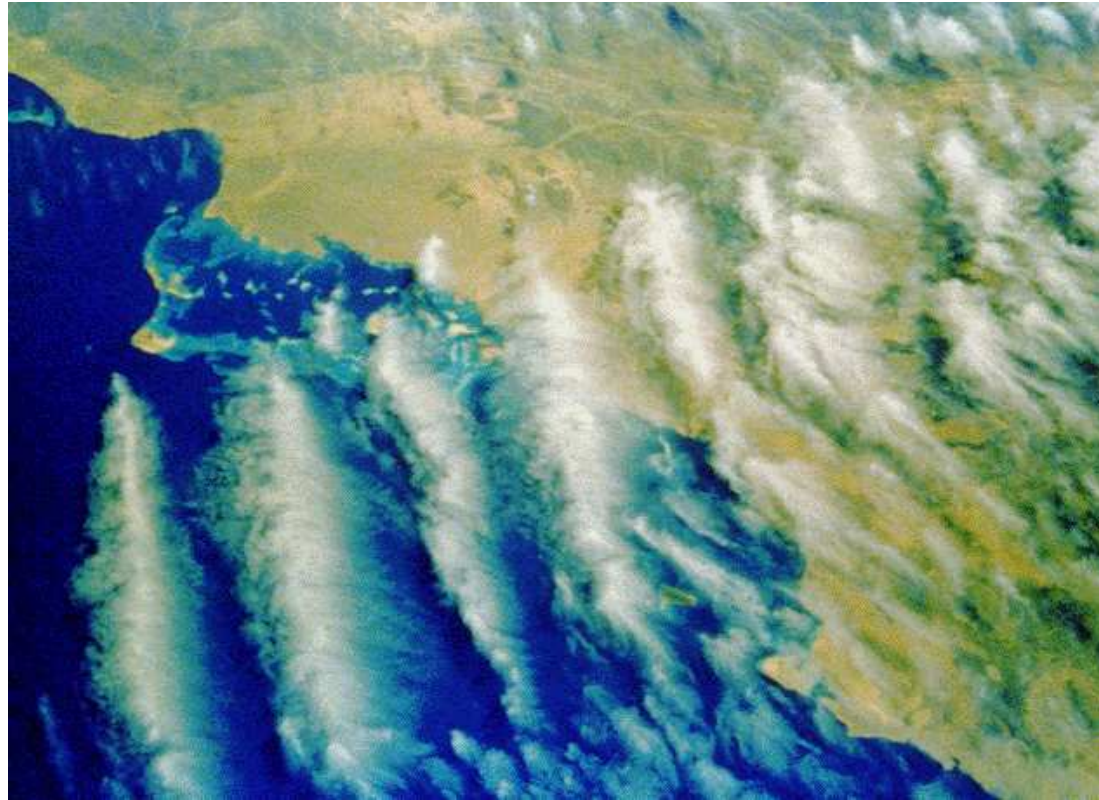


# Vortical-Wave Mode Interactions in Atmosphere-Ocean Flows

L.M. Smith, Y. Lee, M. Remmel, J. Sukhatme, F. Waleffe  
University of Wisconsin, Madison



Roll clouds in the jet stream over Saudi Arabia/Red Sea  
Supported by NSF and DOE

# Vortical-Wave Mode Interactions in Atmosphere-Ocean Flows

- What is a vortical mode, what is a wave mode, what are vortical-wave mode interactions?
- When are vortical-wave mode interactions important?
- Can we use this framework to understand balanced and unbalanced components of geophysical flows?

# Overview:

Rotating stratified fluid flow is a starting point to understand atmosphere-ocean phenomena.

In the linear limit, the governing equations possess:

- a so-called ‘vortical’ linear eigenmode
- wave eigenmodes

The simplest ‘reduced models’ (e.g. QG models)

- keep nonlinear interactions between vortical modes
- neglect wave mode interactions.

This strategy works to describe large-scale flows on short time scales, e.g. short-term weather prediction.

## Overview continued:

However, wave mode interactions contribute in many physical situations, e.g. flow over topography, and influence large-scale coherent flows on long times e.g. they can *generate* jets, vortices and layers.

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We introduce a framework to construct an understanding of all wave-vortical mode interactions.

New, non-perturbative models are derived 'intermediate' between QG and the full governing equations.

## I. Analytical Properties of the Governing Equations

- Solution as a superposition of linear eigenmodes
- Reduced models

## II. Example Numerical Results Depending Crucially on Wave Mode Interactions

- large-scale forcing, small-scale forcing, decay

## III. Derivation of new PDE reduced models to understand wave-vortical mode interactions

## IV. Numerical results for new PDE models

# The rotating Boussinesq equations

Conservation laws for vertically stratified flow rotating about the vertical  $\hat{\mathbf{z}}$ -axis:

$$\text{momentum :} \quad \frac{D\mathbf{u}}{Dt} + f\hat{\mathbf{z}} \times \mathbf{u} = -\nabla\phi - N\theta\hat{\mathbf{z}} + \nu\nabla^2\mathbf{u}$$

$$\text{mass :} \quad \nabla \cdot \mathbf{u} = 0$$

$$\text{energy :} \quad \frac{D\theta}{Dt} - Nw = \kappa\nabla^2\theta, \quad \theta = \frac{g}{N\rho_o}\rho'$$

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$$f = 2\Omega, \quad Ro = \frac{U}{fL}$$

$$\rho = \rho_o - bz + \rho', \quad \rho' \ll \rho_o, |bz|, \quad N^2 = \frac{gb}{\rho_o}, \quad Fr = \frac{U}{NH}$$

# Rossby and Froude numbers in geophysical flows

Pedlosky (1986) estimates:

- $Ro \approx 0.14$  for typical synoptic-scale winds at mid-latitudes

$$U \approx 10 \text{ m s}^{-1}, L \approx 1000 \text{ km}$$

- $Ro \approx 0.07$  in the western Atlantic

$$U \approx 5 \text{ cm s}^{-1}, L \approx 100 \text{ km}$$

Typical values are  $N/f \approx 100$  in the stratosphere  
and  $N/f \approx 10$  in the oceans.

Flows with  $N/f \approx L/H \implies Fr \approx 0.1$  (Burger number unity).

# Solutions in the unforced, linear, inviscid limit

$$[\mathbf{u}, \theta]^T(\mathbf{x}, t; \mathbf{k}) = \phi(\mathbf{k}) \exp \left[ i \left( \mathbf{k} \cdot \mathbf{x} - \sigma(\mathbf{k})t \right) \right] + \text{c.c.}$$

with eigenmodes  $\phi(\mathbf{k})$  and eigenvalues  $\sigma(\mathbf{k})$ .

- Wave modes  $\phi_+(\mathbf{k})$  and  $\phi_-(\mathbf{k})$  with

$$\sigma_{\pm}(\mathbf{k}) = \pm \frac{(N^2 k_h^2 + f^2 k_z^2)^{1/2}}{k}$$

- A non-wave (vortical or geostrophic) mode  $\phi_0(\mathbf{k})$  with

$$\sigma_0(\mathbf{k}) = 0$$



# Slow wave modes (as important as slow vortical modes!)

- Rotation-dominated flows

$$\sigma_{\pm}(\mathbf{k}) \approx \pm \frac{f k_z}{k}$$

slow when  $k_z = 0$ , e.g. vortical columns.

- Stratification-dominated flows

$$\sigma_{\pm}(\mathbf{k}) \approx \pm \frac{N k_h}{k}$$

slow when  $k_h = 0$ , e.g. horizontal shear layers (VSHF)

# Eigenmode representation for nonlinear flows

Since  $\phi_s(\mathbf{k})$ ,  $s = \pm, 0$  form an orthogonal basis

$$[\mathbf{u}, \theta]^T(\mathbf{x}, t) = \sum_{\mathbf{k}} \sum_s b_s(t; \mathbf{k}) \phi_s(\mathbf{k}) \exp \left[ i \left( \mathbf{k} \cdot \mathbf{x} - \sigma_s(\mathbf{k}) t \right) \right]$$

and the equations become

$$\frac{\partial}{\partial t} b_{s_{\mathbf{k}}} = \sum_{\Delta} \sum_{s_{\mathbf{p}}, s_{\mathbf{q}}} C_{\mathbf{k}\mathbf{p}\mathbf{q}}^{s_{\mathbf{k}} s_{\mathbf{p}} s_{\mathbf{q}}} b_{s_{\mathbf{p}}}^* b_{s_{\mathbf{q}}}^* \exp \left[ i \left( \sigma_{s_{\mathbf{k}}} + \sigma_{s_{\mathbf{p}}} + \sigma_{s_{\mathbf{q}}} \right) t \right]$$

27 interaction types, including 3-wave interactions

Exact and near resonances dominate:  $|\sigma_{s_{\mathbf{k}}} + \sigma_{s_{\mathbf{p}}} + \sigma_{s_{\mathbf{q}}}| \ll 1$ .

