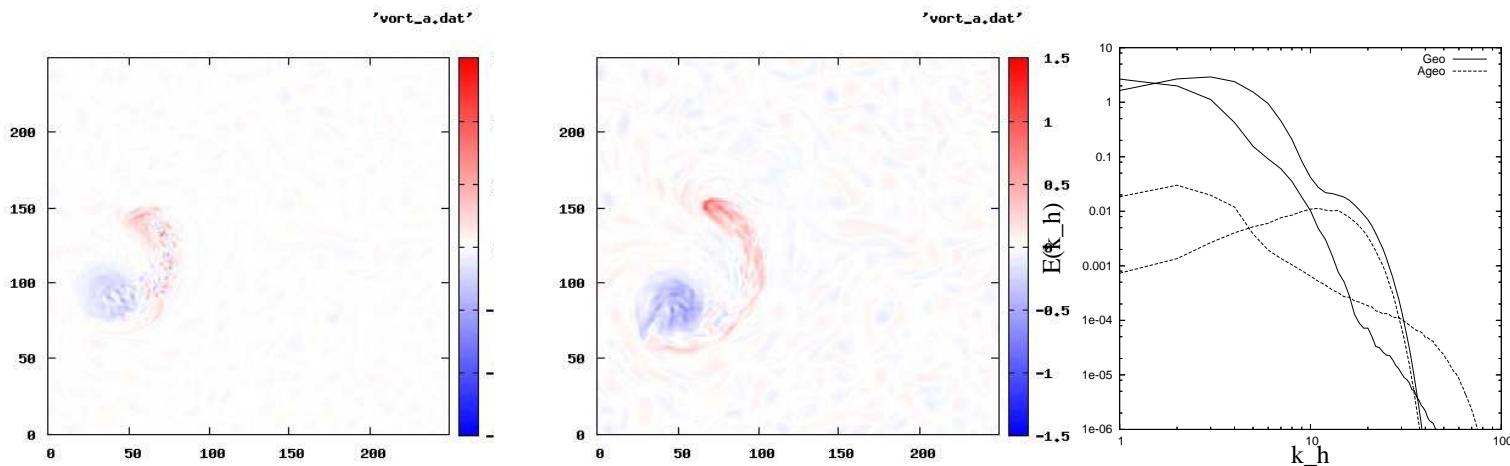
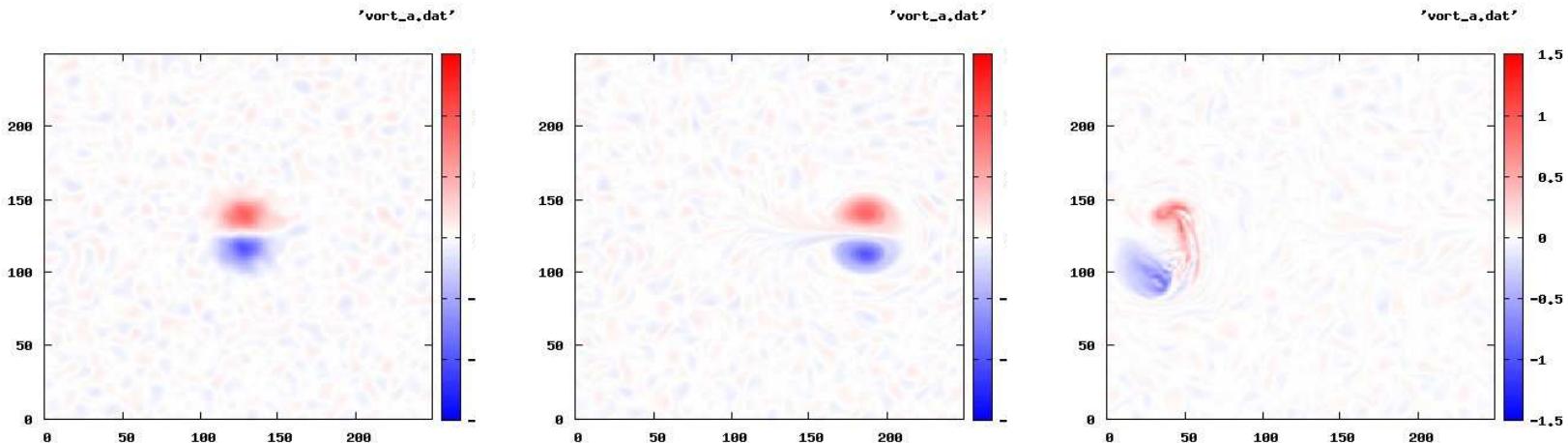
We have used the QG modes and the ω equation to diagnose vertical velocity (left).

The real vertical velocity is on the right.

If balance exists, it isn't this simple.

