

# Effect of near-inertial modes on the midlatitude double gyre problem

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David Straub, McGill

## OUTLINE

1. Motivation
2. Model Setup
3. Results
4. Conclusion

Collaborator: Aaron Gertz

# Motivation: Balanced-to-Unbalanced Energy Transfers

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- Simple example: take 2d turbulence as “balanced flow”; and 3d modes as “unbalanced”
- More involved: take QG turbulence as “balanced”; and inertia-gravity modes as “unbalanced” (e.g., Ngan *et al.*, 2007)

# Motivation: Balanced-to-Unbalanced Energy Transfers

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- Simple example: take 2d turbulence as “balanced flow”; and 3d modes as “unbalanced”
  - unstratified thin (hydrostatic) fluid
  - but allowing  $u$  to vary in  $z$
  - double gyre forcing on depth averaged mode
  - large scale stochastic forcing on vertically varying mode

# Model Equations

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$$u_t + \nabla \cdot (\mathbf{v}u) - fv = -P_x + F^x - D^x$$

$$v_t + \nabla \cdot (\mathbf{v}v) + fu = -P_y - D^y$$

$$\mathbf{v} \equiv (u, v, w) \quad \nabla_3 \cdot \mathbf{v} = 0 \quad P_z = 0$$

- dynamics: hydrostatic (but unstratified)
- domain: 4000km square box ;  $513 \times 513 \times 2$
- $f = f_0 + \beta y$  with  $f_0 = 7.5 \times 10^{-5} \text{s}^{-1}$ ,  $\beta = 2 \times 10^{-11} \text{m}^{-1} \text{s}^{-1}$
- hyperviscosity ( $\nabla^6$ ) with slip conditions at boundaries
- Rayleigh drag on 2d mode
- Forcing: double gyre for 2d mode. Stochastic (and large scale) for 3d modes
- $\delta_i = (U_{Sv}/\beta)^{1/2} \sim 44\text{km} \sim 5.5\text{dx}$

# Model Equations

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$$\nabla^2 \psi_t + J(\psi, \nabla^2 \psi + \beta y) = -\overline{\mathbf{v}' \cdot \nabla \zeta'} - \overline{\mathbf{v}'_x \cdot \nabla v'} + \overline{\mathbf{v}'_y \cdot \nabla u'} + F - D$$

$$u'_t + \mathbf{v} \cdot \nabla u - \overline{\mathbf{v} \cdot \nabla u} - f v' = F^x - D$$

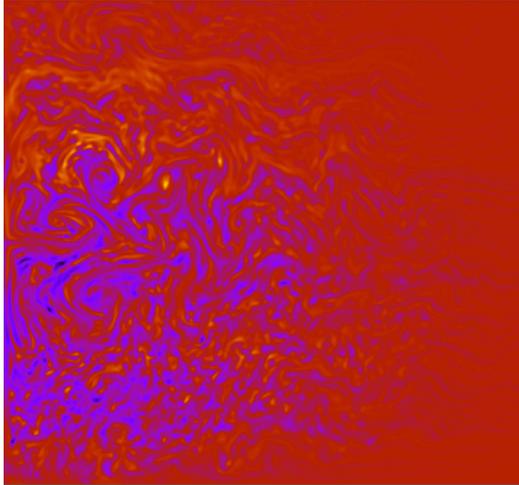
$$v'_t + \mathbf{v} \cdot \nabla v - \overline{\mathbf{v} \cdot \nabla v} + f u' = -D$$

- C-grid
- Multigrid method for elliptic inversion (L-P Nadeau)
- $F^x = a(t) \sin(y/L) e^{-(x^2 + y^2)/(0.5L)^2}$
- $a(t)$  given by an Ornstein-Uhlenbeck process (with characteristic timescale of 1 day)