Reduced models resulting from restriction of the sum

$$\frac{\partial}{\partial t} b_{s_{\mathbf{k}}} = \sum_{\triangle} \sum_{s_{\mathbf{p}}, s_{\mathbf{q}}} C_{\mathbf{k}\mathbf{p}\mathbf{q}}^{s_{\mathbf{k}} s_{\mathbf{p}} s_{\mathbf{q}}} b_{s_{\mathbf{p}}}^* b_{s_{\mathbf{q}}}^* \exp \left[ i \left( \sigma_{s_{\mathbf{k}}} + \sigma_{s_{\mathbf{p}}} + \sigma_{s_{\mathbf{q}}} \right) t \right]$$

automatically conserve energy because each triad  $(\mathbf{k}, \mathbf{p}, \mathbf{q})$  satisfies the detailed balance:

$$C_{\mathbf{k}\mathbf{p}\mathbf{q}}^{s_{\mathbf{k}} s_{\mathbf{p}} s_{\mathbf{q}}} + C_{\mathbf{p}\mathbf{q}\mathbf{k}}^{s_{\mathbf{p}} s_{\mathbf{q}} s_{\mathbf{k}}} + C_{\mathbf{q}\mathbf{k}\mathbf{p}}^{s_{\mathbf{q}} s_{\mathbf{k}} s_{\mathbf{p}}} = 0$$

## 3D Pure Rotation

keeping only slow wave modes with  $k_z = 0 \implies$  symmetric 2D flow for  $(u,v)$ ;  $w$  is a passive scalar.

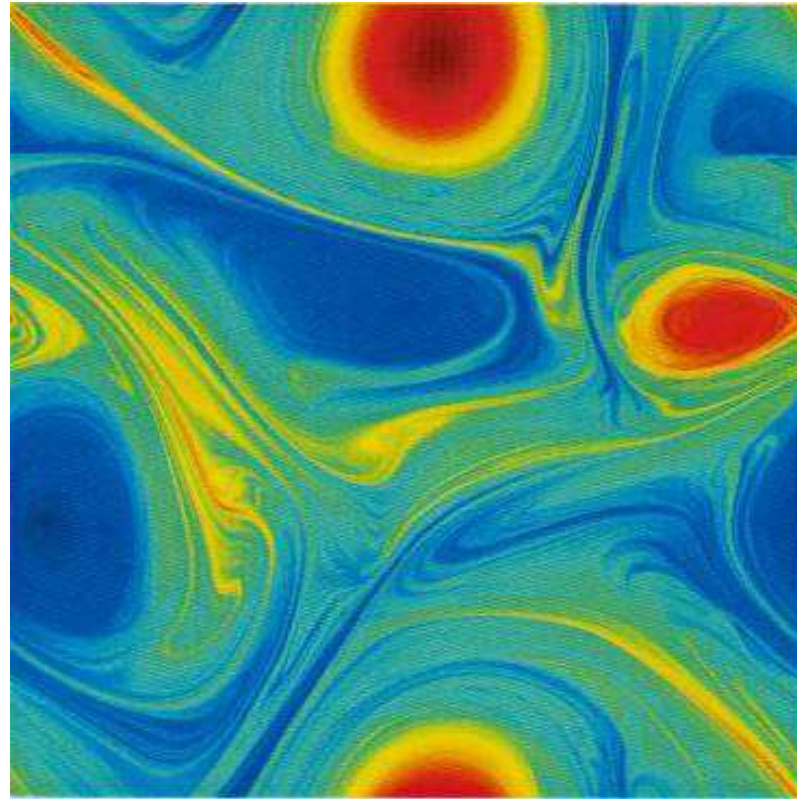
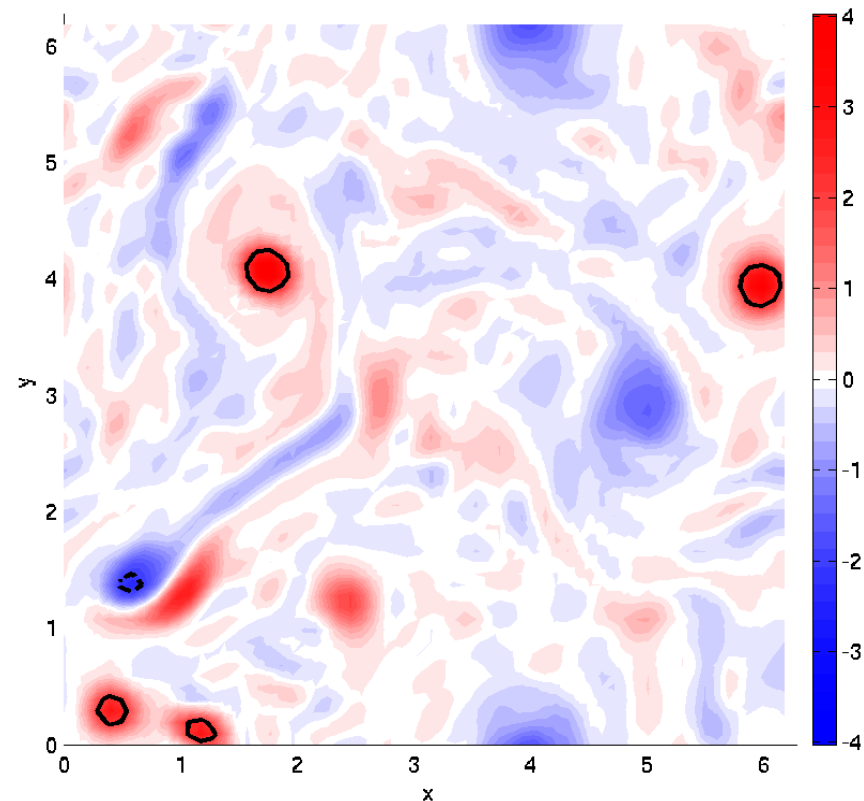
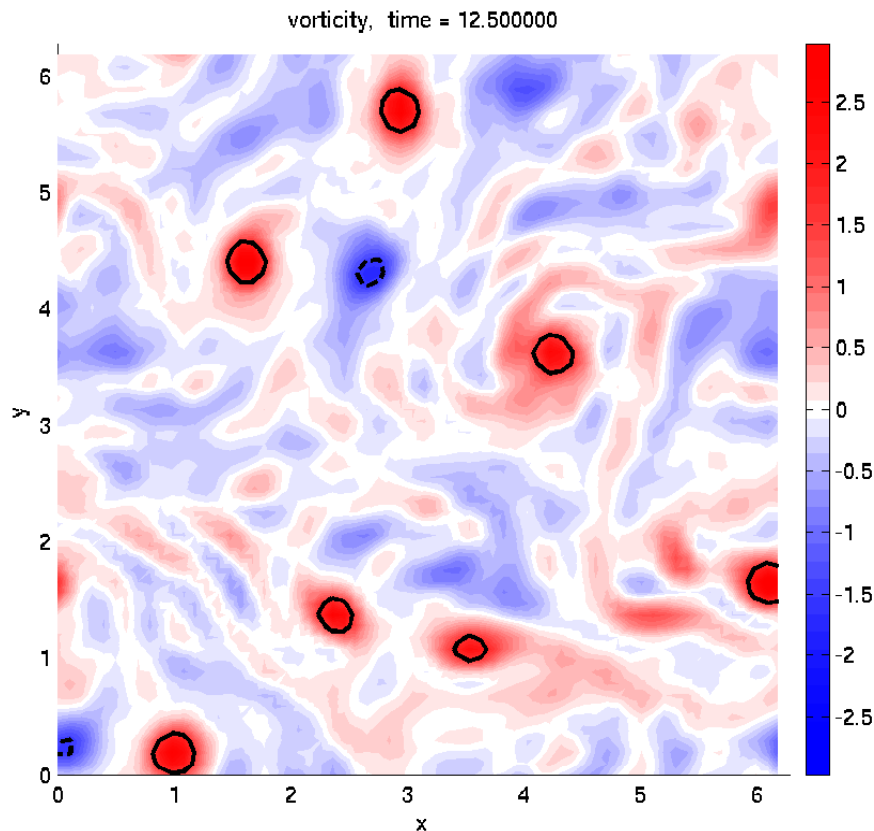


Fig. shows (symmetric) 2D decay with large-scale drag.

Embid & Majda 1996, 1998, Babin et al. 2002

## 3D Pure Rotation

all interactions with  $|\sigma_{s_k} + \sigma_{s_p} + \sigma_{s_q}| < Ro \implies$  cyclone dominance (but is not a PDE ! ).



## Reduced Models for 3D Boussinesq

Keeping only slow vortical mode interactions  $\Rightarrow$  the symmetric 3D quasi-geostrophic equation.

In Fourier space

$$\frac{d}{dt} b_0(t; \mathbf{k}) = \sum_{\Delta} C_{\mathbf{k}\mathbf{p}\mathbf{q}}^{000} b_0^*(\mathbf{p}) b_0^*(\mathbf{q})$$

