

Leslie Smith HW 3

#1 For 2-D non-rotating flow,

$$\begin{cases} \bar{u}_t + \underline{u} \cdot \nabla \underline{u} = -\nabla P & \text{① use } \\ \nabla \cdot \underline{u} = 0 & \text{②} \end{cases} \quad \begin{cases} u = -\frac{\partial \Psi}{\partial y} \\ v = \frac{\partial \Psi}{\partial x} \end{cases} \quad \text{and } w_3 = \nabla^2 \Psi$$

$$\text{For ①, use } \bar{u}_t + \bar{w} \bar{x} \bar{u} = -\nabla P^* \quad \text{③} \quad P^* = P + \frac{uu}{2}$$

$$\underline{u} \cdot \text{③}, \Rightarrow E = \frac{\underline{u} \cdot \underline{u}}{2}$$

$$\frac{\partial E}{\partial t} \text{ ④} = \underline{u} \cdot (-\nabla P^*) = -\nabla \cdot [P^* \underline{u}] + P \nabla \cdot \underline{u}$$

$$\int_V \frac{\partial E}{\partial t} = - \int_V -\nabla \cdot [P^* \underline{u}] = - \int_A P^* \underline{u} \cdot \hat{n} dA = 0$$

$$\Rightarrow E = \frac{\underline{u} \cdot \underline{u}}{2} \text{ is conserved}$$

use stream function Ψ , ① can be written as

$$\partial_t \nabla^2 \Psi + J(\Psi, \nabla^2 \Psi) = 0 \quad \text{⑤}$$

$$\nabla^2 \Psi \cdot \text{⑤}, \quad F = \frac{w_3^2}{2} = \frac{(\nabla^2 \Psi)^2}{2}, \quad - \int_V (\underline{u} \cdot \nabla \Psi) \nabla^2 \Psi = \int_V \nabla \cdot (\underline{u} \nabla \Psi) \nabla^2 \Psi \quad \text{dil}$$

$$\int_V \frac{\partial}{\partial t} F = - \int V J(\Psi, \nabla^2 \Psi) \cdot \nabla^2 \Psi \quad \text{[crossed out]} = 0$$

$$\text{so enstrophy } F = \frac{w_3^2}{2} \text{ is conserved.}$$

in Fourier, means $\sum_k E(k)$ and $\sum_k k^2 E(k)$ are conserved.

Using Kolmogorov's assumption, since energy is conserved,
let ε be the energy transfer rate,

$$E(k) = C \varepsilon^{1/3} k^{\beta}, \quad \text{since } [E] = \frac{\varepsilon^{1/3}}{t^{2/3}}, \quad [\varepsilon] = \frac{\varepsilon^2}{t^{5/3}}$$

$$\Rightarrow E(k) = C \varepsilon^{2/3} k^{-5/3}$$

now let η be enstrophy transfer rate, $[\eta] = [k^2 \varepsilon] = \frac{1}{t^{2/3}}$

$$\text{assume } k^2 E(k) = C_0 \eta^{1/3} k^{\beta},$$

$$\text{we have } k^2 E(k) = C_0 \eta^{2/3} k^{-1} \Rightarrow E(k) = C_0 \eta^{2/3} k^{-3}$$

now we have two scales, $E(k) \propto k^{-5/3}$, $E(k) \propto k^{-3}$

~~!!!!~~,

$$\text{For QG: } \left(\frac{\partial}{\partial t} + \underline{u}_H \cdot \nabla_H \right) \left(\nabla_H^2 + \frac{F^2}{N^2} \frac{\partial^2}{\partial z^2} \right) \Psi = 0$$

$$\text{where } u = -\frac{\partial \Psi}{\partial y}, \quad v = \frac{\partial \Psi}{\partial x}, \quad \theta = -\frac{F}{N} \frac{\partial^2 \Psi}{\partial z^2} \quad \text{⑥}, \quad \Psi = \left(\nabla_H^2 + \frac{F^2}{N^2} \frac{\partial^2}{\partial z^2} \right) \Psi$$

$$\Psi \cdot \text{⑥}, \quad \text{let } E = \frac{\underline{u}_H \cdot \underline{u}_H}{2} + \frac{\theta \theta}{2} = \frac{\nabla_H^2 \Psi}{2} + \frac{F^2}{N^2} (\Psi_z)^2 \quad \text{[crossed out]}$$

$$\Rightarrow \frac{\partial}{\partial t} E = - \int_V \underline{u}_H \cdot \nabla_H (\nabla_H^2 + \frac{F^2}{N^2} \frac{\partial^2}{\partial z^2}) \Psi \cdot \Psi$$

$$\therefore = - \int_V (\underline{u}_H \cdot \nabla_H g) \frac{\hat{z}}{4} = - \int_V \nabla_H \cdot (\underline{u}_H g \frac{\hat{z}}{4}) = \int_A \cancel{\underline{g} \cdot \underline{n}} \underline{u}_H \cdot \hat{n} dA \\ = 0$$

$$(\nabla_H \cdot (\underline{u}_H g) \frac{\hat{z}}{4}) = (\cancel{\underline{g} \cdot \underline{u}_H}) g \frac{\hat{z}}{4} + \underline{u}_H \cdot \nabla (g \frac{\hat{z}}{4}) = (\underline{u}_H \cdot \nabla g) \frac{\hat{z}}{4} + \cancel{(\underline{u}_H \cdot \cancel{\nabla_H^2 g})} \cdot g \\ = (\underline{u}_H \cdot \nabla g) \frac{\hat{z}}{4}$$

so Energy $E = \frac{\underline{u}_H \cdot \underline{u}_H + \theta^2}{2}$ is conserved.