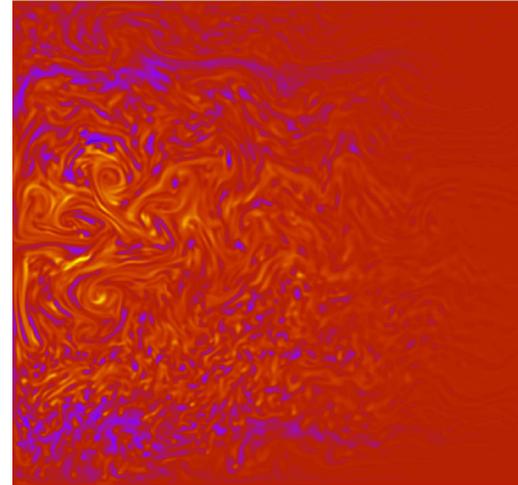
# A day in the life of $u'$

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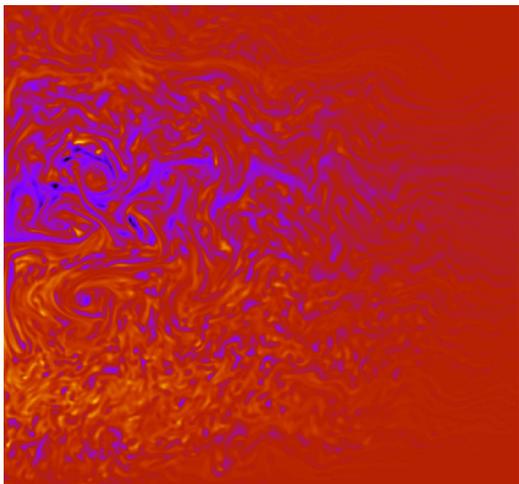
NOON



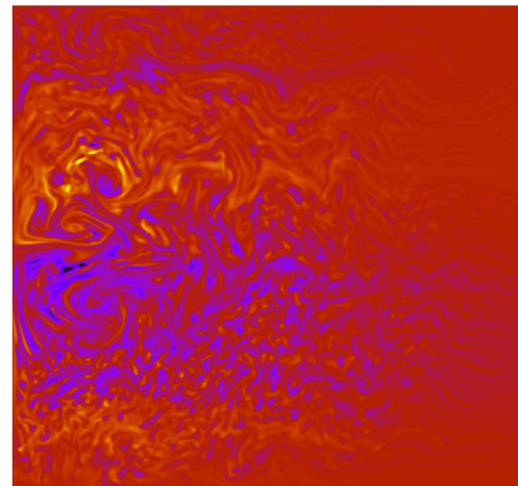
8PM



4AM



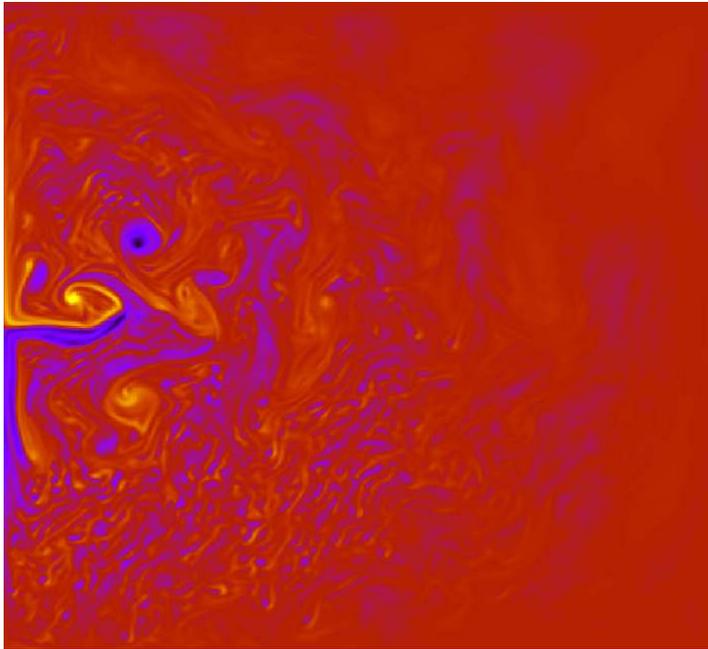
NOON (next day)



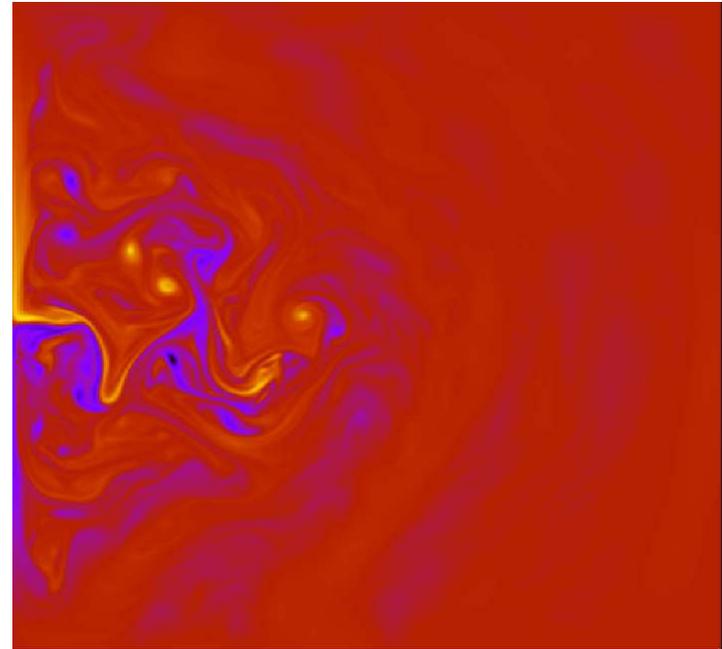
# Vertical Vorticity ( $r_{\text{Rayleigh}} = 5 \times 10^{-8} \text{s}^{-1}$ )