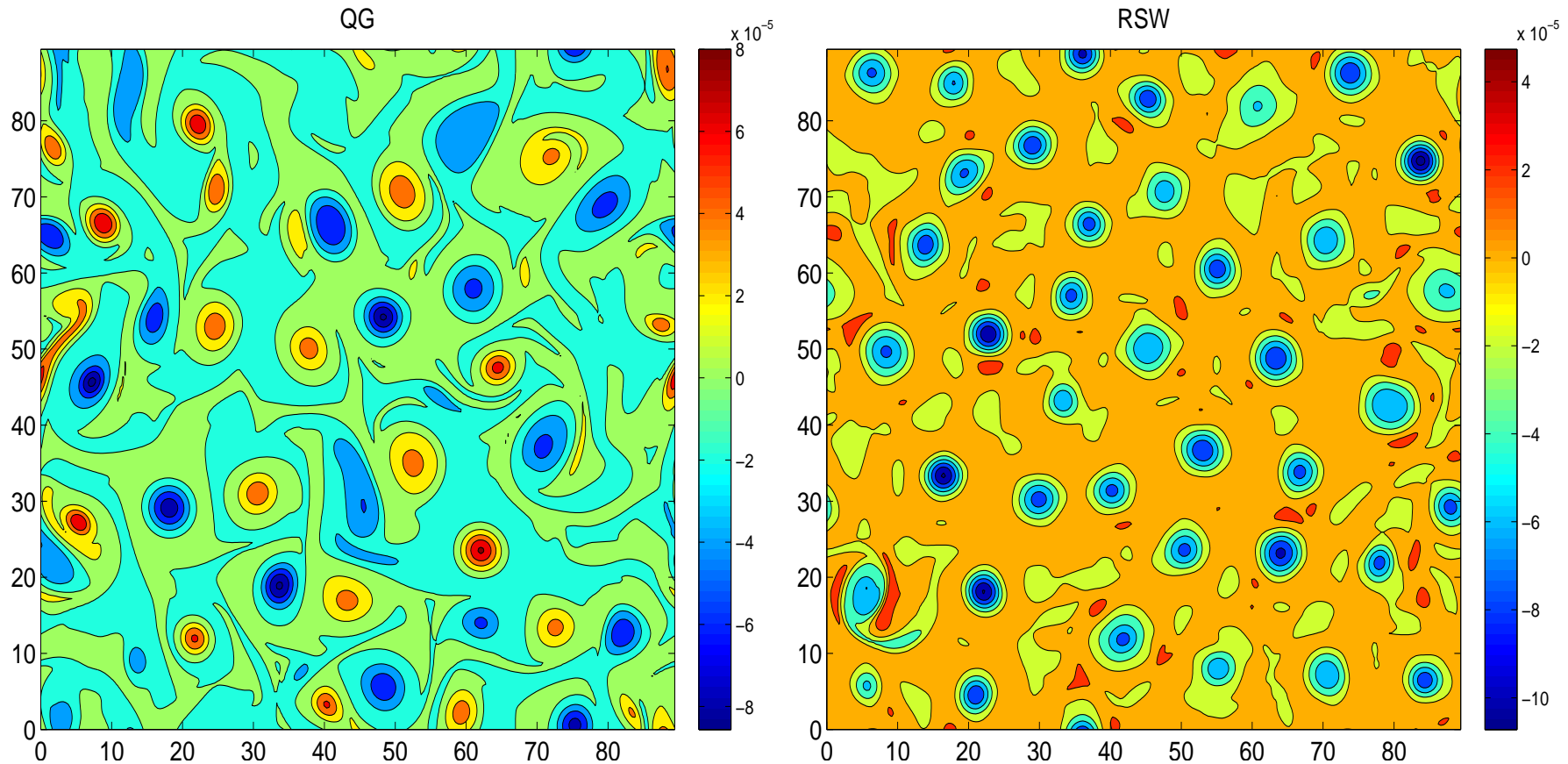
An inverse transform gives 3DQG

$$\left( \frac{\partial}{\partial t} + \mathbf{u}_H \cdot \nabla \right) q = 0, \quad q = \left( \nabla_H^2 + \frac{f^2}{N^2} \frac{\partial^2}{\partial z^2} \right) \psi(\mathbf{x}, t)$$

$$\nabla_H^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}, \quad \mathbf{u}_H = \hat{\mathbf{z}} \times \nabla \psi, \quad \theta = -\frac{f}{N} \frac{\partial \psi}{\partial z}$$

## Part II. Phenomena not captured by QG

### Example 1: Anticyclone dominance in RSW decay



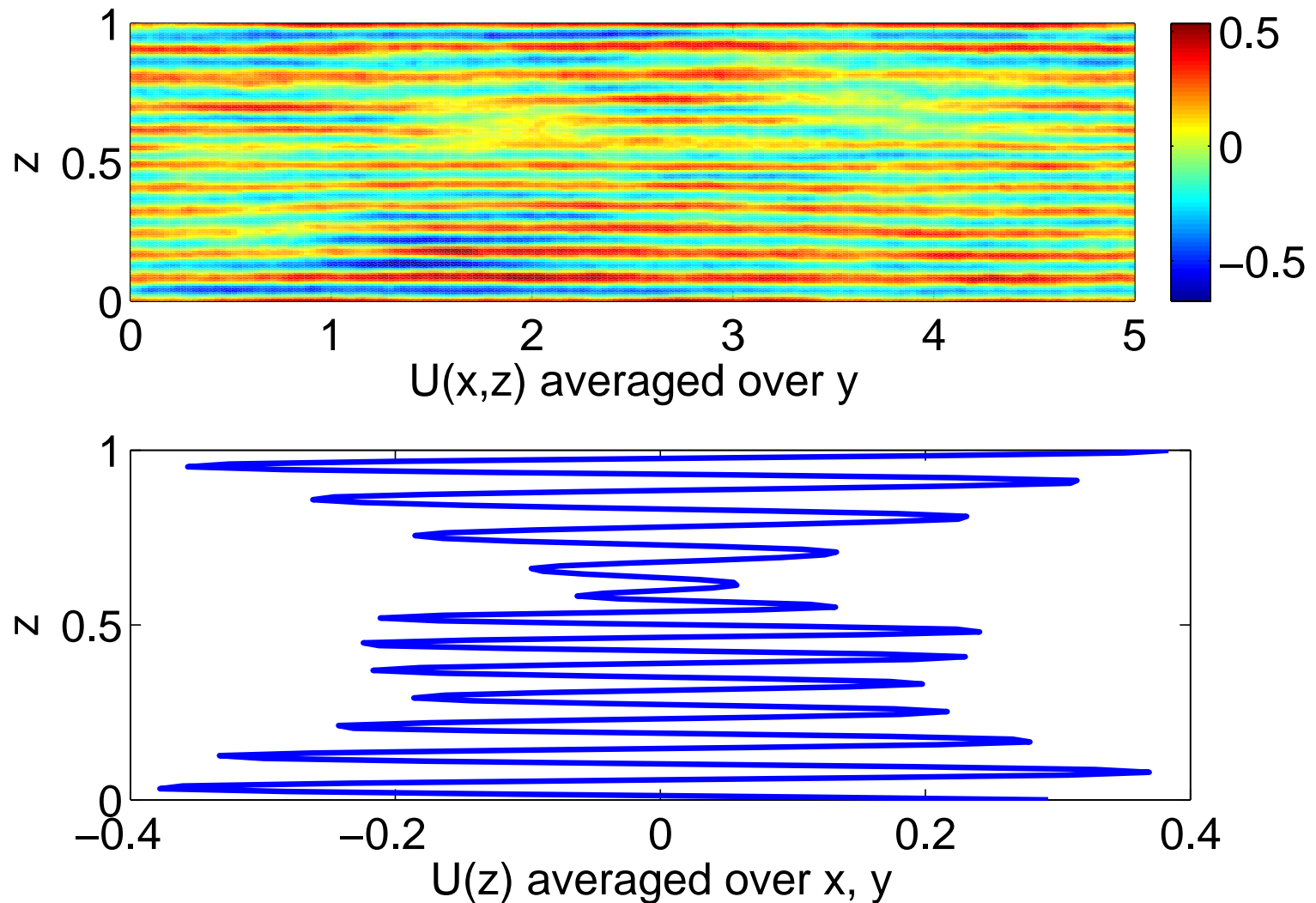
2DQG model

Full RSW

$$Ro = U(fL)^{-1} = 0.4, \quad Fr = U(gH)^{-1/2} = 0.25$$

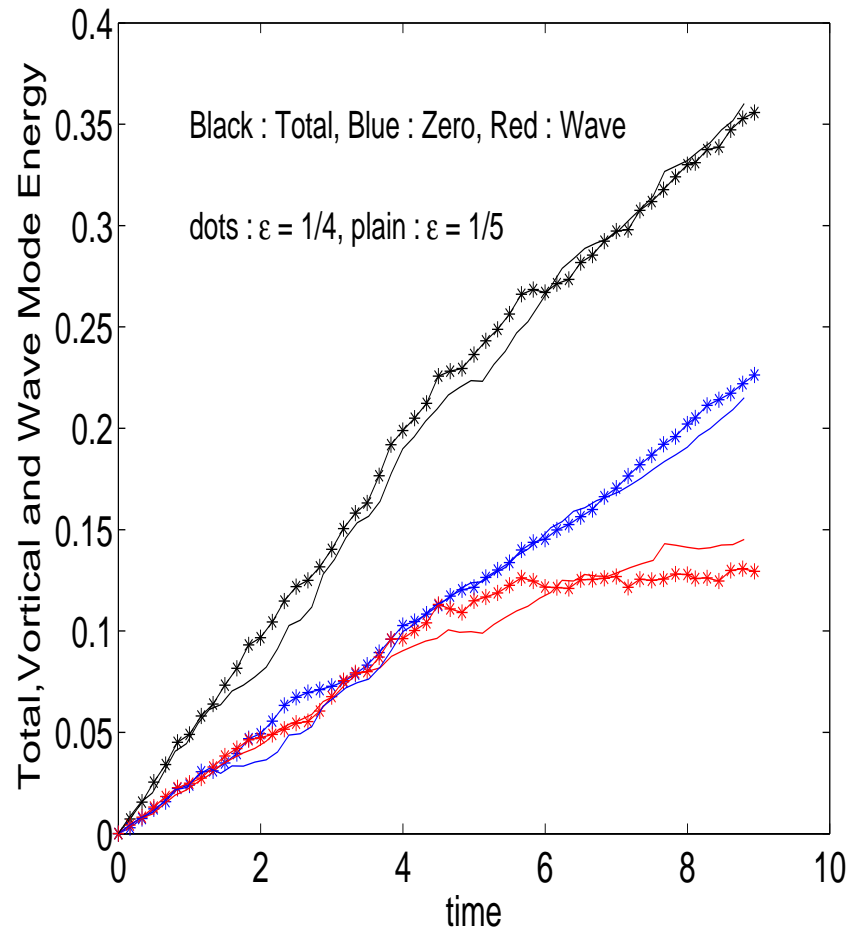
## Example 2: Generation of VSHF; $Bu = fL/(NH)=1$ ; $H/L=1/5$