Regarding recent dipoles... (see Snyder *et al.*)

- Boussinesq with $N/f = 10$, $\text{Ro}=0.7$. Grid: $256 \times 256 \times 16$ with $\Delta x = \Delta y = \Delta z$
- Initial flow is a 2D dipole + a little 3D noise.



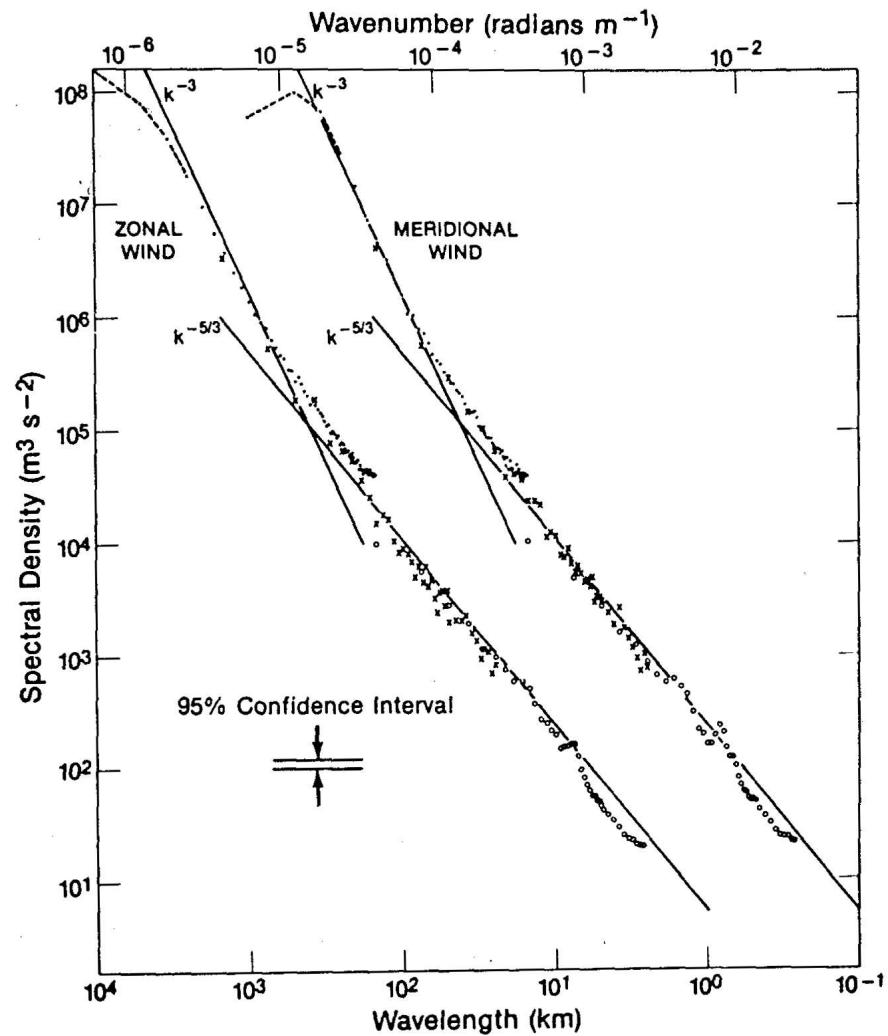
Summary of the New Stuff

- Stratified Turbulence without Rotation
 - Vortical forcing:
 - Direct cascade.
 - Simulations confirm $H \sim U/N$.
 - Wave forcing:
 - $N^2 k_z^{-3}$ not observed.
 - Both: When H not resolved by Δz weird things happen.
- Stratified Turbulence with rotation
 - Growth in the vertical integral scale as Ω increases
 - Bands of ageostrophic energy remain localised near vortices.
 - Steep vortical spectra, shallow ageo spectra *very robust*