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With 3d Forcing

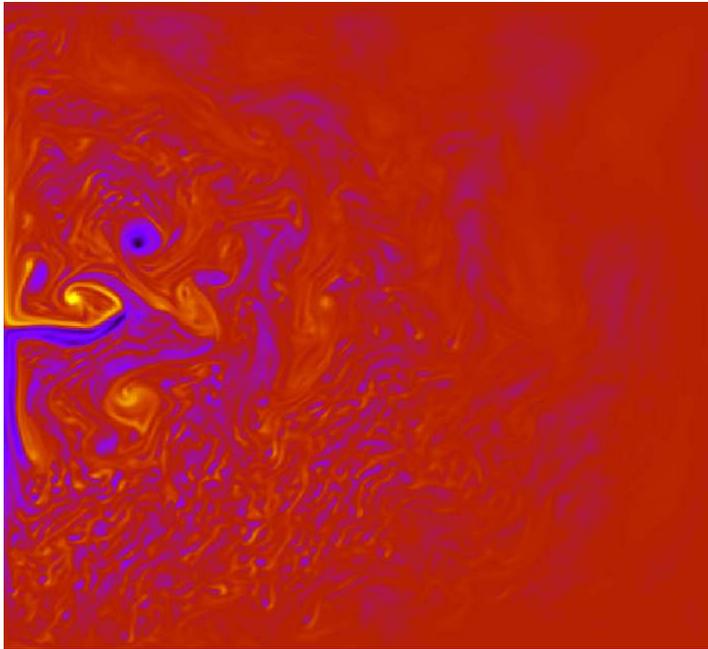


Without 3d Forcing

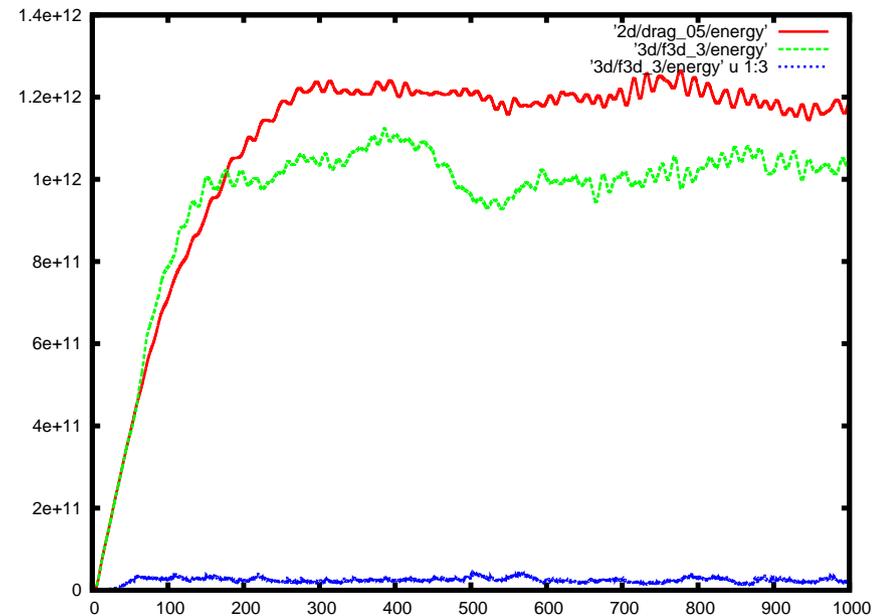


# Vertical Vorticity and Energy Time Series

With 3d Forcing