$$Ro = U/(fL) = 0.1; Fr = U/(NH) = 0.1$$

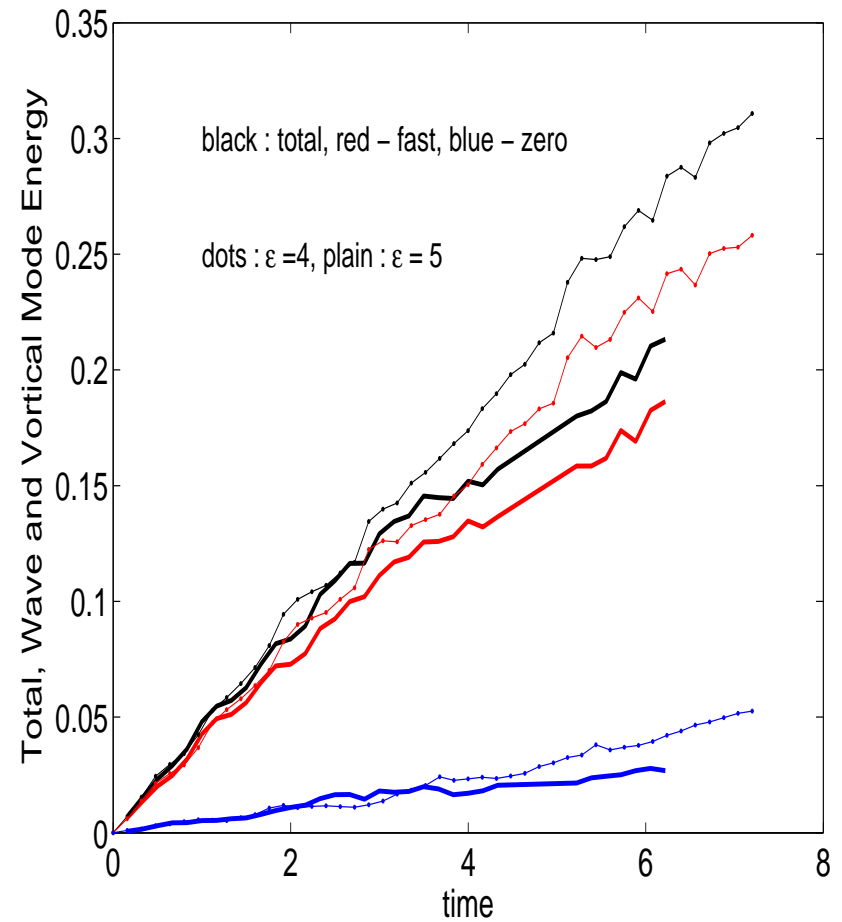




### Example 3: Asymmetry with respect to $f/N = 1$ (QG-like)



$f/N = 1/4, 1/5$



$f/N = 4, 5$

Bartello 1995, Sukhatme & Smith 2008

## Part III. PDE Intermediate Models

Intermediate models add more physics to QG

Improving upon 2DQG: Allen, Barth & Newberger 1990; Spall & McWilliams 1992; Yavneh & McWilliams 1994; Warn, Bokhove, Shepherd & Vallis 1995; Vallis 1996

Improving upon 3DQG: Allen 1991, 1993; Muraki, Snyder & Rotunno 1999; Muraki & Hakim 2001

Previous intermediate models are perturbative in nature

# A hierarchy of NEW Intermediate Models

- Derived by adding subsets of wave-vortical mode interactions to QG
- Non-perturbative
- Include near-resonant triads
- Provides a framework for understanding the coupling between balanced and unbalanced flow components  
e.g. Kuo, Allen, Polvani 99; Ford, McIntyre, Norton 00; Majda 03

## The Full Equations:

$$0 \quad | \quad 00 \quad \oplus \quad 0+ \quad \oplus \quad 0- \quad \oplus \quad ++ \quad \oplus \quad +- \quad \oplus \quad -- \quad (1)$$

$$+ \quad | \quad 00 \quad \oplus \quad 0+ \quad \oplus \quad 0- \quad \oplus \quad ++ \quad \oplus \quad +- \quad \oplus \quad -- \quad (2)$$

$$- \quad | \quad 00 \quad \oplus \quad 0+ \quad \oplus \quad 0- \quad \oplus \quad ++ \quad \oplus \quad +- \quad \oplus \quad -- \quad (3)$$

QG (vortical mode interactions only):

$$0 \quad | \quad 00$$

## Two NEW Models

PPG (add to QG interactions involving exactly 1 wave):

$$\begin{array}{l}
 0 \quad | \quad 00 \quad \oplus \quad 0 + \quad \oplus \quad 0 - \quad (ppg1) \\
 + \quad | \quad 00 \quad \text{and} \quad - \quad | \quad 00 \quad (ppg2, 3)
 \end{array}$$

P2G (add to PPG interactions involving exactly 2 waves):

$$\begin{array}{l}
 0 \quad | \quad 00 \quad \oplus \quad 0 + \quad \oplus \quad 0 - \quad \oplus \quad + + \quad \oplus \quad + - \quad \oplus \quad - - \quad (p2g1) \\
 + \quad | \quad 00 \quad \oplus \quad 0 + \quad \oplus \quad 0 - \quad (p2g2) \\
 - \quad | \quad 00 \quad \oplus \quad 0 + \quad \oplus \quad 0 - \quad (p2g3)
 \end{array}$$

# QG and PPG Rotating Shallow Water (RSW) Equations

$$\text{QG: } \partial Q / \partial t + J(\Psi, Q) = 0$$

$$\text{PPG : } \frac{\partial \nabla^2 \chi}{\partial t} - \nabla^2 V = 2J\left(\frac{\partial A}{\partial x}, \frac{\partial A}{\partial y}\right), \quad (1)$$

$$\frac{\partial Q}{\partial t} + J(\Psi, Q) + \nabla \chi \cdot \nabla Q + \langle u \rangle \frac{\partial Q}{\partial x} + \langle v \rangle \frac{\partial Q}{\partial y} + Q \nabla^2 \chi = 0, \quad (2)$$

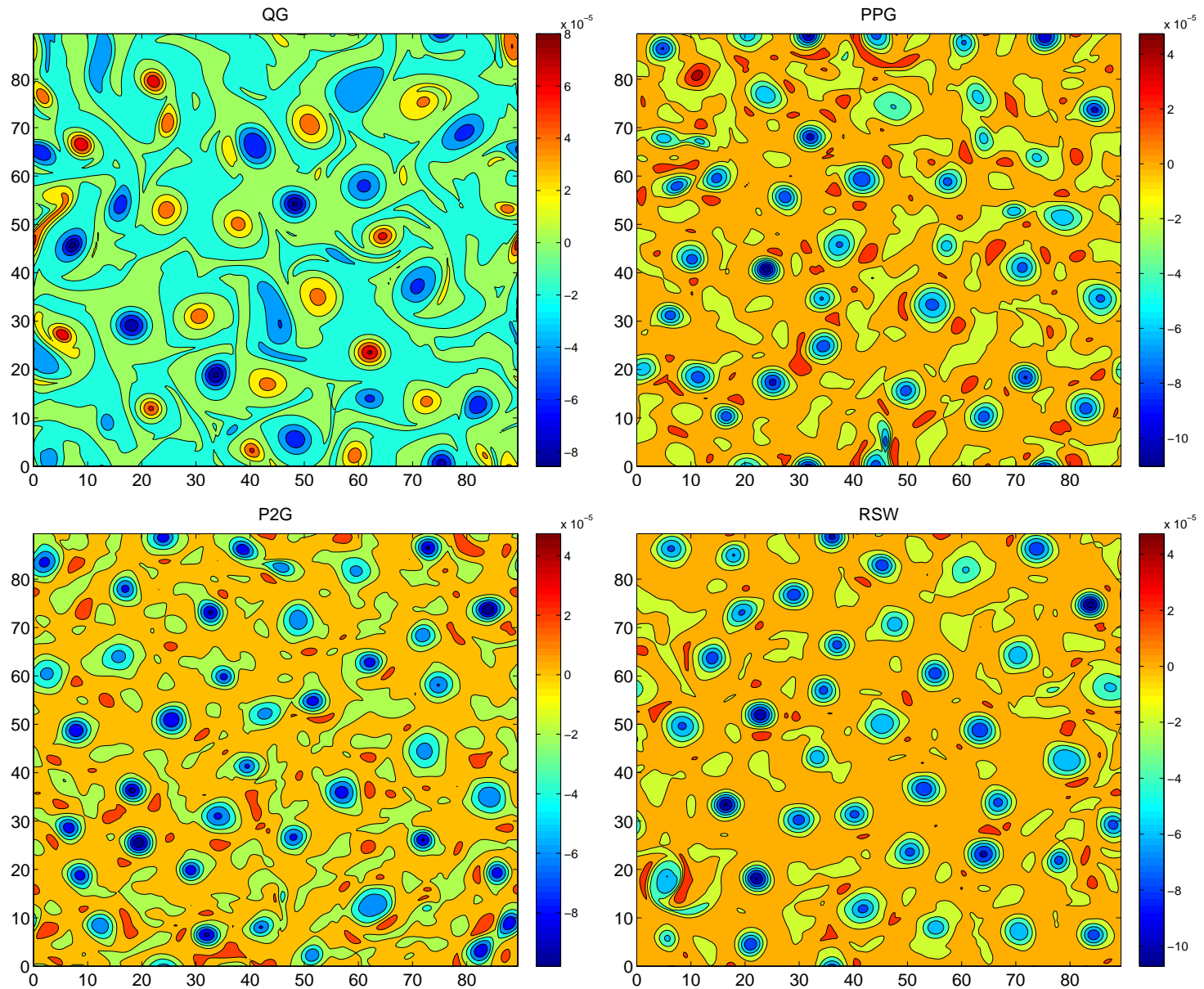
$$\frac{\partial \nabla^2 V}{\partial t} - c^2 \nabla^4 \chi + f^2 \nabla^2 \chi = f J(A, Q) \quad (3)$$

$$Q = \left(\nabla^2 - \frac{f^2}{gH}\right)\Psi, \quad u = \chi_x - \Psi_y, \quad v = \chi_y + \Psi_x, \quad \nabla^2 \chi = u_x + v_y$$