Extra stuff

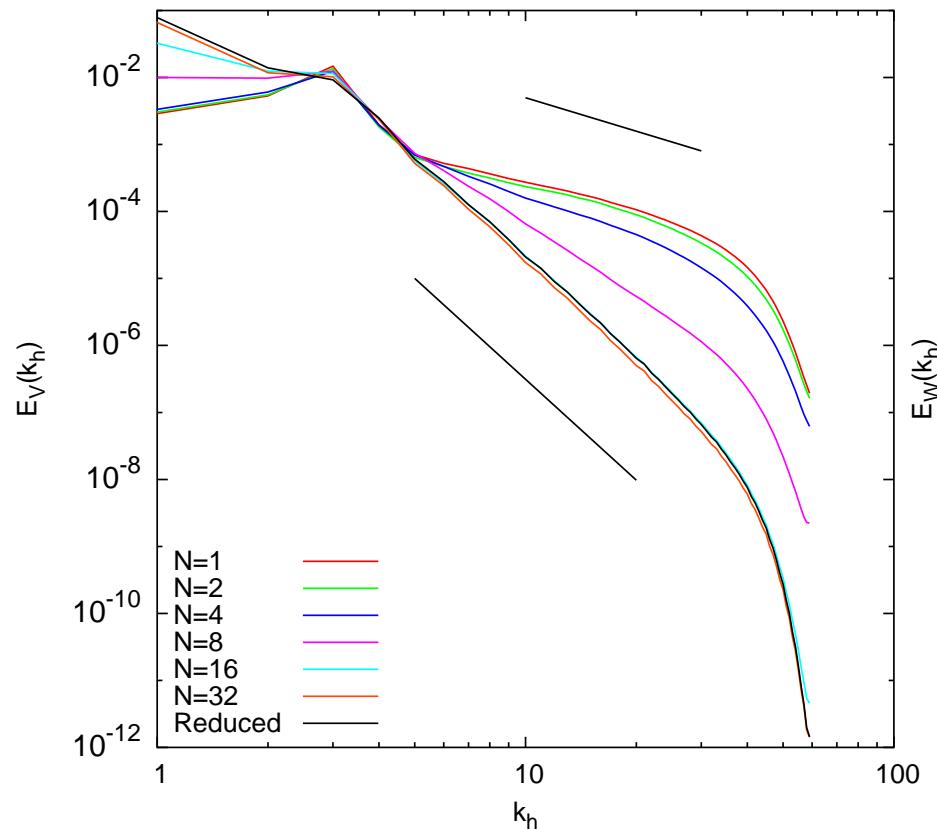
Observations: The Atmosphere

- GARP 1979
- Instruments on 7900 commercial flights
- Nastrom & Gage (1986)