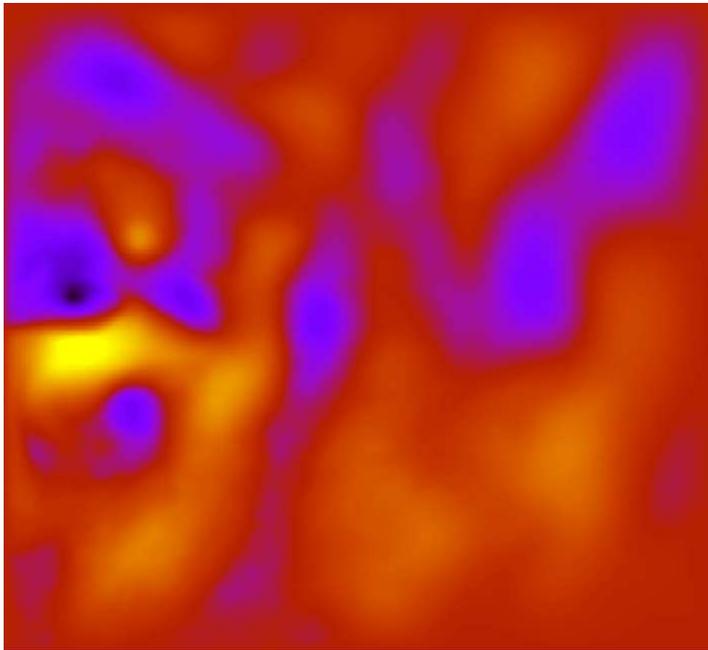
$E(t)$



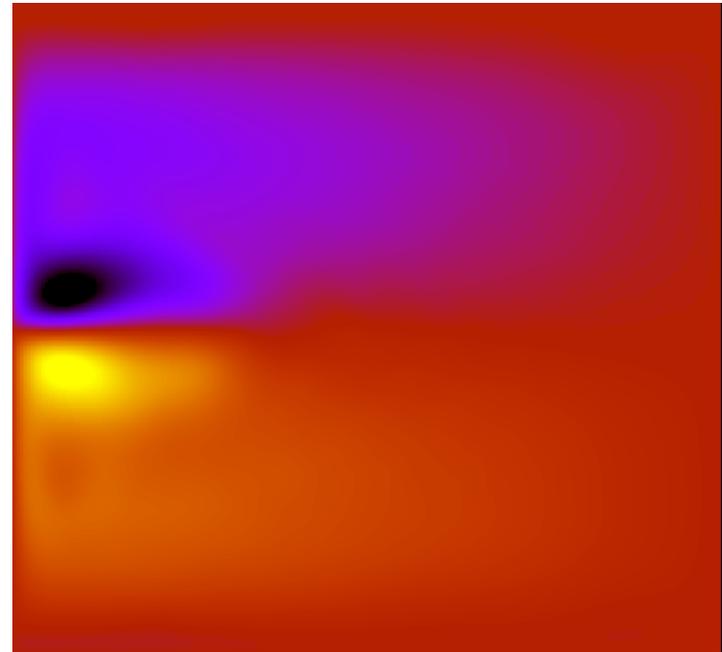
# Streamfunctions

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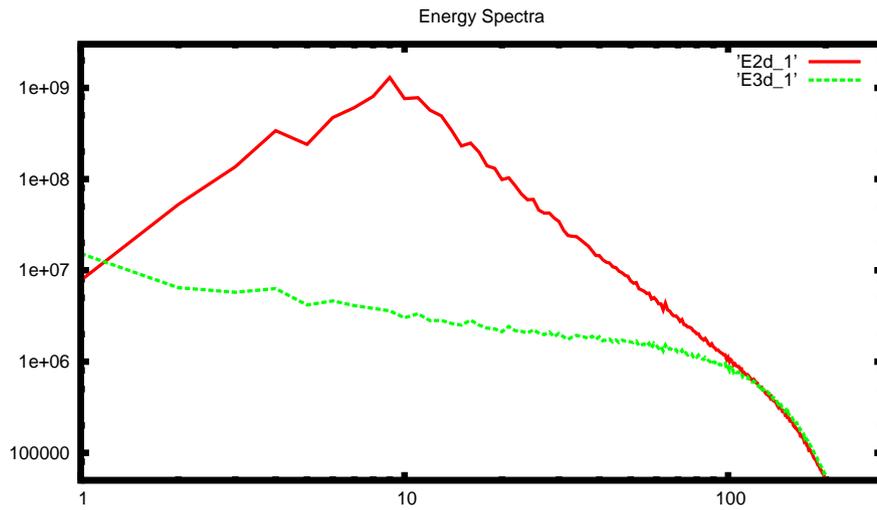
Typical Streamfunction