For ⑤, $(\frac{\partial}{\partial t} + \underline{u}_H \cdot \nabla_H) g = 0$, $g \cdot \underline{g} \cdot \underline{g}$.

$$\text{let } F = \frac{\theta^2}{2} = \left[\left(\nabla_H^2 + \frac{F^2}{N^2} \frac{\partial^2}{\partial z^2} \right) \frac{\hat{z}}{4} \right]^2,$$

$$\text{we have } \frac{\partial}{\partial t} F = - \int_V (\underline{u}_H \cdot \nabla_H g) g = - \int_V \nabla_H \cdot (\underline{u}_H g^2) = 0$$

so potential enstrophy $F = \left(\left(\nabla_H^2 + \frac{F^2}{N^2} \frac{\partial^2}{\partial z^2} \right) \frac{\hat{z}}{4} \right)^2$ is conserved.

$$\text{now if rescale } \zeta = \frac{N}{F} z \Rightarrow d\zeta = \frac{N}{F} dz$$

we will have the ~~similar~~ similar result as 2-D flow

$$E(K) = \frac{1}{2}(u^2 + v^2 + \theta^2) \text{ and } F(K) = \frac{1}{2} K^2 E(K) \text{ are conserved}$$

so with similar argument, $E(K) \propto K^{-5/3}$ and $E(K) \propto K^{-3}$

$$\#2: \text{Since } C_{Kpq} = \frac{iN(p \times q) \cdot \hat{z}}{\sigma_K K \sigma_p p \sigma_q q} (\bar{\sigma}_q^2 q^2 - \bar{\sigma}_p^2 p^2), \\ \text{with } \sigma_K K = (N^2 k_n^2 + F^2 k_z^2)^{1/2}$$

Now if $K + p + q = 0$.

$$(q \times k) \cdot \hat{z} = (q \times (-p - q)) \cdot \hat{z} = -(q \times p) \cdot \hat{z} = (p \times q) \cdot \hat{z}$$

$$\text{so } p \times q = q \times k = k \times p$$

$$\text{so } C_{Kpq} + C_{pqk} + C_{qkp}$$

$$= \frac{iN(p \times q) \cdot \hat{z}}{\sigma_K K \sigma_p p \sigma_q q} (\bar{\sigma}_q^2 q^2 - \bar{\sigma}_p^2 p^2) + \frac{iN(q \times k) \cdot \hat{z}}{\sigma_p p \sigma_q q \sigma_k k} (\bar{\sigma}_q^2 k^2 - \bar{\sigma}_q^2 q^2)$$

$$+ \frac{iN(k \times p) \cdot \hat{z}}{\sigma_q q \sigma_k k \sigma_p p} (\bar{\sigma}_p^2 p^2 - \bar{\sigma}_k^2 k^2)$$

$$= \frac{iN(p \times q) \cdot \hat{z}}{\sigma_K K \sigma_p p \sigma_q q} (\bar{\sigma}_q^2 q^2 - \bar{\sigma}_p^2 p^2 + \bar{\sigma}_k^2 k^2 - \bar{\sigma}_q^2 q^2 + \bar{\sigma}_p^2 p^2 - \bar{\sigma}_k^2 k^2) = 0$$

group #6

$$\text{and } \bar{\tau}_k^2 k^2 C_{kpq} + \bar{\tau}_p^2 p^2 C_{pqp} + \bar{\tau}_q^2 q^2 C_{qkp}$$

$$= \frac{iN(p \times q) \cdot \hat{z}}{\bar{\tau}_k k \bar{\tau}_p p \bar{\tau}_q q} \left[\bar{\tau}_k^2 k^2 (\bar{\tau}_q^2 q^2 - \bar{\tau}_p^2 p^2) + \bar{\tau}_p^2 p^2 (\bar{\tau}_k^2 k^2 - \bar{\tau}_q^2 q^2) + \bar{\tau}_q^2 q^2 (\bar{\tau}_p^2 p^2 - \bar{\tau}_k^2 k^2) \right]$$

$$= 0$$

$$\text{so } C_{kpq} + C_{pqp} + C_{qkp} = 0$$

$$\text{and } \bar{\tau}_k^2 k^2 C_{kpq} + \bar{\tau}_p^2 p^2 C_{pqp} + \bar{\tau}_q^2 q^2 C_{qkp} = 0$$

in one trial, $\underline{k} + \underline{p} + \underline{q} = 0$.

$$\begin{aligned} a_k^* \cdot \frac{\partial}{\partial t} a_k &= C_{kpq} a_p^* a_q^* \\ a_p^* \cdot \frac{\partial}{\partial t} a_p &= C_{pqp} a_q^* a_k^* \\ a_q^* \cdot \frac{\partial}{\partial t} a_q &= C_{qkp} a_k^* a_p^* \end{aligned} \Rightarrow \begin{cases} \frac{\partial}{\partial t} \left(\frac{a_k^2}{2} \right) = C_{kpq} a_k^* a_p^* a_q^* \\ \frac{\partial}{\partial t} \left(\frac{a_p^2}{2} \right) = C_{pqp} a_p^* a_q^* a_k^* \\ \frac{\partial}{\partial t} \left(\frac{a_q^2}{2} \right) = C_{qkp} a_k^* a_p^* a_q^* \end{cases}$$

$$\text{take sum, } \Rightarrow \frac{\partial}{\partial t} \left(\frac{a_k^2 + a_p^2 + a_q^2}{2} \right) = 0$$

$\Rightarrow |a_k^0|^2 + |a_p^0|^2 + |a_q^0|^2$ is conserved

similarly, $\bar{\tau}_k^2 k^2 |a_k^0|^2 + \bar{\tau}_p^2 p^2 |a_p^0|^2 + \bar{\tau}_q^2 q^2 |a_q^0|^2$ is conserved