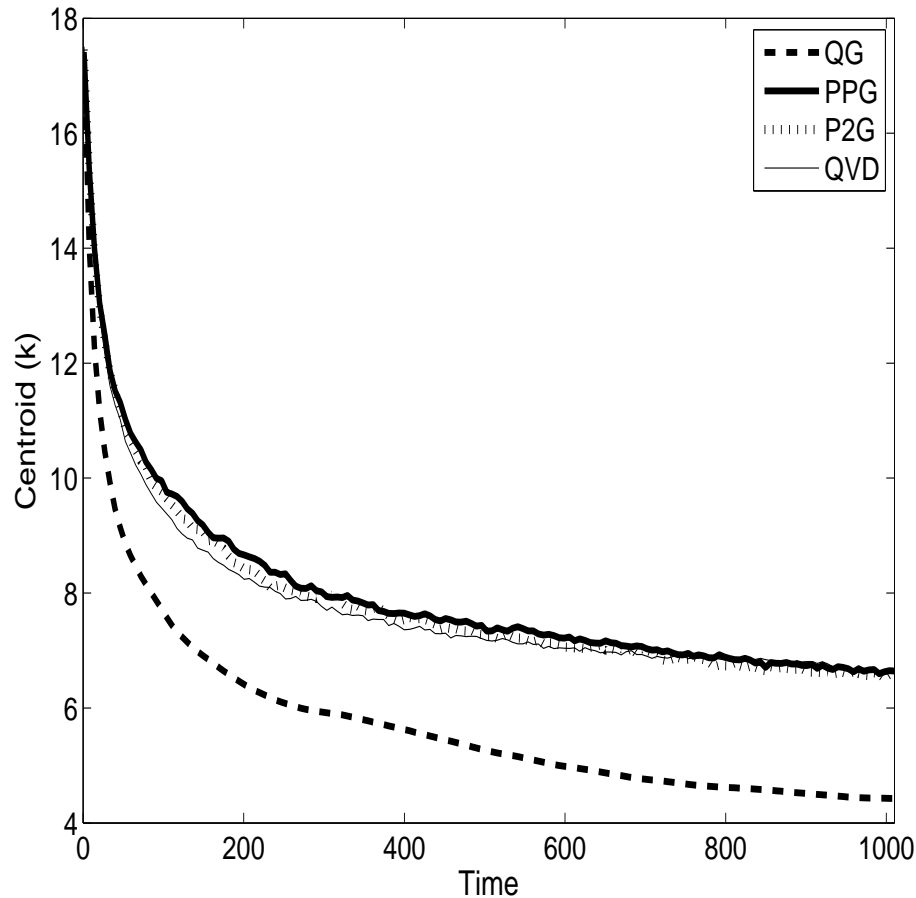
$\nabla^2 V = \nabla^2(f\Psi - gh)$  is a measure of geostrophic imbalance (Vallis 96); also called geostrophic departure (Warn 95); ageostrophic vorticity (Mohebalhojeh & Dritschel 01)

$$A \equiv (f^2 - c^2 \nabla^2)^{-1} c^2 Q$$

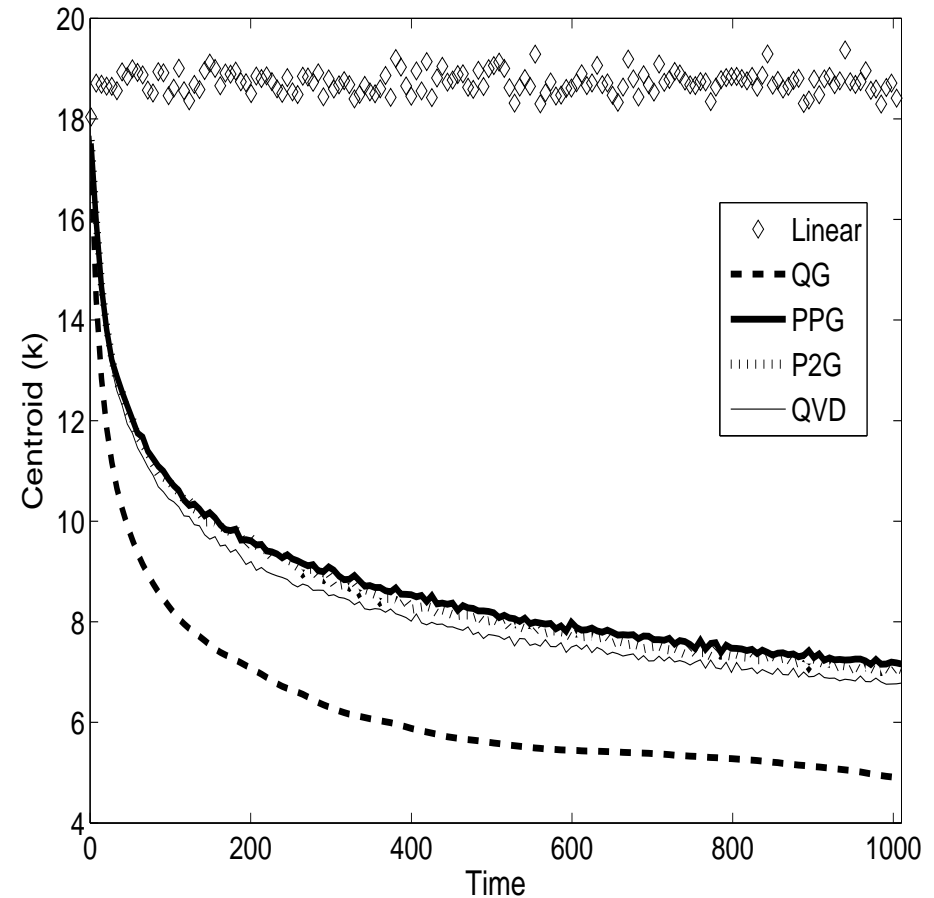
# RSW decay, $Ro=0.4$ , $Fr = 0.25$ , divergence-free unbalanced i.c.