Time mean

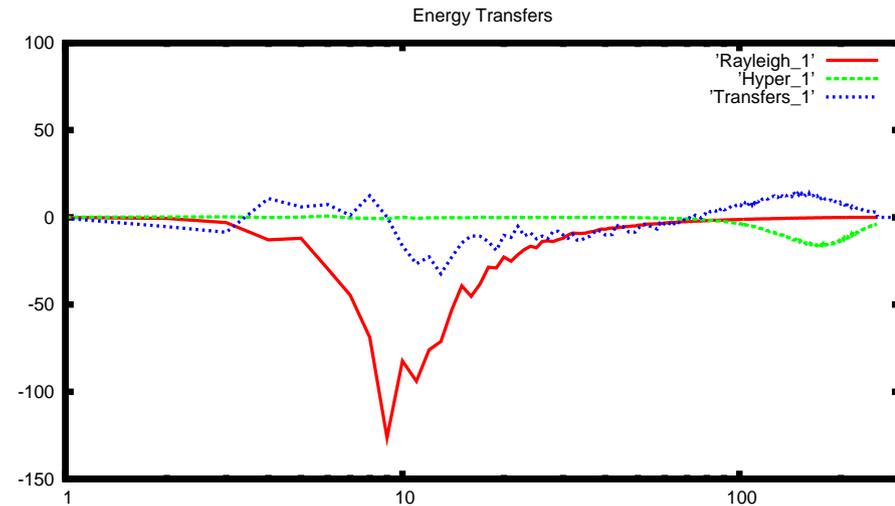


# Power, dissipation and transfer spectra



Red:  $E_{2d}(k)$

Green:  $E_{3d}(k)$