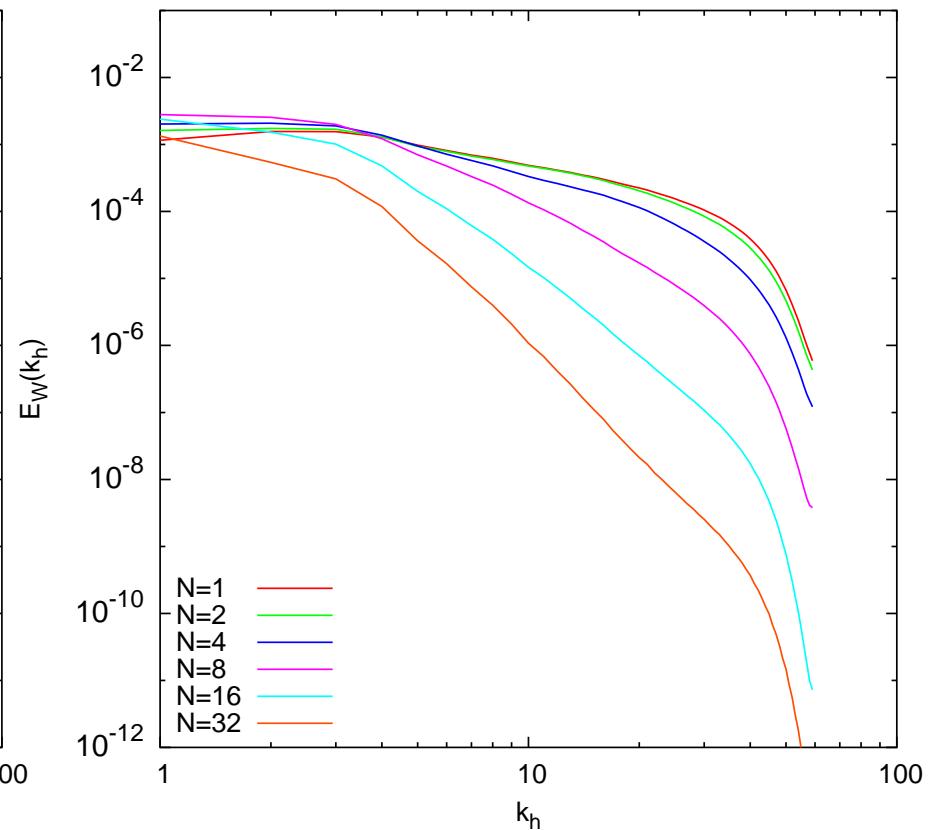


Vortical forcing ($\Omega = 0$): k_h spectra

Vortical energy

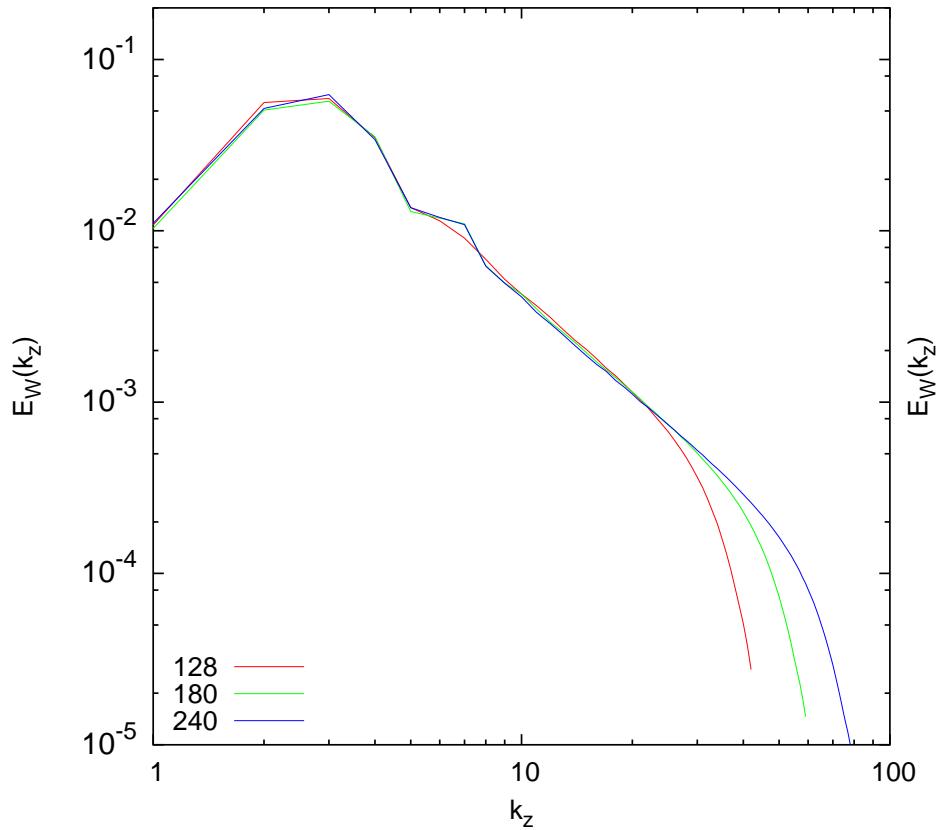


Wave energy

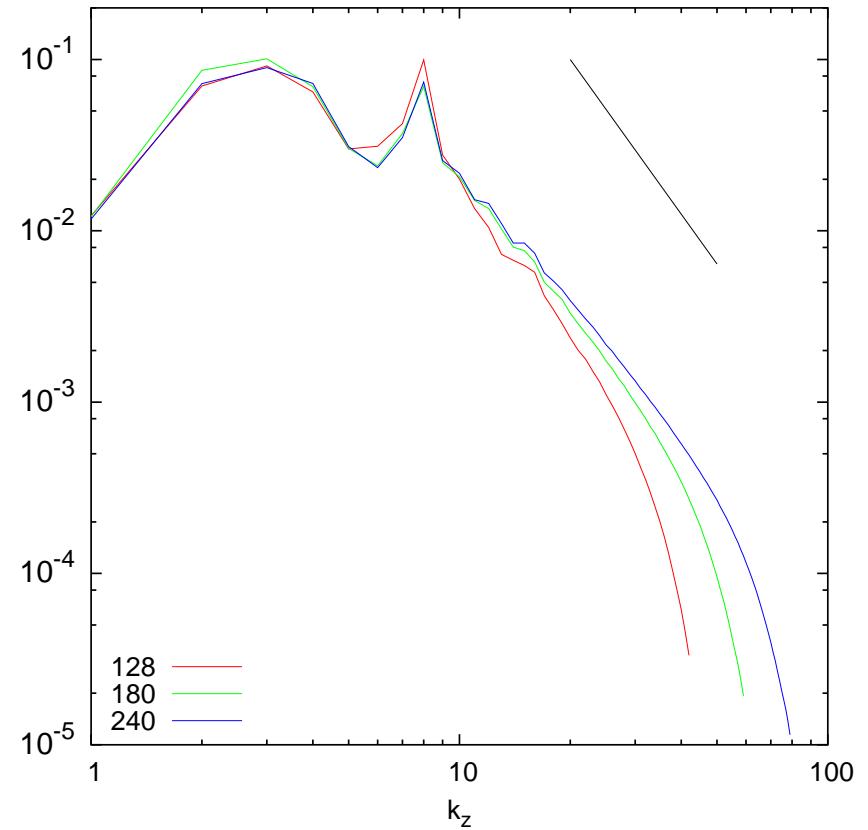


Wave forcing ($\Omega = 0$): dependence on Re

$Fr_h \sim 1$



$Fr_h \sim 0.1$



- Decreasing $Fr_h \rightarrow$ steeper spectra \rightarrow nonlocal interactions.

Bottleneck at Δ^4