# Centroid in RSW decay; divergence-free unbalanced i.c.



$Ro = 0.4, Fr = 0.25$

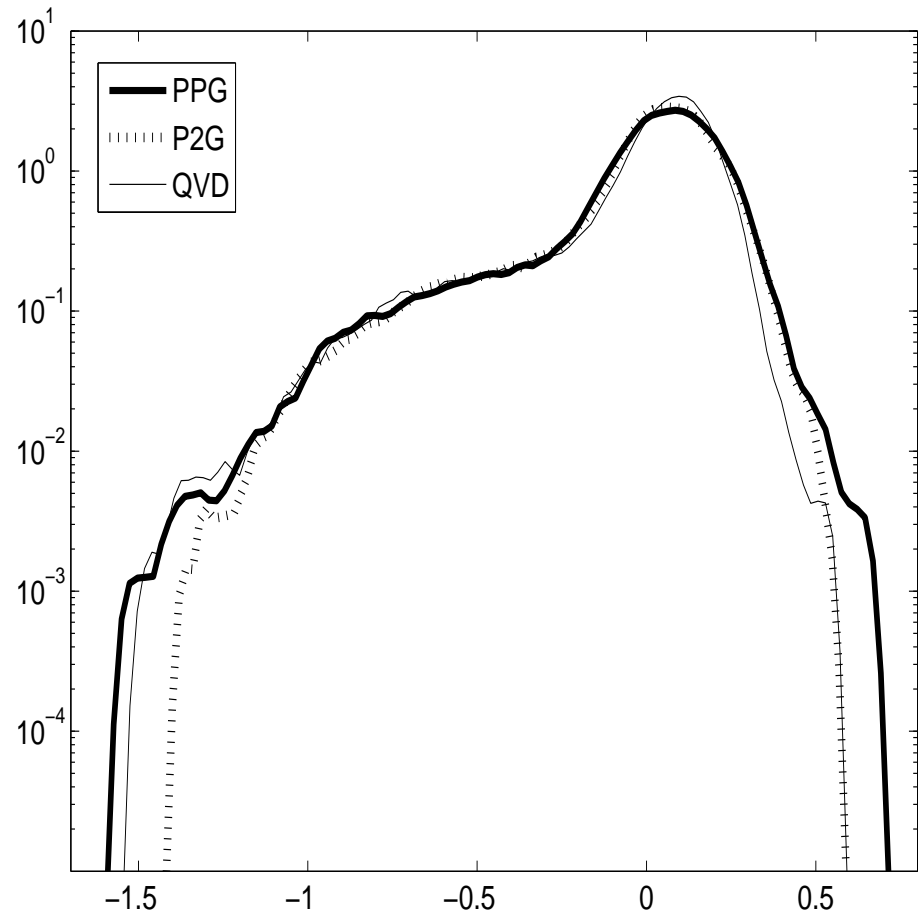
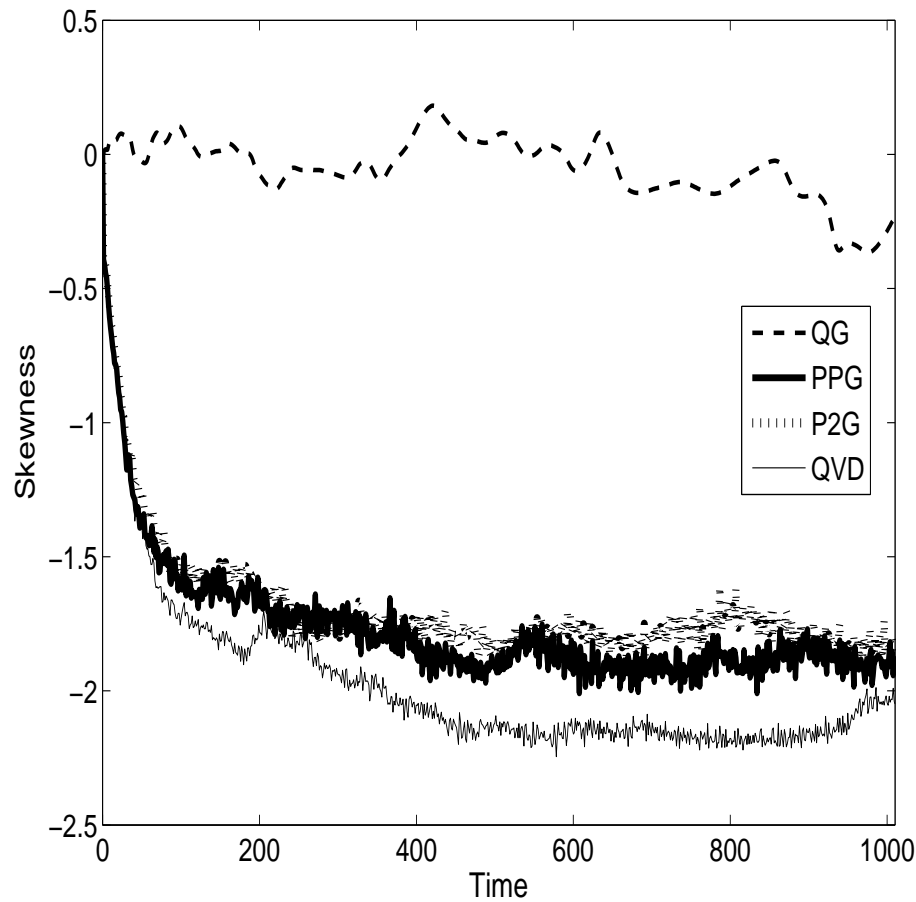


$Ro = 0.25, Fr = 0.2$

$$\text{Cent}(k) = (\sum_{\mathbf{k}} k(|u_{\mathbf{k}}|^2 + |v_{\mathbf{k}}|^2)) / \sum_{\mathbf{k}} (|u_{\mathbf{k}}|^2 + |v_{\mathbf{k}}|^2)$$



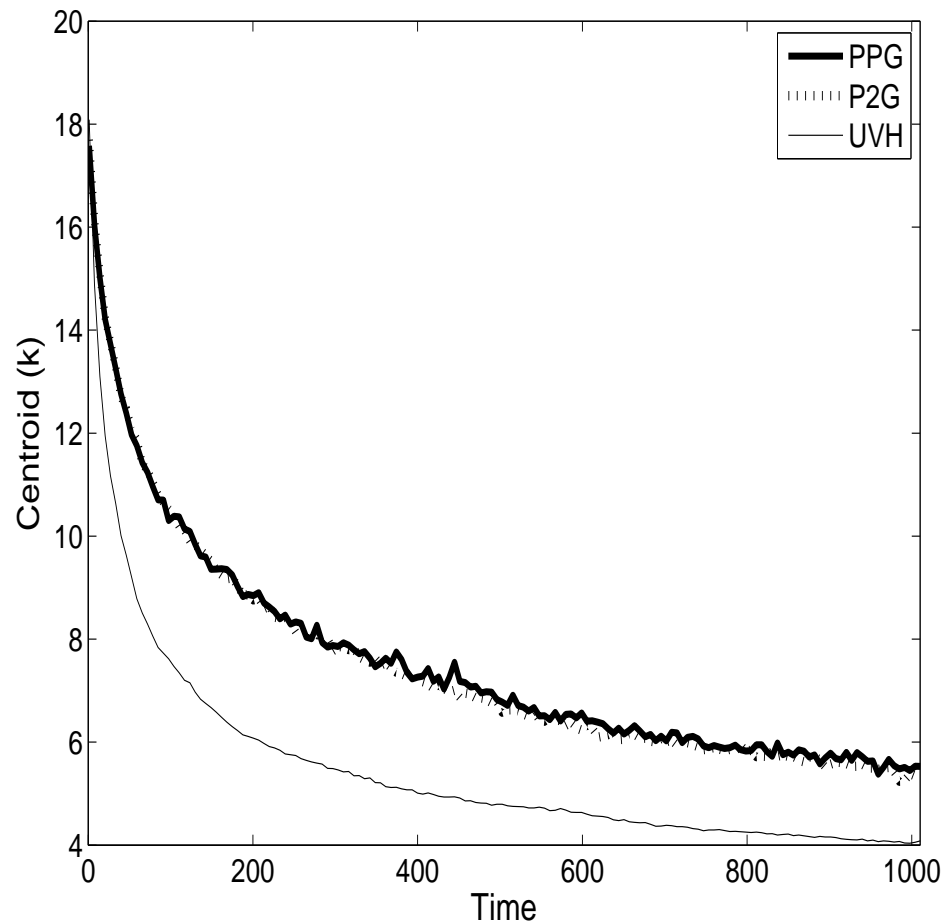
# Vorticity skewness in RSW decay; divergence-free unbalanced i.c.