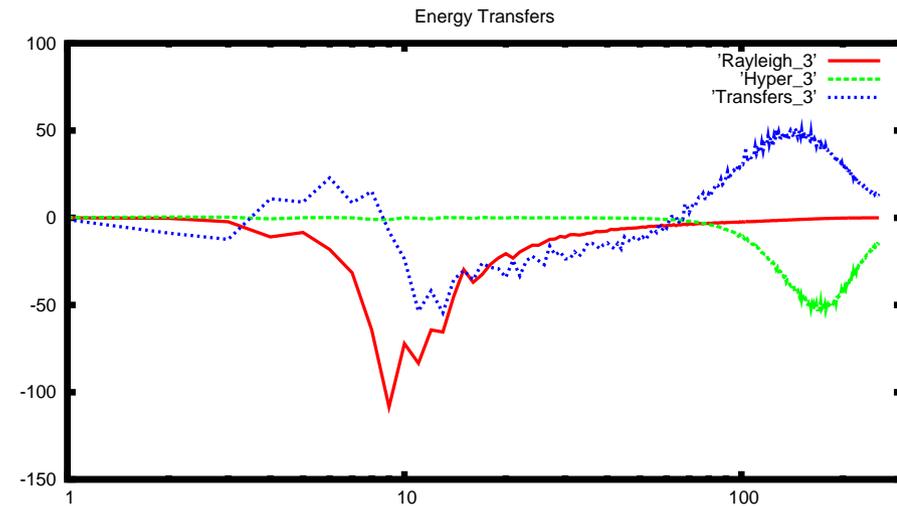
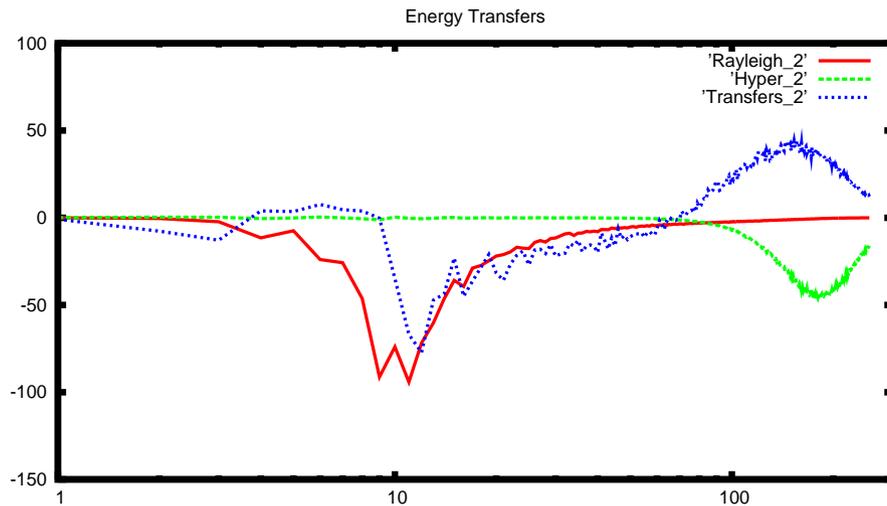
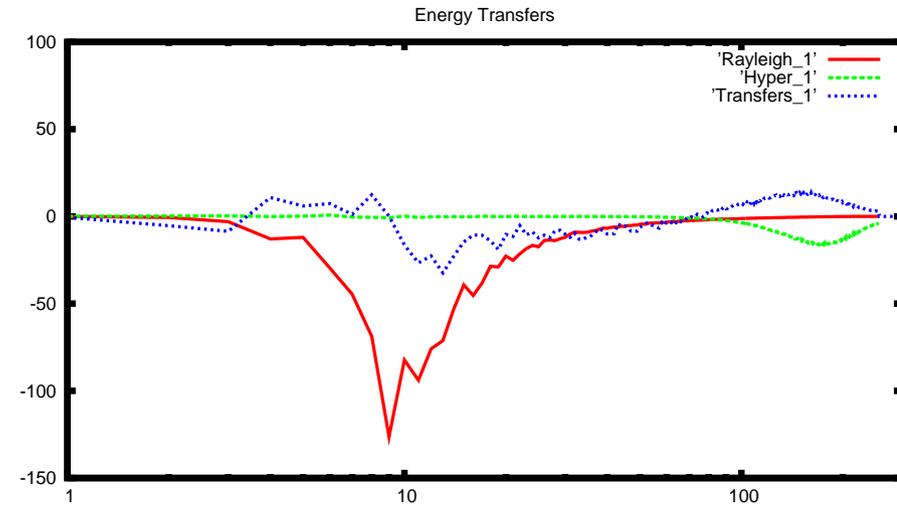
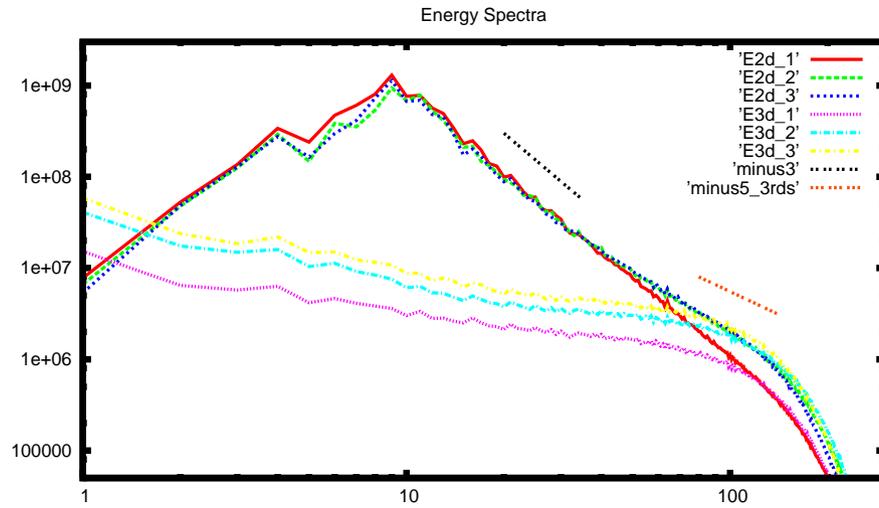


Red: Rayleigh Drag

Blue: 2d-to-3d transfers

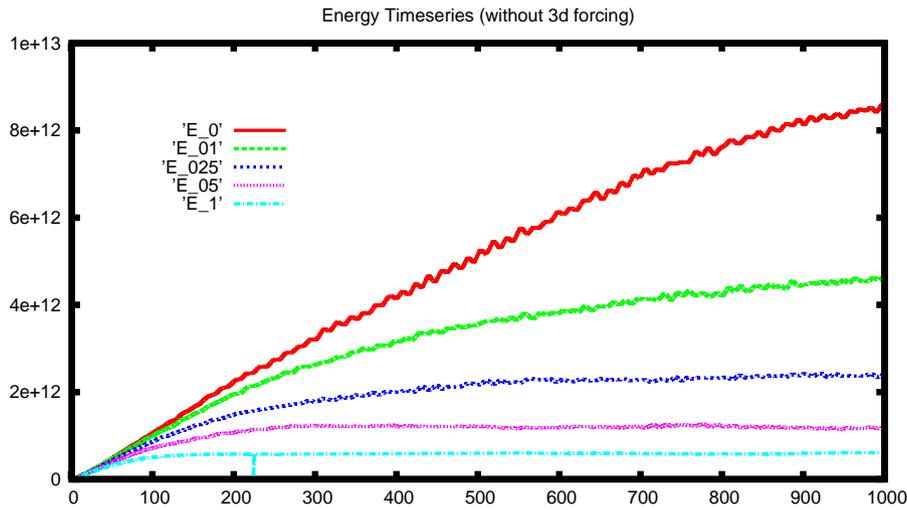
Green: Hyperviscosity

# Dependence on level of 3d forcing

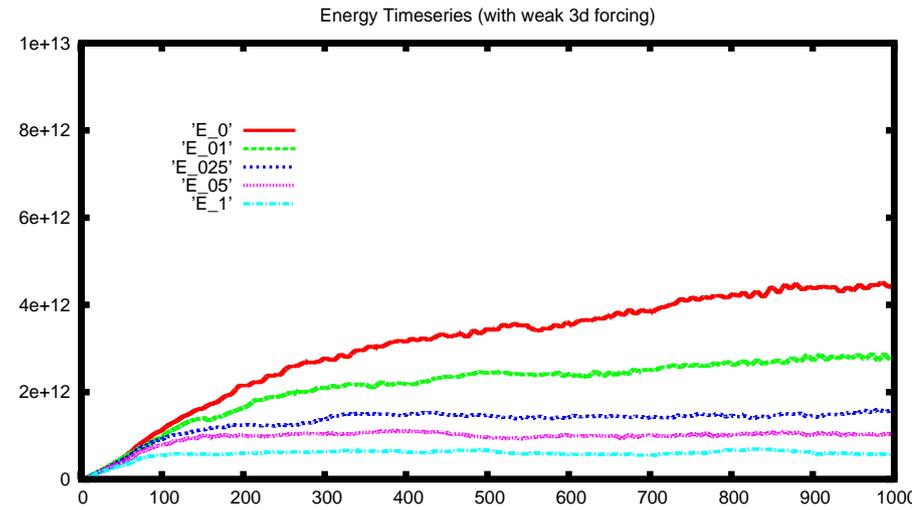


$$\tau_{\text{Rayleigh}} = (0, 1, 2.5, 5, 10) \times 10^{-8} \text{s}^{-1}$$

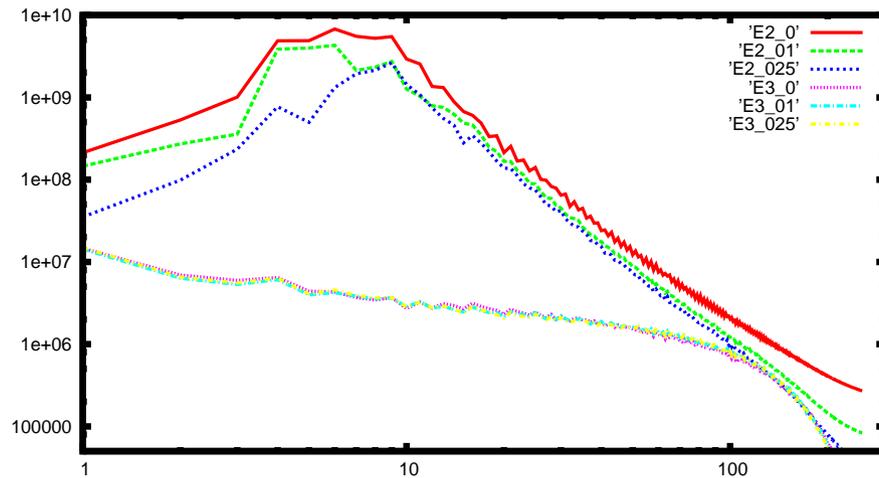
Without 3d forcing



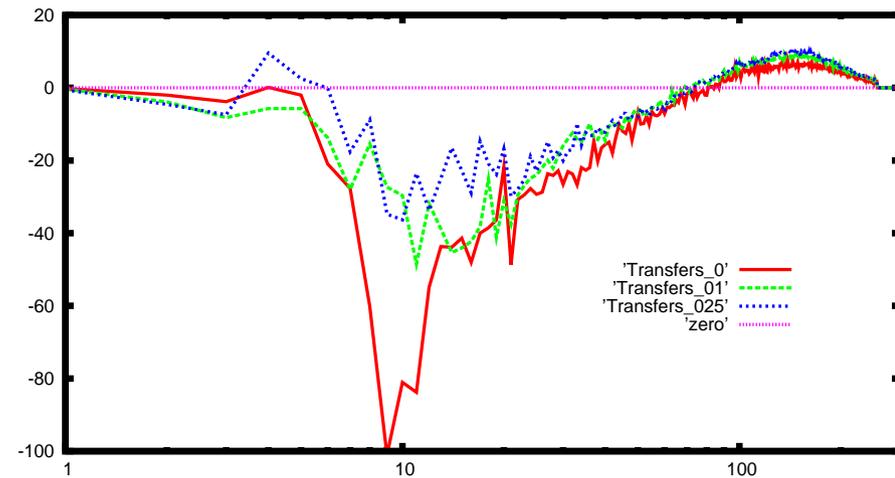
With (weak) 3d forcing



Energy Spectra (Different Rayleigh Drags)