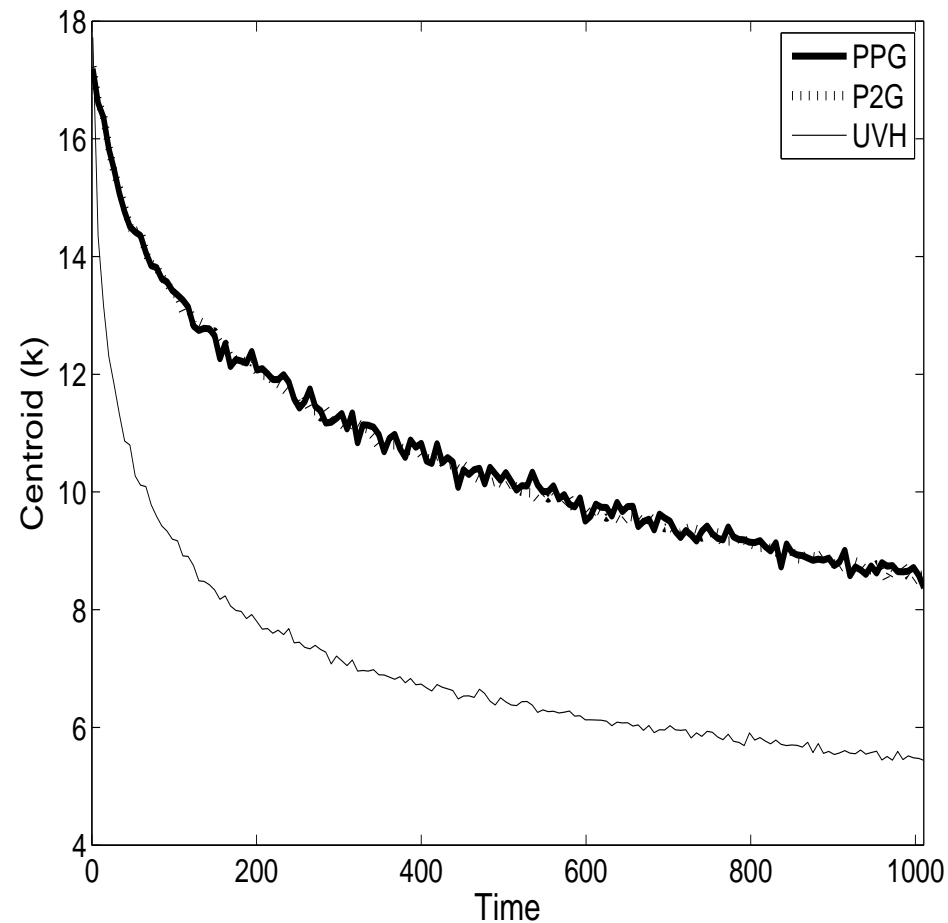
$$Ro = 0.4, Fr = 0.25$$

# Centroid in RSW decay with divergent initial conditions

$$Ro = 1, Fr = 0.3$$



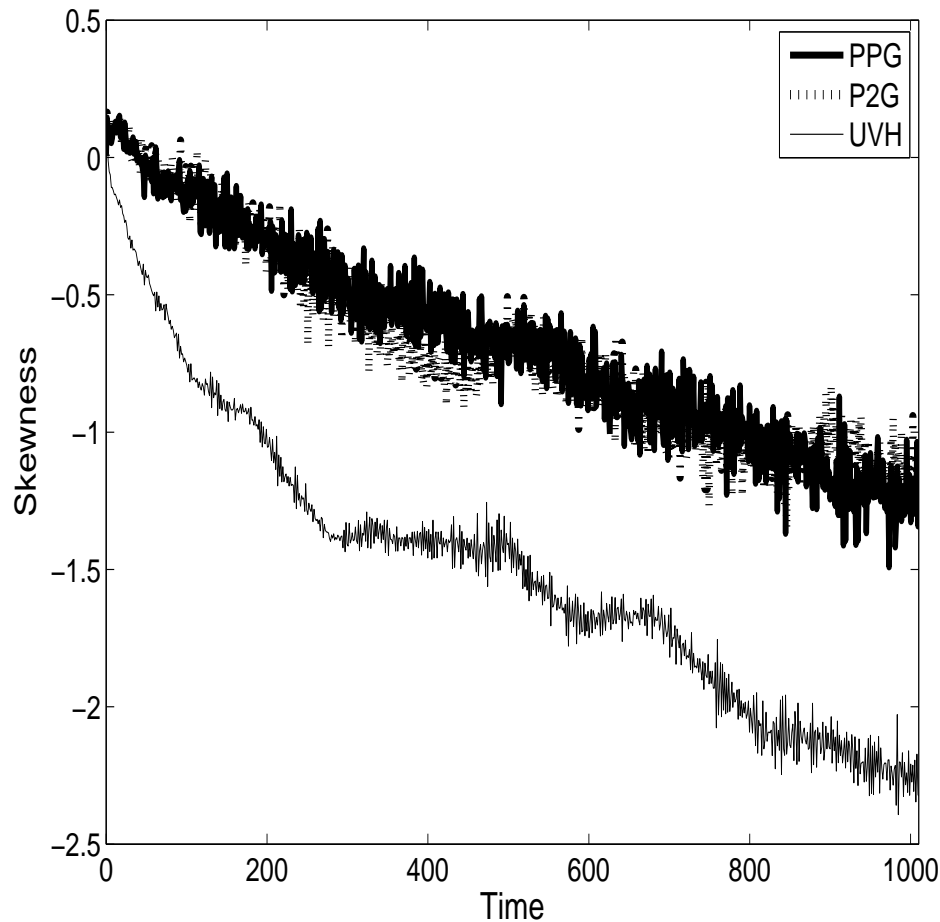
25% divergence-free



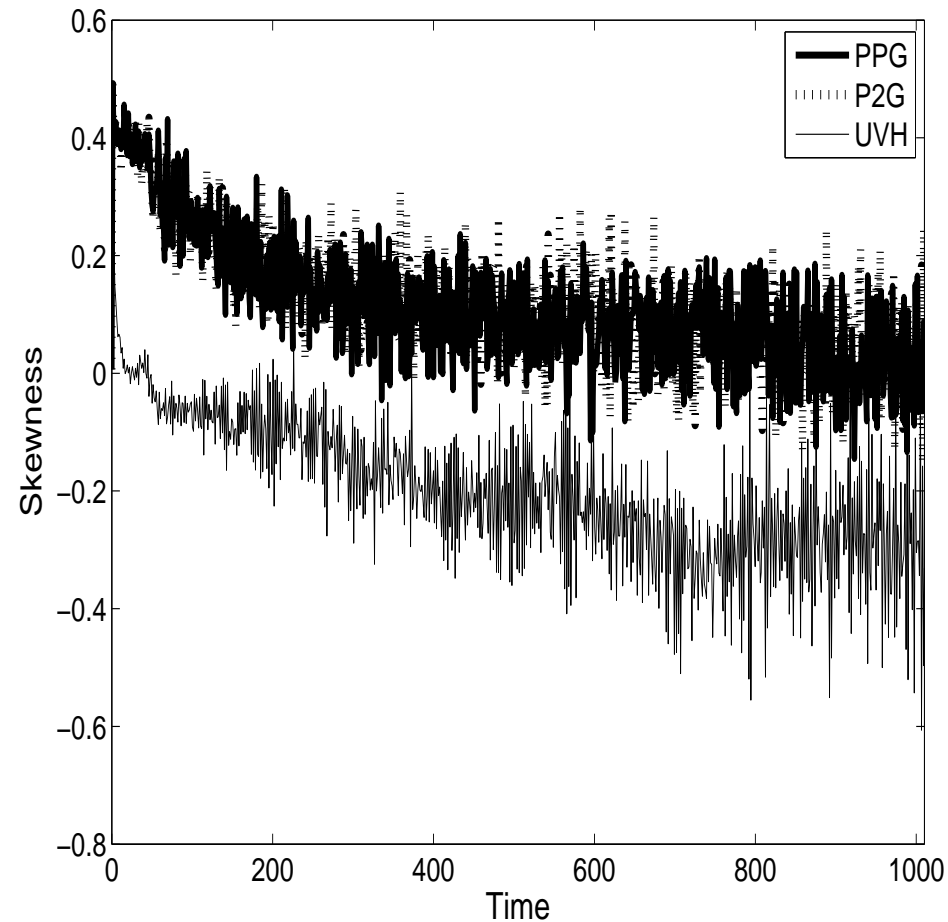
5% divergence-free

# Skewness in RSW decay with divergent initial conditions

$$Ro = 1, Fr = 0.3$$



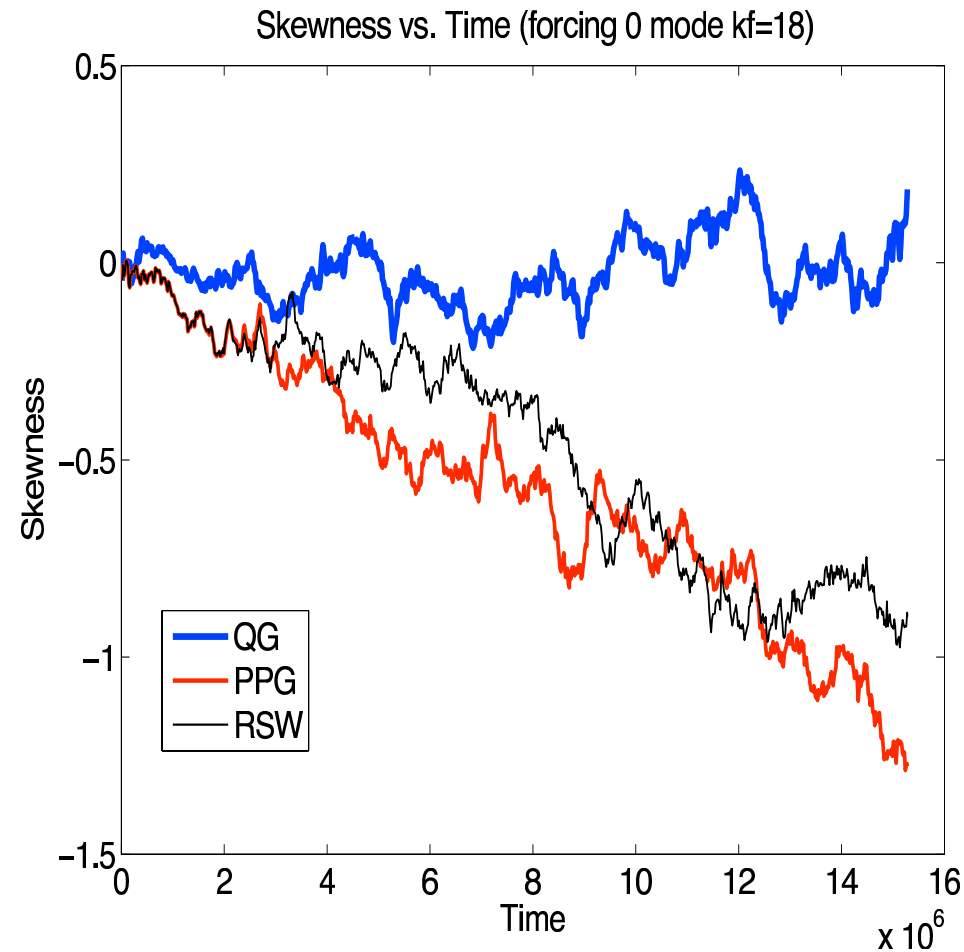
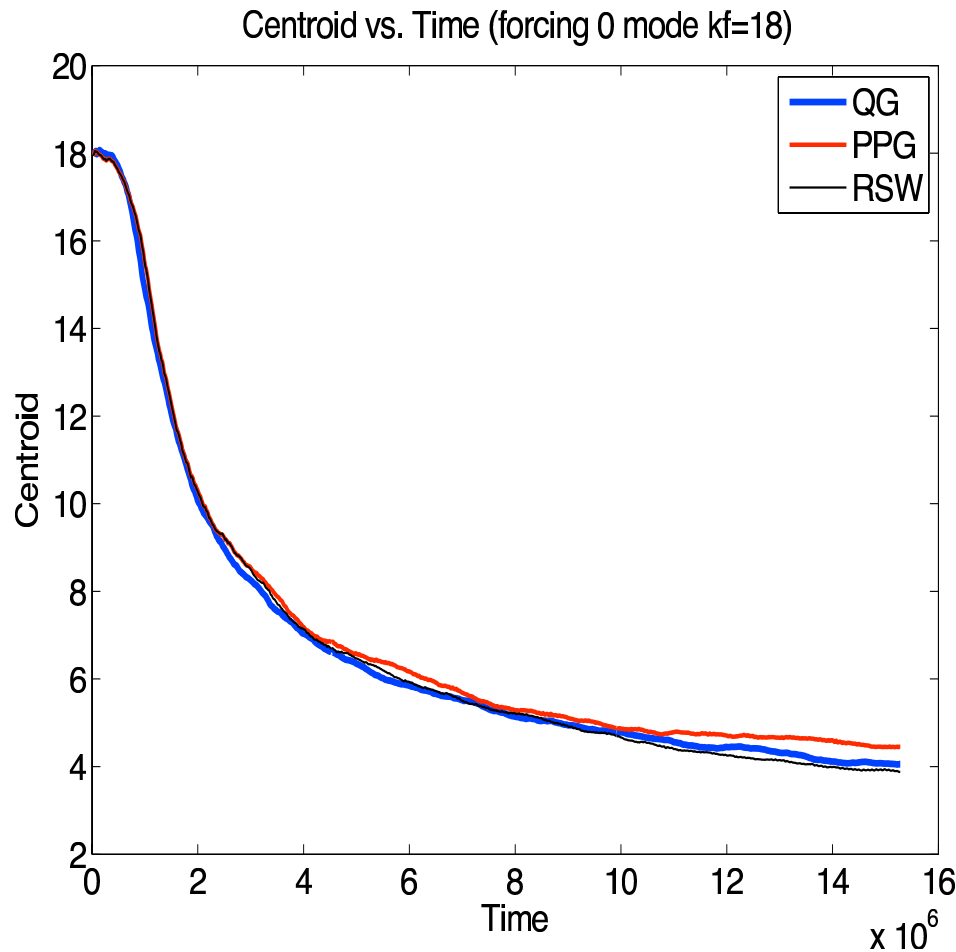
25% divergence-free



5% divergence-free

# Anticyclones in forced RSW (white-noise forcing of PV modes)

$$Ro = \varepsilon^{1/3} (k_f / \pi)^{2/3} f^{-1} = 0.05, \quad Fr = (\varepsilon \pi)^{1/3} k_f^{-1/3} (gH)^{-1/2} = 0.03$$



## Differences between RSW and RB

- Unlike for RB, truncations of RSW are not guaranteed to maintain energy as an integral invariant (Warn 86).
- RSW intermediate models with both energy conservation and Lagrangian invariance of PV do not perform as well as models conserving only PV (Allen 93).
- Model performance is not necessarily linked to conservation properties (Allen 93).

## Differences between RSW and RB

- In RSW, two wave modes cannot transfer energy into or out of a vortical mode, i.e.

$$C_{\mathbf{k}\mathbf{p}\mathbf{q}}^{0+-} = C_{\mathbf{k}\mathbf{p}\mathbf{q}}^{0++} = C_{\mathbf{k}\mathbf{p}\mathbf{q}}^{0--} = 0$$

(true only for exact resonances in RB)

- In RSW the catalytic resonances with  $(s_{\mathbf{k}}, s_{\mathbf{p}}, s_{\mathbf{q}}) = (\pm, 0, \pm)$  only exchange energy in the same wavenumber shell
- In RSW, there are no exact 3-wave resonances (Majda 96)

# Rotating Boussinesq Equations

- In RB, two wave modes can transfer energy into or out of a vortical mode, except for exact resonances.
  - In RB the catalytic resonances with  $(s_{\mathbf{k}}, s_{\mathbf{p}}, s_{\mathbf{q}}) = (\pm, 0, \pm)$  exchange energy between scales (on a cone in wavevector space)
  - In RB, there are exact 3-wave resonances
-

# Rotating Boussinesq Equations

- In RB, two wave modes can transfer energy into or out of a vortical mode, except for exact resonances.
- In RB the catalytic resonances with  $(s_{\mathbf{k}}, s_{\mathbf{p}}, s_{\mathbf{q}}) = (\pm, 0, \pm)$  exchange energy between scales (on a cone in wavevector space)
- In RB, there are exact 3-wave resonances

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Richer behavior in RB from interactions involving one vortical mode and two wave modes (including catalytic exact/near resonances), and from 3-wave exact/near resonances.



## Testable prediction for RB

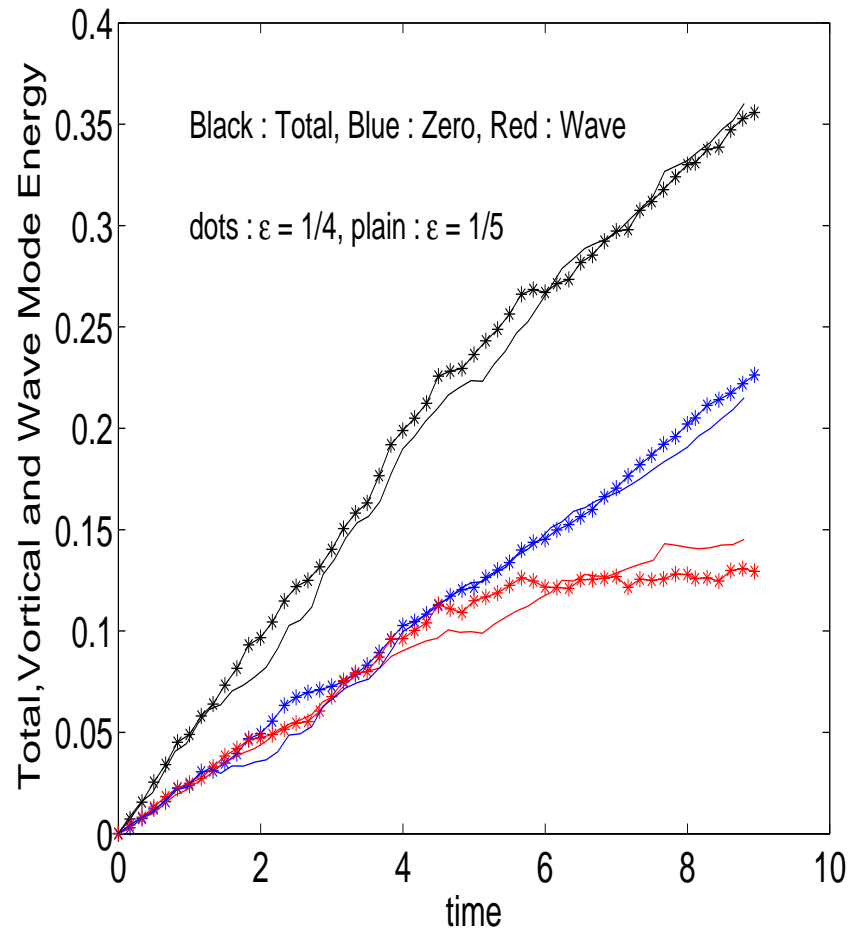
We predict the P2SG model can reproduce asymmetry with respect to  $f/N = 1$  in the forward transfer range:

$$0 \quad | \quad 00 \quad \oplus \quad + + \quad \oplus \quad + - \quad \oplus \quad - - \quad (1)$$

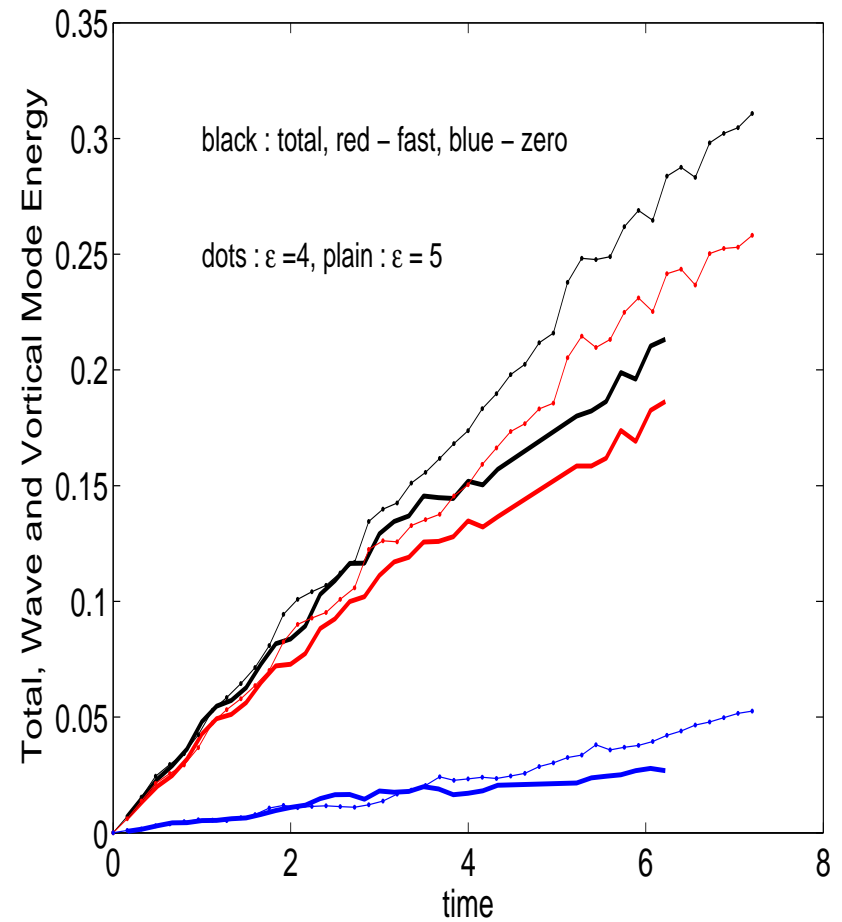
$$+ \quad | \quad 0 + \quad \oplus \quad 0 - \quad (2)$$

$$- \quad | \quad 0 + \quad \oplus \quad 0 - \quad (3)$$

### Example 3: Asymmetry with respect to $f/N = 1$ (QG-like)



$f/N = 1/4, 1/5$



$f/N = 4, 5$

## Conclusions/Discussion

- A nonperturbative method to derive a hierarchy of models including wave-vortical interactions
- Provides a framework for complete understanding of balanced and imbalanced flow components
- Interactions involving exactly one gravity wave responsible for anticyclone dominance in RSW
- Stay tuned for results on RB