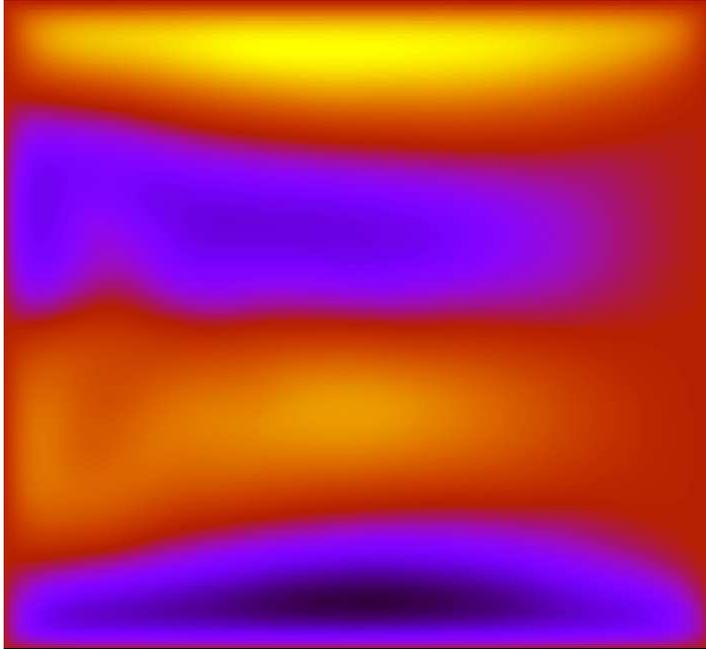


Energy Transfers (different Rayleigh drags)

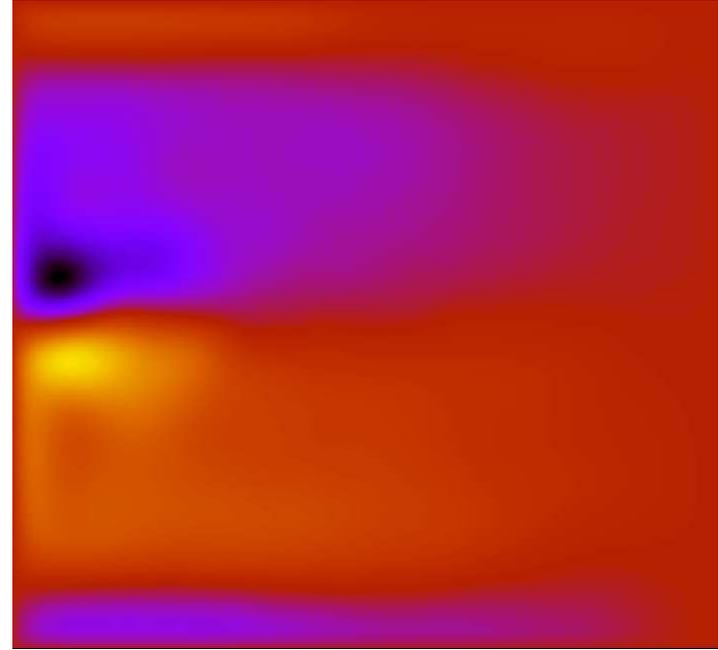


# Zero Rayleigh Drag (mean streamfunctions)

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Without 3d Forcing



With 3d Forcing

## Conclusions

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- When externally forced, geostrophic-to-inertial mode transfers can have a significant damping effect on the geostrophic motion.
- The effect is larger for stronger forcing of the inertial modes and weaker Rayleigh friction
- Stratified case: high vertical resolution is needed (i.e.,  $k_z > N/U$  or  $Nh/fL \leq O(1)$  for  $O(1)R_0$ )

## Rotating / Stratified Case

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- thin aspect ratio  $\rightarrow$  hydrostatic
- Horz. PGF  $\sim g\tilde{\rho}k_h/\rho_0k_z$
- HPGF small if

$$N^2 \ll k_z^2 U^2 \quad (\text{assuming } T \sim L/U)$$

or

$$N^2 k_h^2 \ll f^2 k_z^2 \quad (\text{assuming } R_0 \sim O(1))$$

# Model Equations

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$$\nabla^2 \psi_t + J(\psi, \nabla^2 \psi + \beta y) = -\overline{\mathbf{v}' \cdot \nabla \zeta'} - \overline{\mathbf{v}'_x \cdot \nabla v'} + \overline{\mathbf{v}'_y \cdot \nabla u'} + F - D$$

$$u'_t - (f + \nabla^2 \psi)v' - \zeta' \bar{v} = -(\bar{u}u' + \bar{v}v')_x + F^x - D^x$$

$$v'_t + (f + \nabla^2 \psi)u' - \zeta' \bar{u} = -(\bar{u}u' + \bar{v}v')_y - D^y$$

- C-grid
- Multigrid method for elliptic inversion (L-P Nadeau)
- $F^x = a(t)\sin(y/L)$  or  $F^x = a(t)\sin(y/L)e^{-0.25(x^2+y^2)/L^2}$
- $a(t)$  given by an Ornstein-Uhlenbeck process (with characteristic timescale of 1 day)